

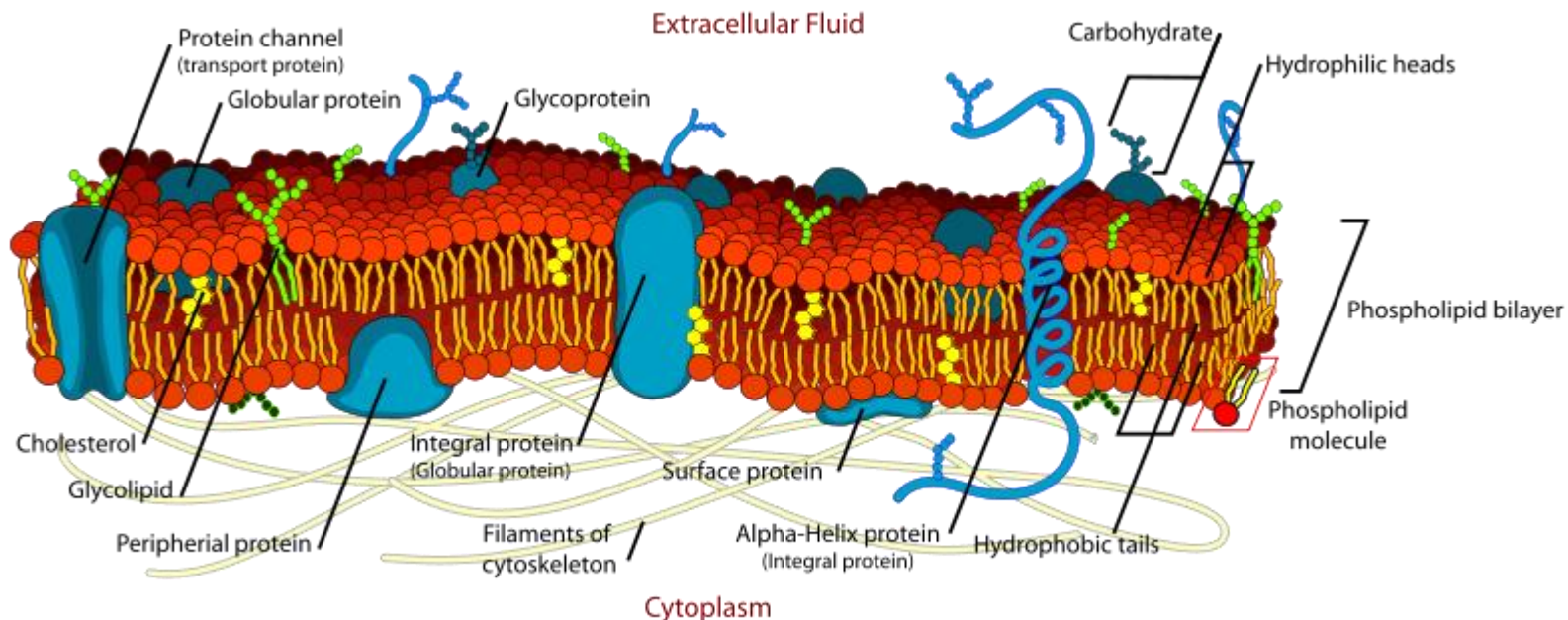
# Lipid Bilayers and Membrane Dynamics: Insight into Thickness Fluctuation

Andrea C. Woodka,<sup>1</sup> Paul D. Butler,<sup>1</sup> Lionel Porcar,<sup>2</sup> Bela Farago,<sup>2</sup> and  
Michihiro Nagao,<sup>1,2,3</sup>

<sup>1</sup> NIST Center for Neutron Research, National Institute of Standards and Technology, Gaithersburg, MD 20899-6102, USA,

<sup>2</sup> Institut Laue Langevin, 6 rue Jules Horowitz, BP 156-38042, Grenoble Cedex 9, France,

<sup>3</sup> Center for Exploration of Energy and Matter, Indiana University, Bloomington, IN 47408, USA

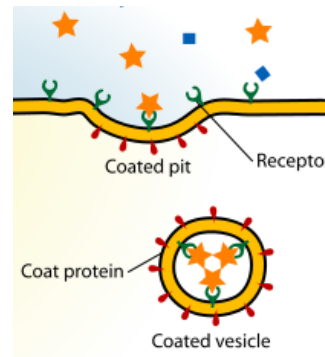


# Motivation

Lipid membranes are self-assembled highly flexible structures that have the ability to undergo an array of conformational and dynamic transitions which are essential for many biological functions.

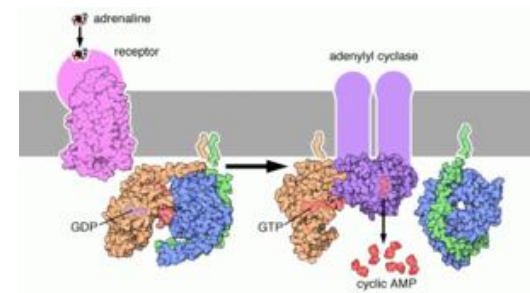
## Microscopic length scale:

Membrane stiffness and fluidity have been shown to have a large impact on cellular uptake and release.<sup>(2)</sup>



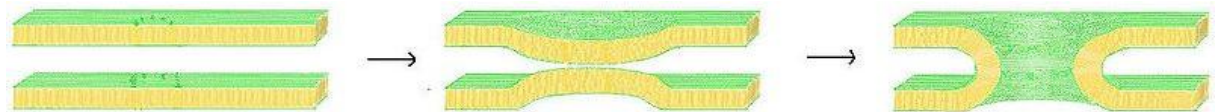
## Spectroscopic length scale:

Cell signal transduction is affected by molecular lateral diffusion within the lipid membrane.<sup>(1)</sup>



## Intermediate length scale :

Membrane thickness fluctuations have been proposed as a mechanism for pore formation.<sup>(3)</sup>

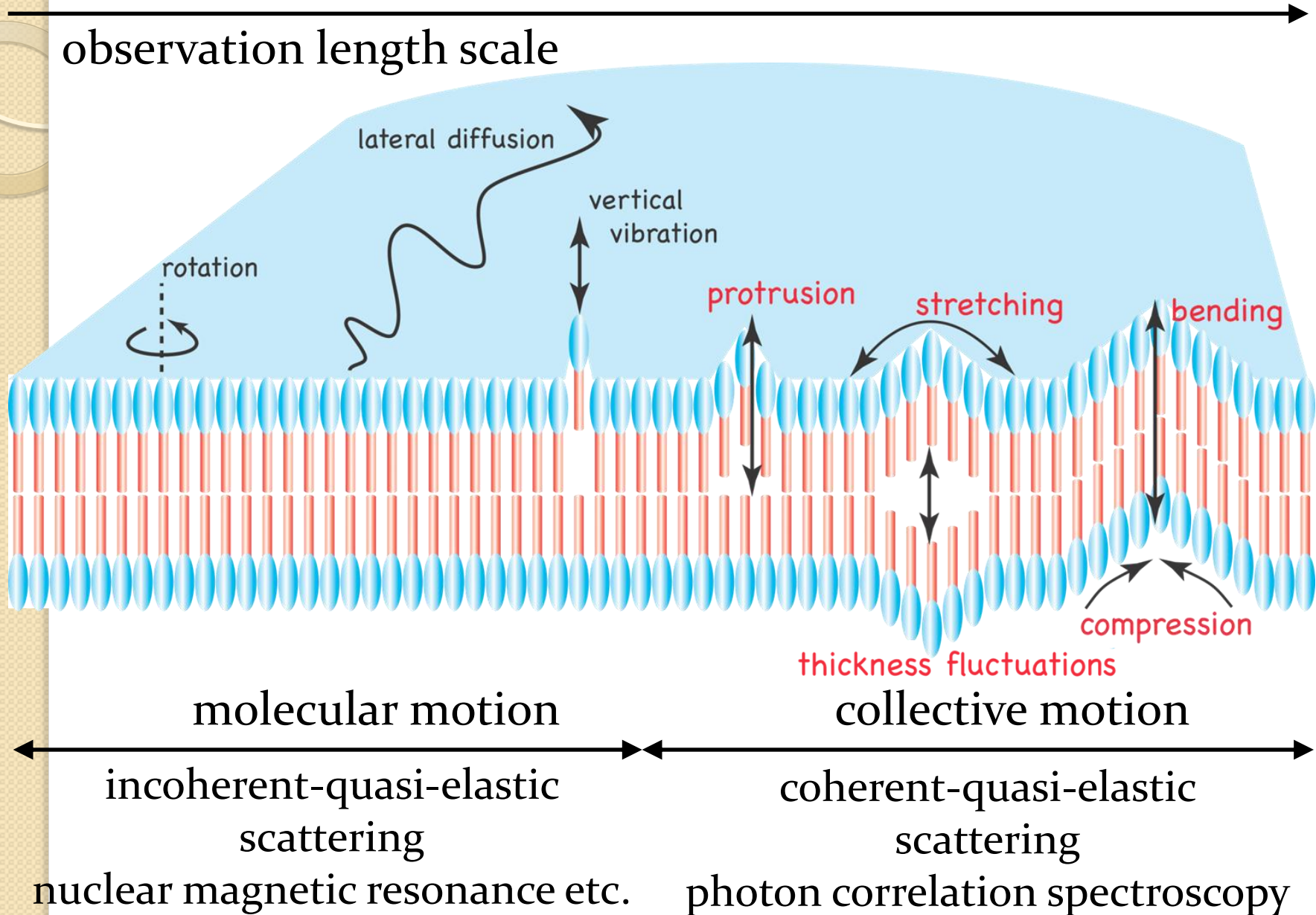


(1) D Marguet, et al. EMBO J, 25, 3446 (2006)

(2) P. Weber, et. al. Adv Med Eng, 114, 377(2006)

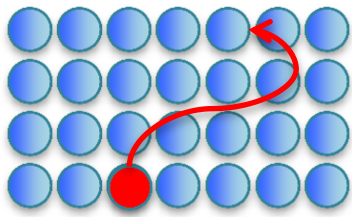
(3) L. Movilenu, et al. Bull Math Biol, 58, 1231 (2006)

# Membrane Dynamics

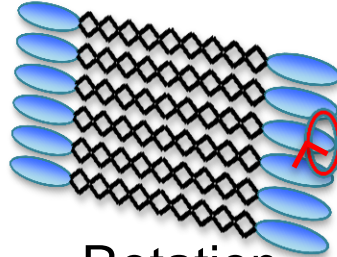


# The Dynamics in Lipid Bilayers

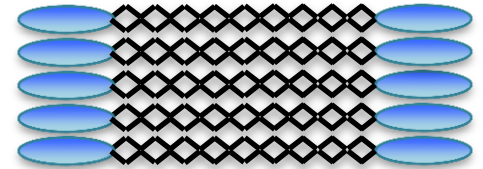
Molecular movements (incoherent movements of molecules)



Lateral diffusion

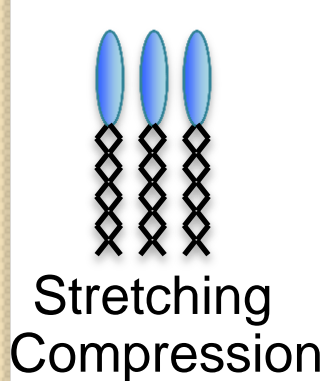


Rotation



Vertical vibration

Coherent movements of molecular assemblies

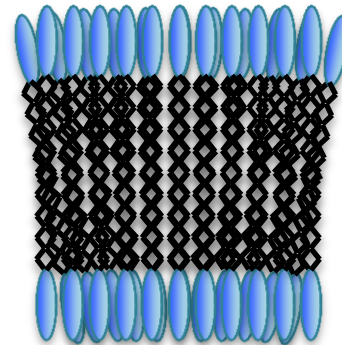


Stretching  
Compression

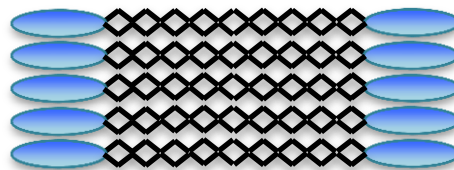


Stretching

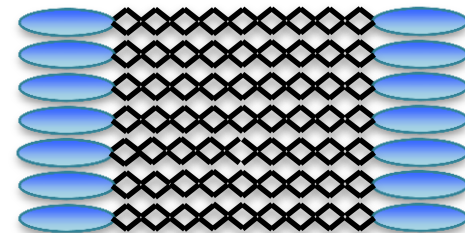
Compression



Bending



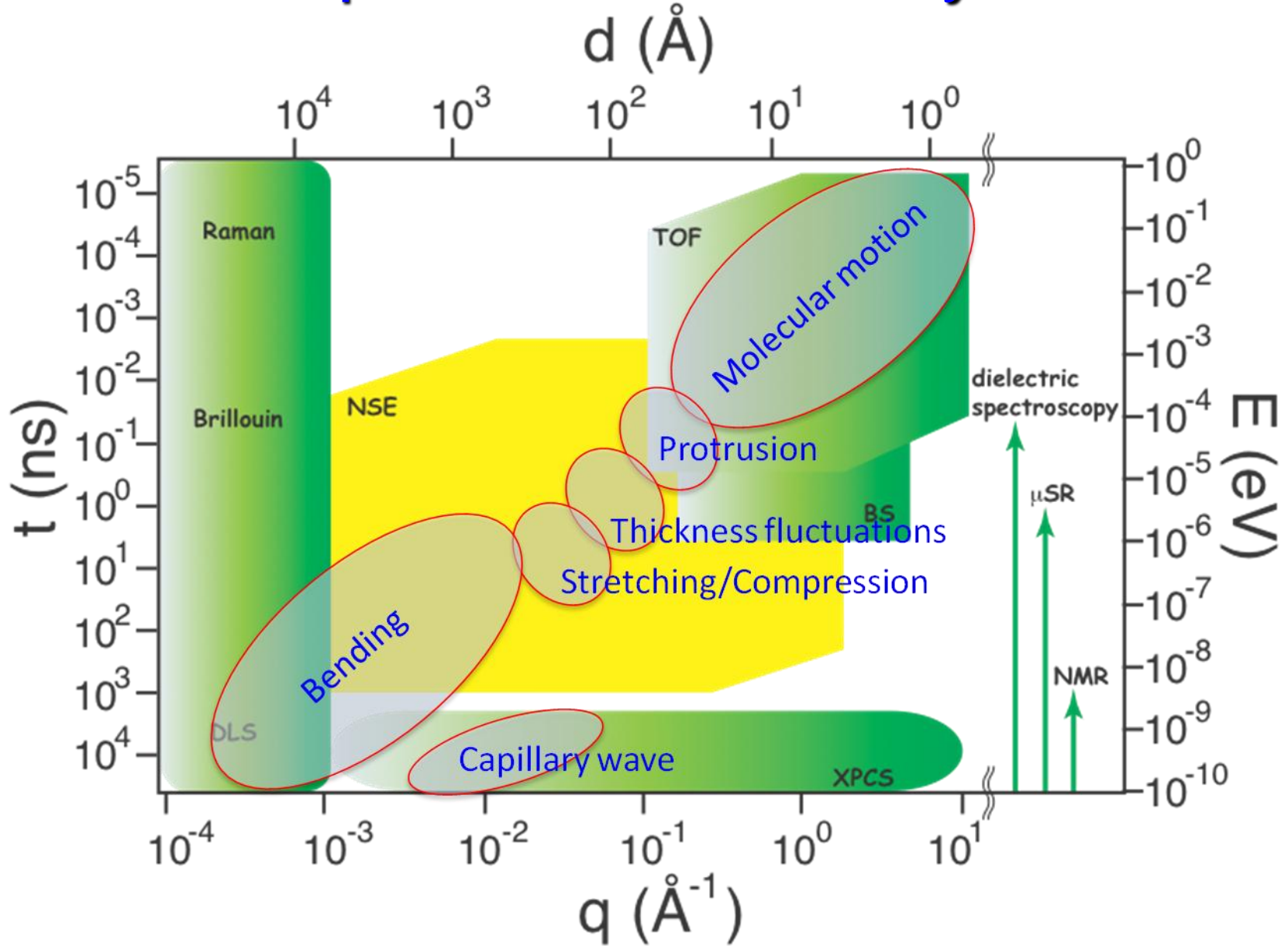
Protrusion



Thickness fluctuations

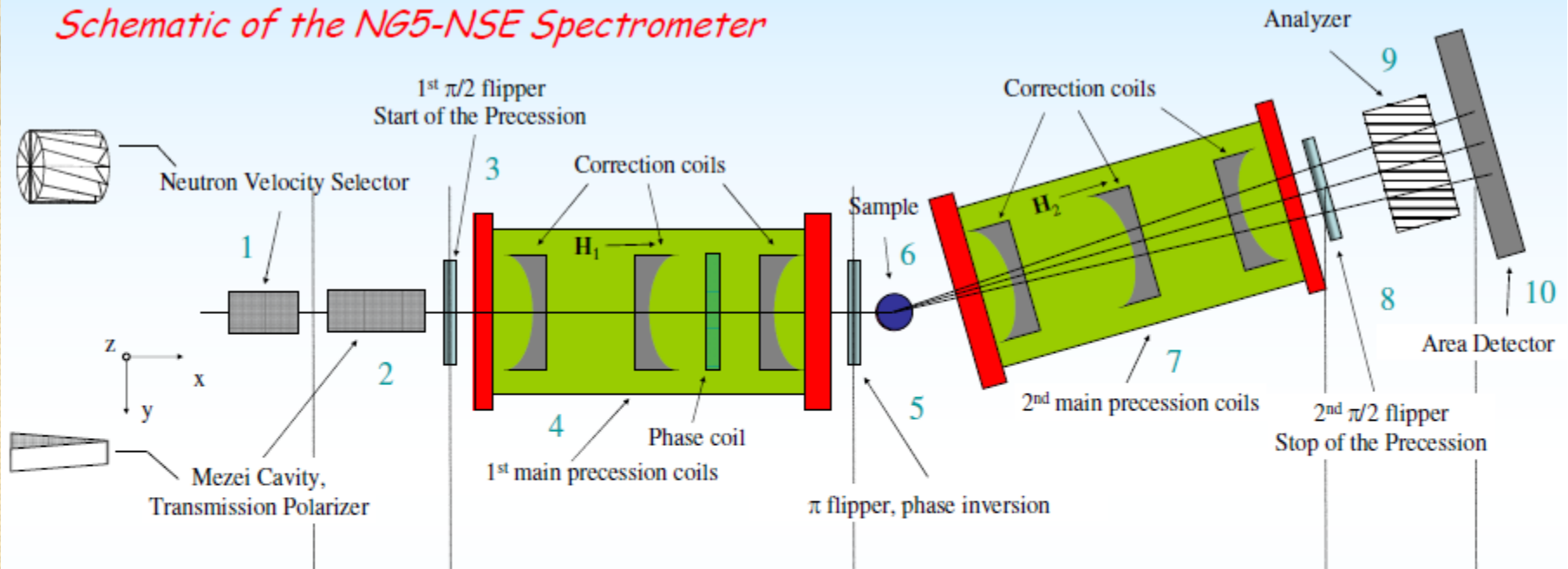


# Techniques to Measure Dynamics

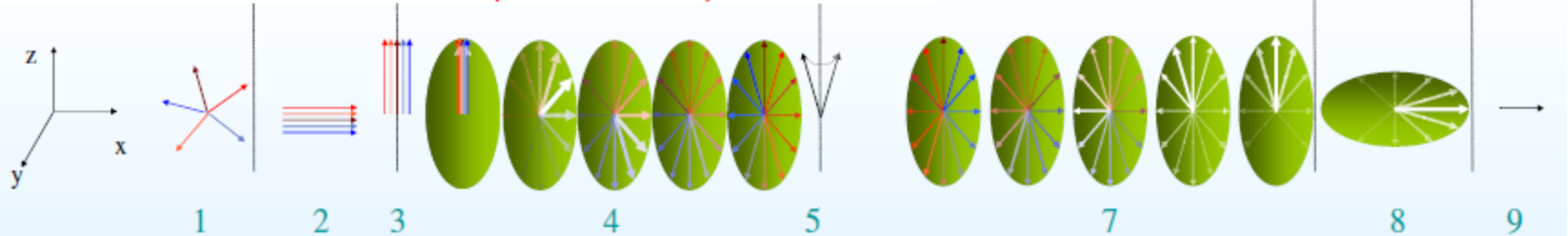


# Neutron Spin Echo

*Schematic of the NG5-NSE Spectrometer*



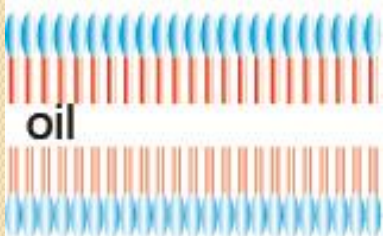
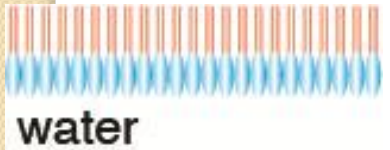
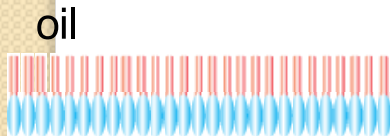
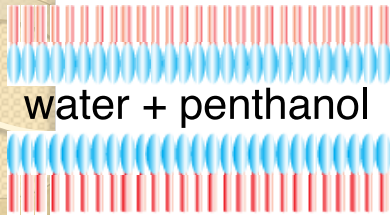
*Motion of the neutron beam spins in the spectrometer*



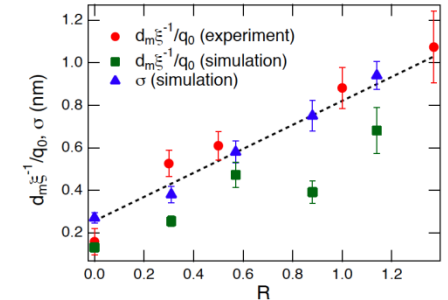
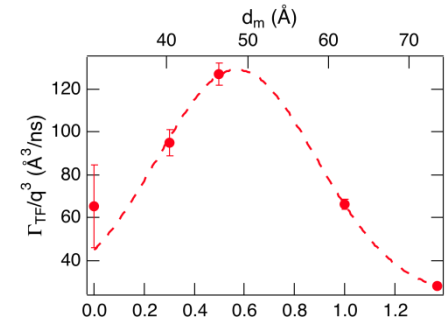
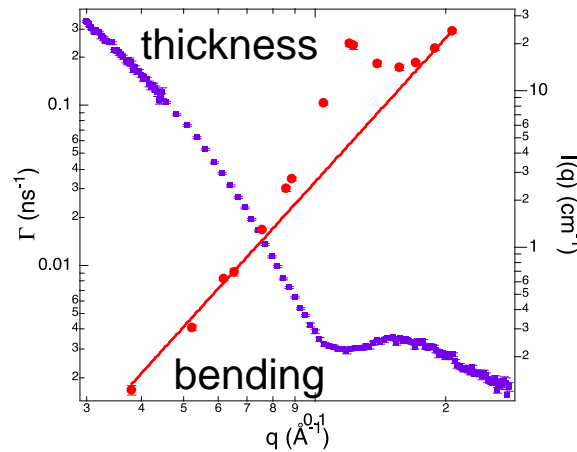
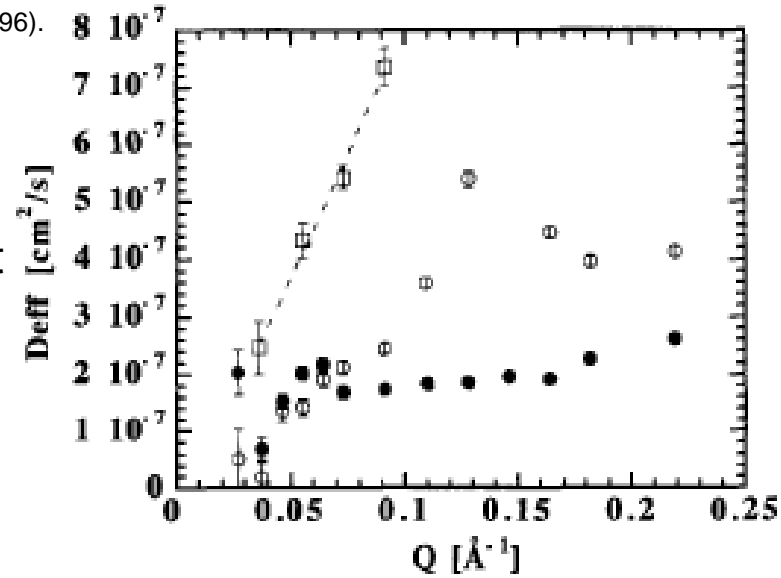
Elastic Scattering

# Surfactant Membranes

Farago et al, *Physica B* **213&214**, 712 (1995).; Farago, *Physica B* **226**, 51 (1996).



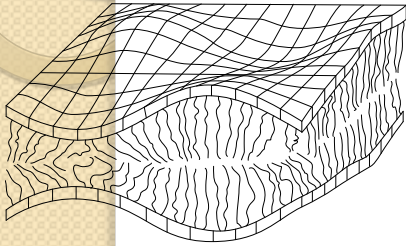
- NSE measured  $Q$  dependence of  $D_{\text{eff}}$
- showing an increase at the membrane thickness length scale



# Thickness fluctuations in lipid bilayers (theoretical studies)

## Breathing model of a lipid bilayer by Miller

Miller, *Top. Bioelectrochem. Bioenerg.* **4**, 161 (1981).; Bach and Miller, *Biophys. J.* **29**, 183 (1980).; Miller, *Biophys. J.* **45**, 643 (1984).



Amplitude of the fluctuations reaches  $\approx 15 \text{ \AA}$  or more from the geometrical constraints (volume conservation)

## Thickness fluctuations by Hladky and Gruen

Hladky and Gruen, *Biophys. J.* **38**, 251 (1982).

Thickness fluctuations occur, but the amplitude is small.

Long wavelength fluctuation amplitude is negligible

Short wavelength fluctuations ( $< 30 \text{ \AA}$ ) are severely limited

Intermediate wavelength fluctuation amplitude  $< 10 \text{ \AA}$

## Deformation free energy of bilayer membranes by Huang

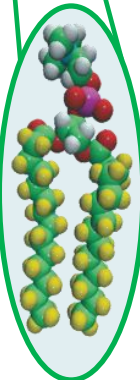
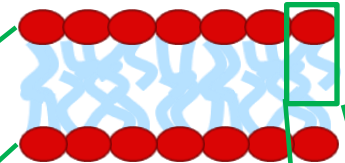
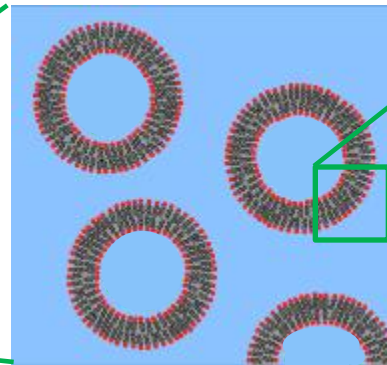
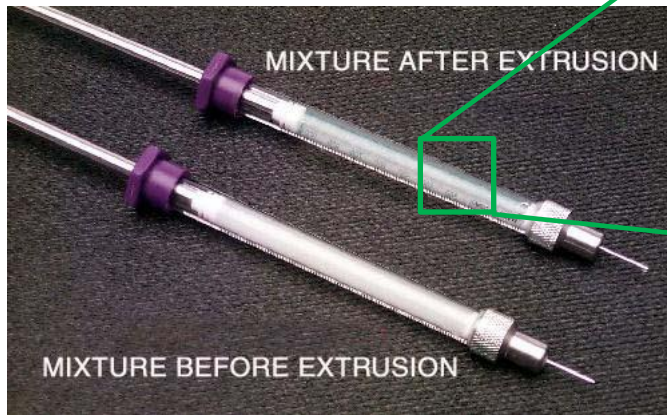
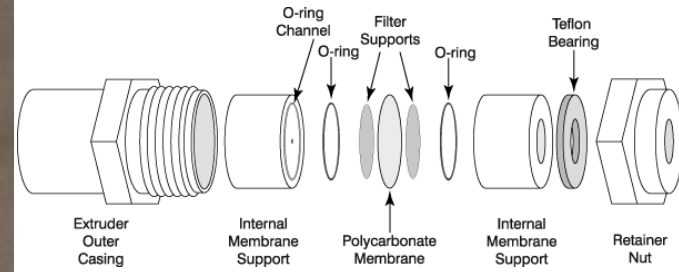
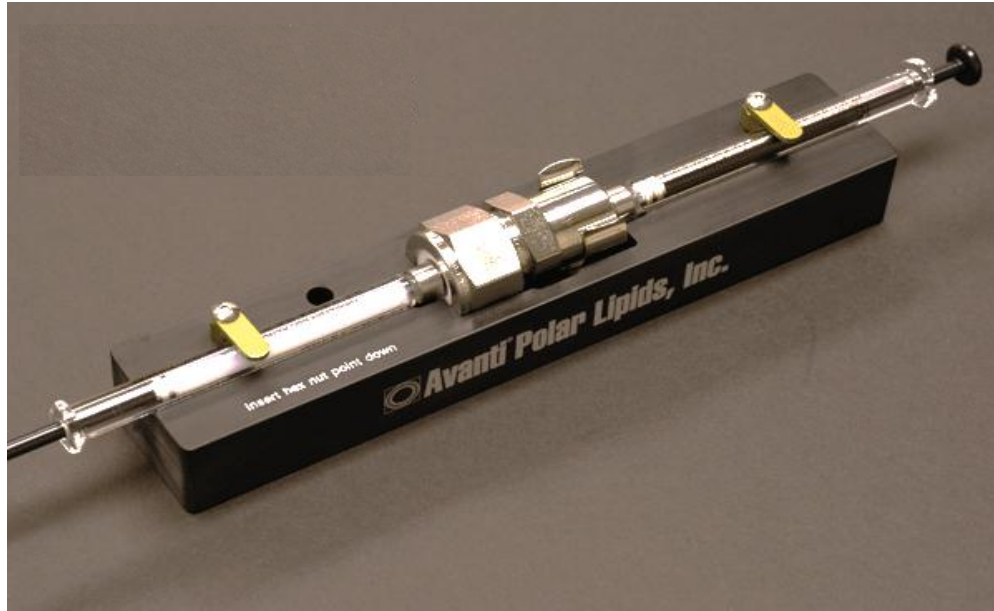
Huang, *Biophys. J.* **50**, 1061 (1986).

$$\sqrt{\langle \text{Amplitude}^2 \rangle} = \frac{k_B T}{\pi \gamma} \left\{ \tan^{-1} \left[ \frac{16\pi^2 K_1 + h\gamma}{2h\sqrt{K_1 \bar{B}}} \right] - \tan^{-1} \left[ \frac{\gamma}{2\sqrt{K_1 \bar{B}}} \right] \right\} \approx 4.5 \text{ \AA}$$

Theoretically, thickness fluctuations exist, their amplitude is very small



# Bilayer Preparation



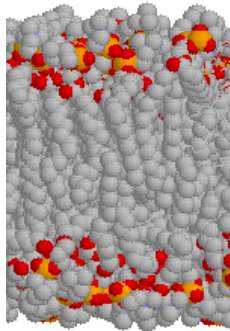
(D54)DMPC = C14

(D62)DPPC = C16

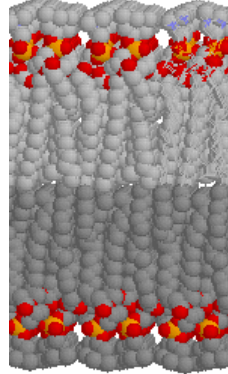
(D70)DSPC = C18

# Phospholipid Melting Temperature

Fluid phase  
above  $T_m$

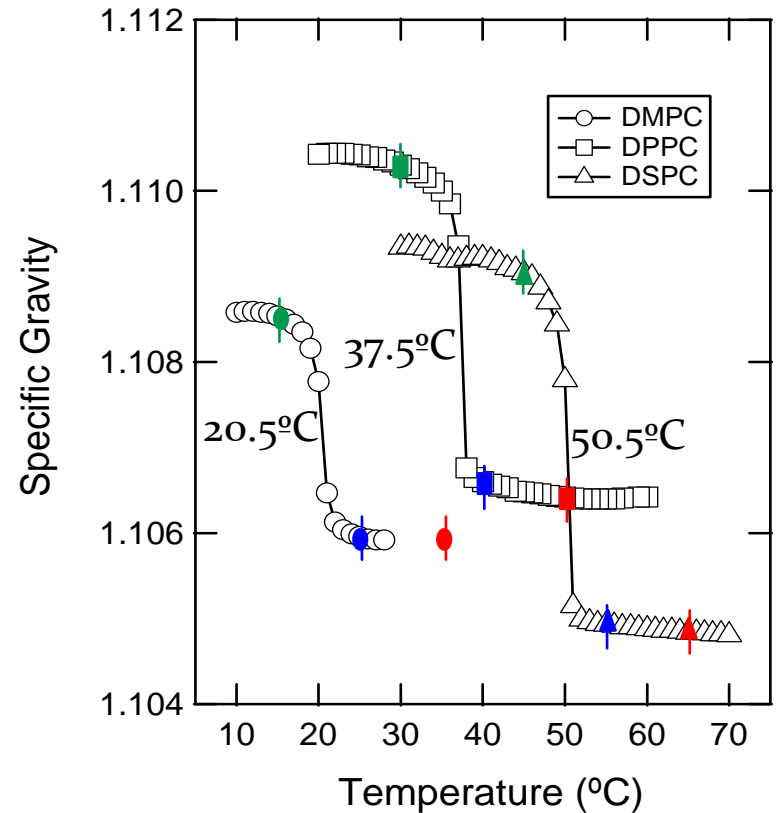


Gel phase,  
below  $T_m$

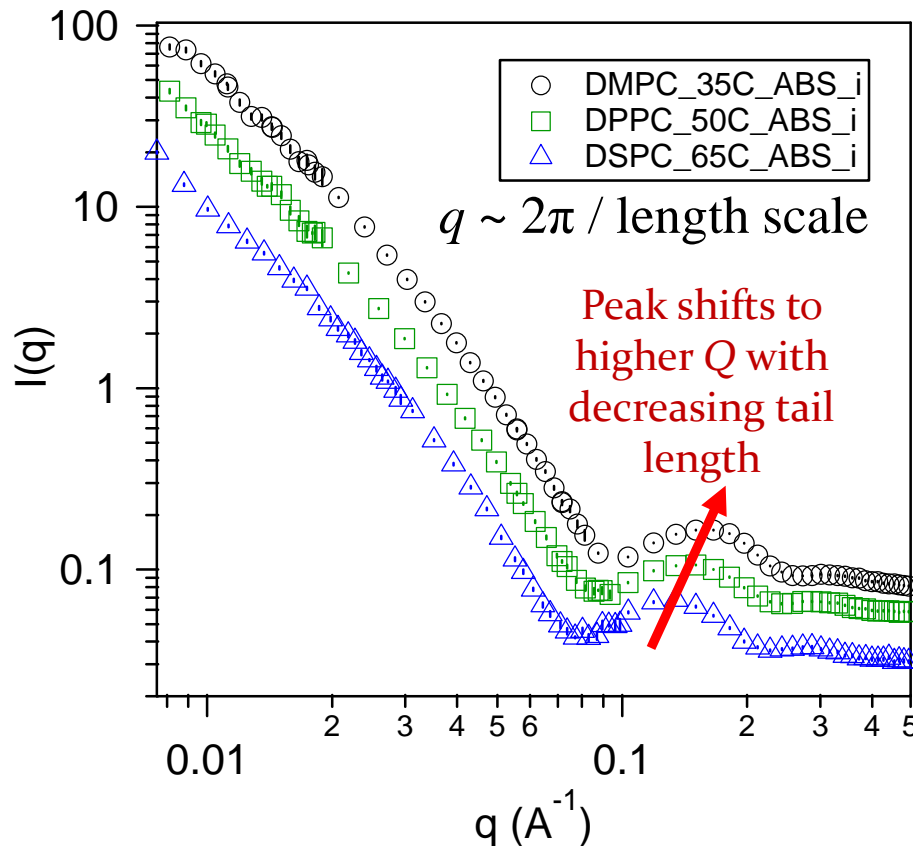
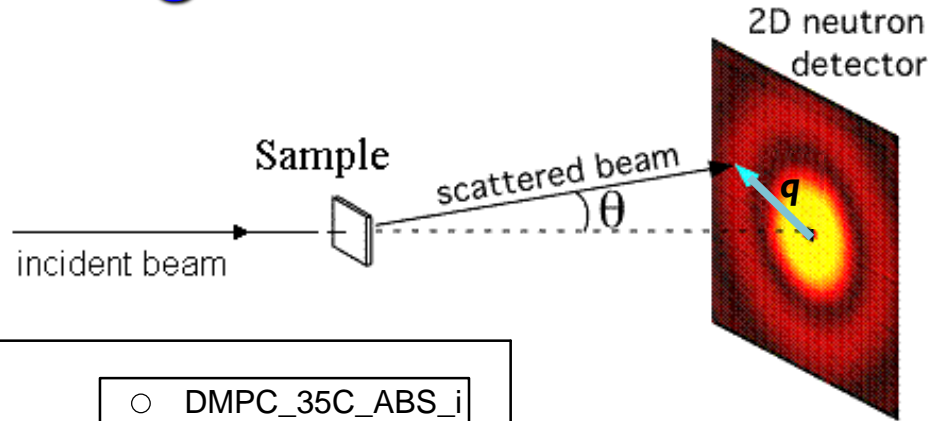


Above  $T_m$  lipid tails highly fluid, disordered and in constant motion

Below  $T_m$  transition to a gel state, tails are fully extended with highly ordered packing



# Small Angle Neutron Scattering



\*Scattering intensity is offset to highlight peak shift

SANS and NSE are complementary techniques

## SANS

static “snapshot”  
elastic scattering

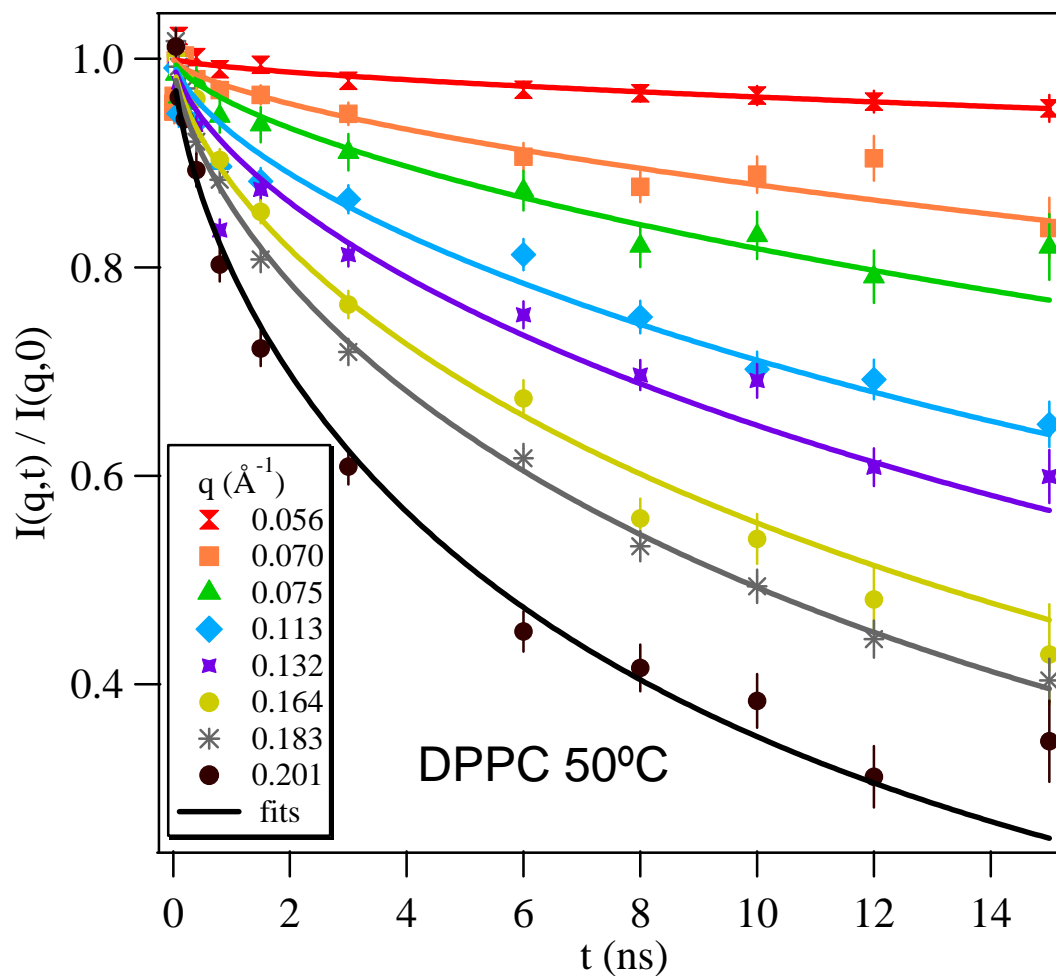
$$I(q) = \int S(q, \omega) d\omega$$

## NSE

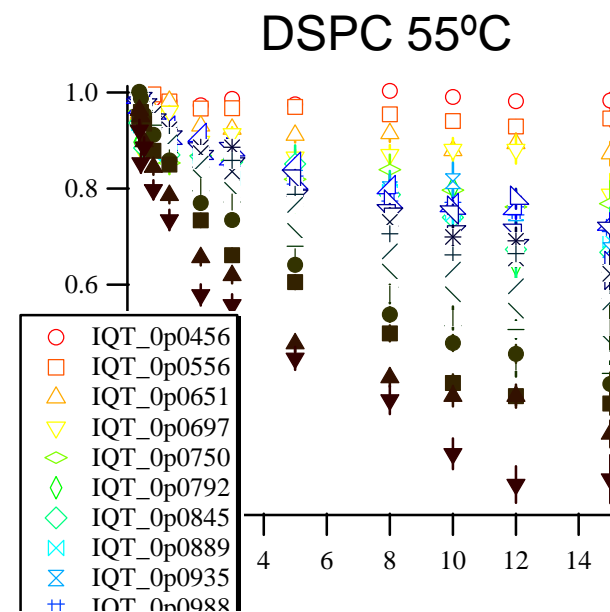
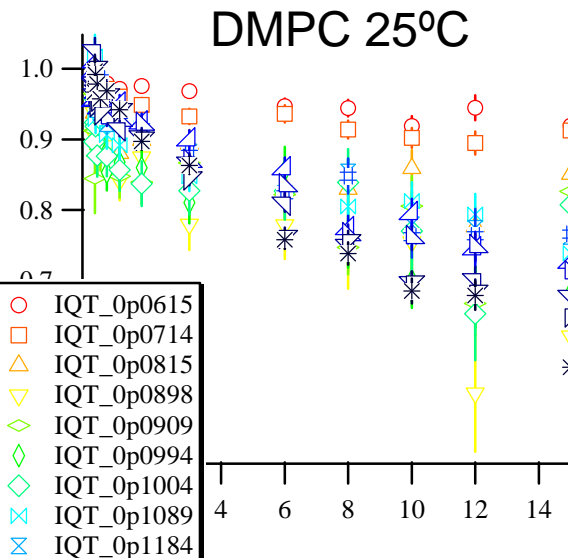
dynamic “snapshot”  
quasielastic scattering

$$I(q, t) = \int S(q, \omega) \cos(\omega t) d\omega$$

# NSE : $I(q,t)$



$$\frac{I(q,t)}{I(q,0)} = \exp[-(\Gamma t)^{2/3}]$$



# Surfactant Membrane Dynamic (bending)

## Helfrich bending energy

Helfrich, *Z. Natureforsch.* **28**, 693 (1973).

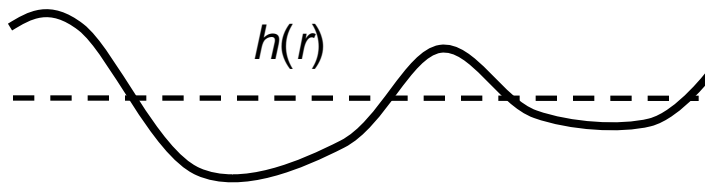
Assuming the membrane is thin enough sheet, which is undulating

## Zilman-Granek theory

Zilman and Granek, *Phys. Rev. Lett.* **77**, 4788 (1996).

Zilman and Granek, *Chem. Phys.* **284**, 195 (2002).

Dynamics of a planar non-interacting Helfrich sheet



$$\frac{I(q,t)}{I(q,0)} = \exp\left[-(Gt)^{2/3}\right]$$

$$H = \frac{1}{2} \kappa \int d^2r (\nabla^2 h(\vec{r}))^2 = \frac{1}{2\xi^2} \sum_{\vec{q}} \kappa q^4 h_{\vec{q}} h_{-\vec{q}}$$

$$G = 0.025 g_c \frac{\kappa k_B T \eta^{1/2}}{h} q^3$$

$\kappa$ : bending modulus

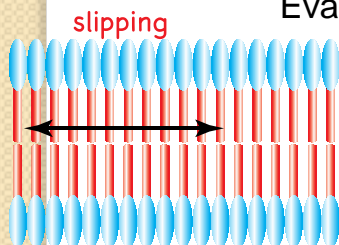
$\eta$ : solvent viscosity

## Watson-Brown theory

Watson and Brown, *Biophys. J.* **98**, L09 (2010).

Extension of ZG theory including slipping of each monolayer

Evans and Yeung, *Chem. Phys. Lipids.* **73**, 39 (1994).; Seifert and Langer, *Europhys. Lett.* **23**, 7



Inter-monolayer friction plays a role, where lateral compressibility  $k_m$  of membrane appears in dynamical equation

$$\kappa \rightarrow \tilde{\kappa} = \kappa + 2d^2 k_m$$



# Fitting $I(q, t)$

Bending motion is explained as a single membrane dynamics model

$$\frac{I(q, t)}{I(q, 0)} = \exp \left[ - (\Gamma t)^\beta \right]$$

Zilman and Granek, *Phys. Rev. Lett.* **77**, 4788 (1996).; Zilman and Granke, *Chem. Phys.* **184**, 195 (2002).

$\Gamma$ : decay rate,  $\beta=2/3$

$$\frac{\Gamma_{\text{Bend}}}{q^3} = 0.025\alpha \sqrt{\frac{k_B T}{\tilde{\kappa}} \frac{k_B T}{\eta_{\text{D}_2\text{O}}}}$$

$\tilde{\kappa}$ : effective bending modulus,  
 $\eta$ : solvent viscosity,  $\alpha \approx 1$

Considering slipping friction

$$\tilde{\kappa} = \kappa + 2d^2 k_m$$

Watson and Brown, *Biophys. J.* **98**, L9 (2010).

$$k_m = \frac{24\kappa}{d_t}$$

Rawicz et al., *Biophys. J.* **79**, 328 (2000).

$$\frac{\Gamma_{\text{Bend}}}{q^3} = 0.0058 \sqrt{\frac{k_B T}{\kappa} \frac{k_B T}{\eta_{\text{D}_2\text{O}}}}$$

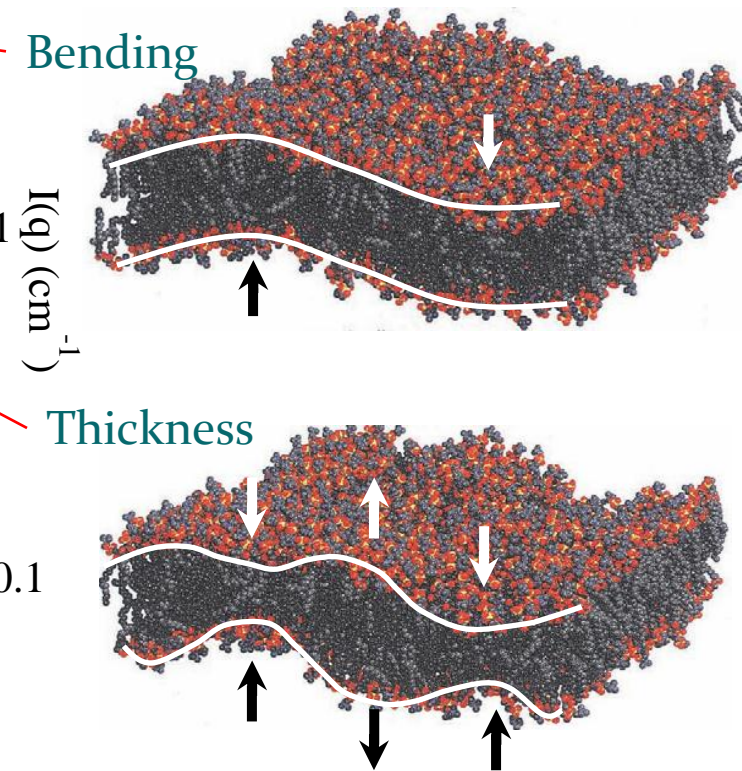
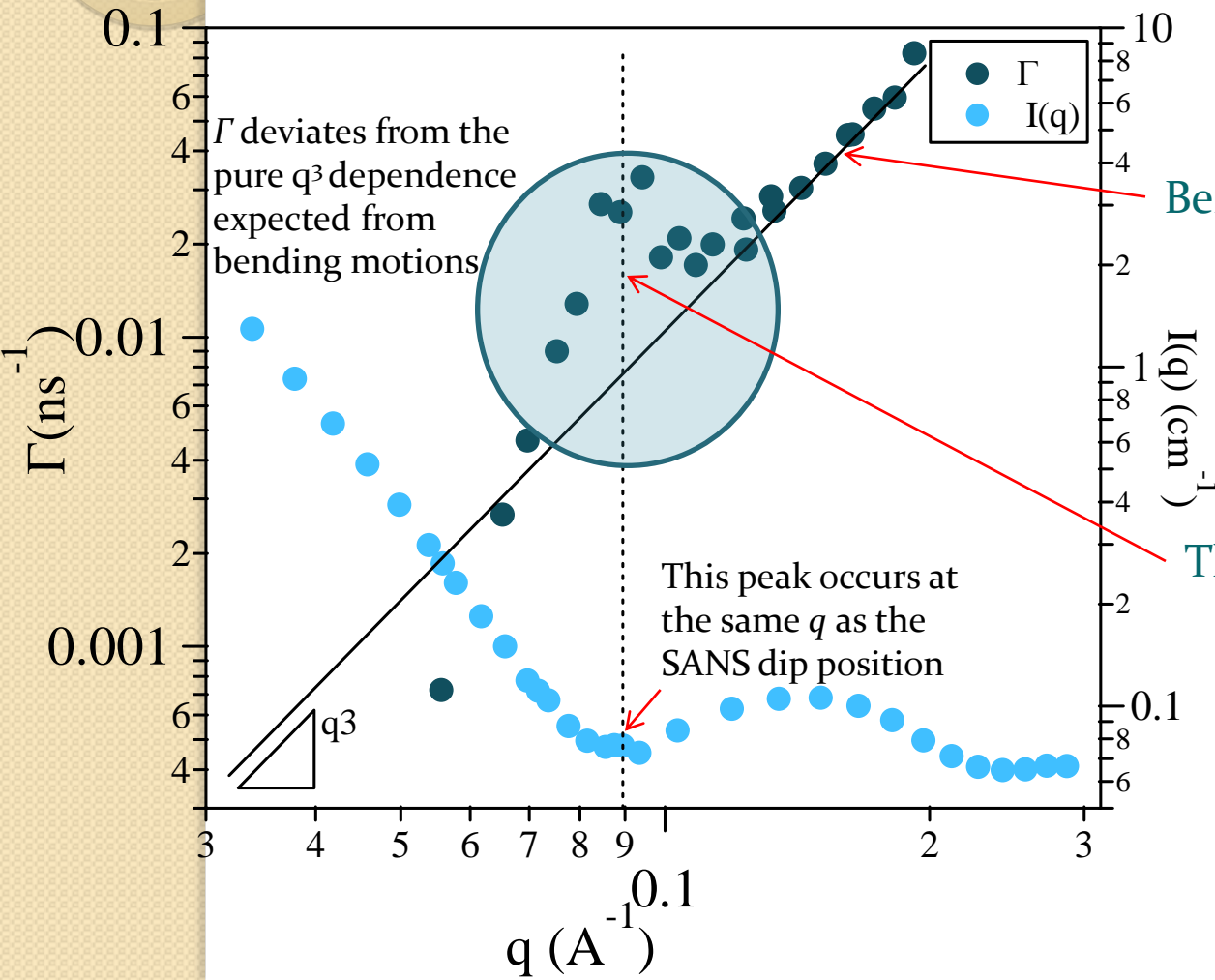
Lee et al., *Phys. Rev. Lett.* **105**, 038101 (2010).

# NSE : $\Gamma$ vs. $q$ DPPC @ 50°C

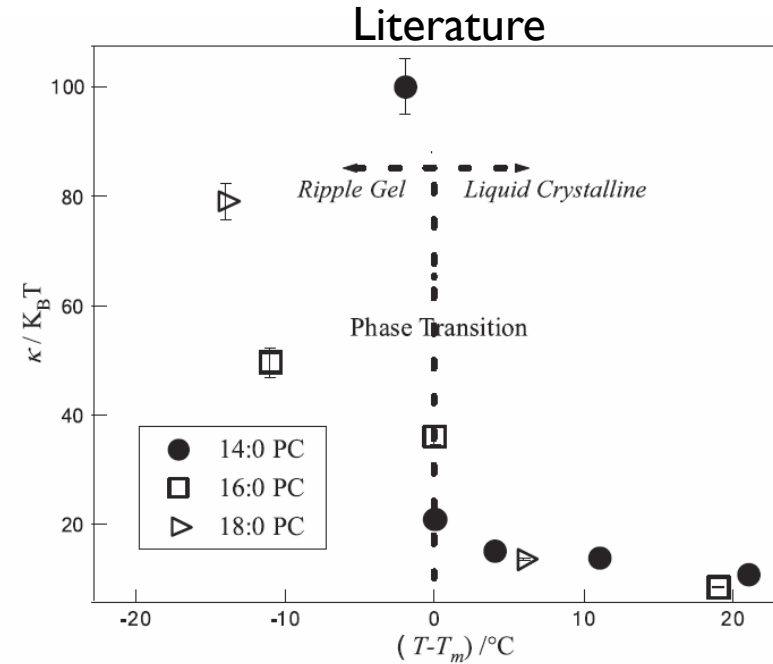
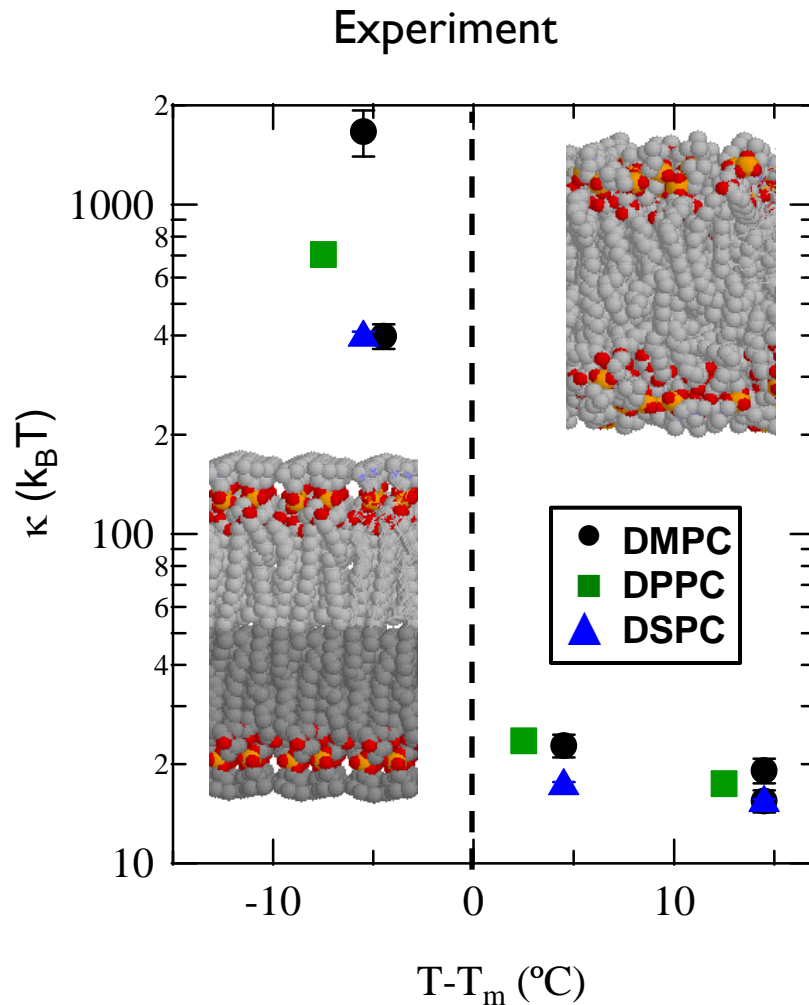
$$\Gamma_{BEND} = 0.025 \left( \frac{k_B T}{\kappa} \right)^{1/2} \frac{k_B T}{3\eta_{D_2O}} q^3 \rightarrow \frac{\Gamma}{q^3} = \frac{\Gamma_{BEND}}{q^3} + \frac{\Gamma_{TF}}{q^3} \frac{1}{1 + (q - q_0)^2 \xi^2}$$

Zilman-Granek

M. Nagao



# NSE : *Bending Modulus*

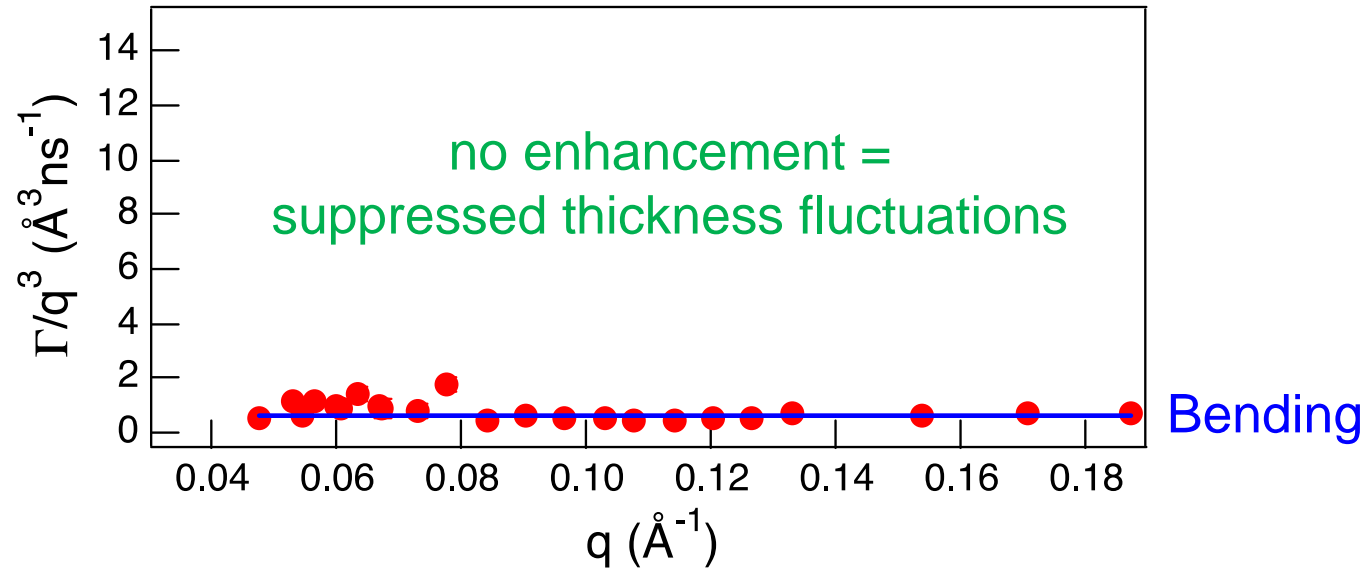


Z. Yi, et. al, J. Phys.: Cond. Mater, **21**, 155104 (2009).

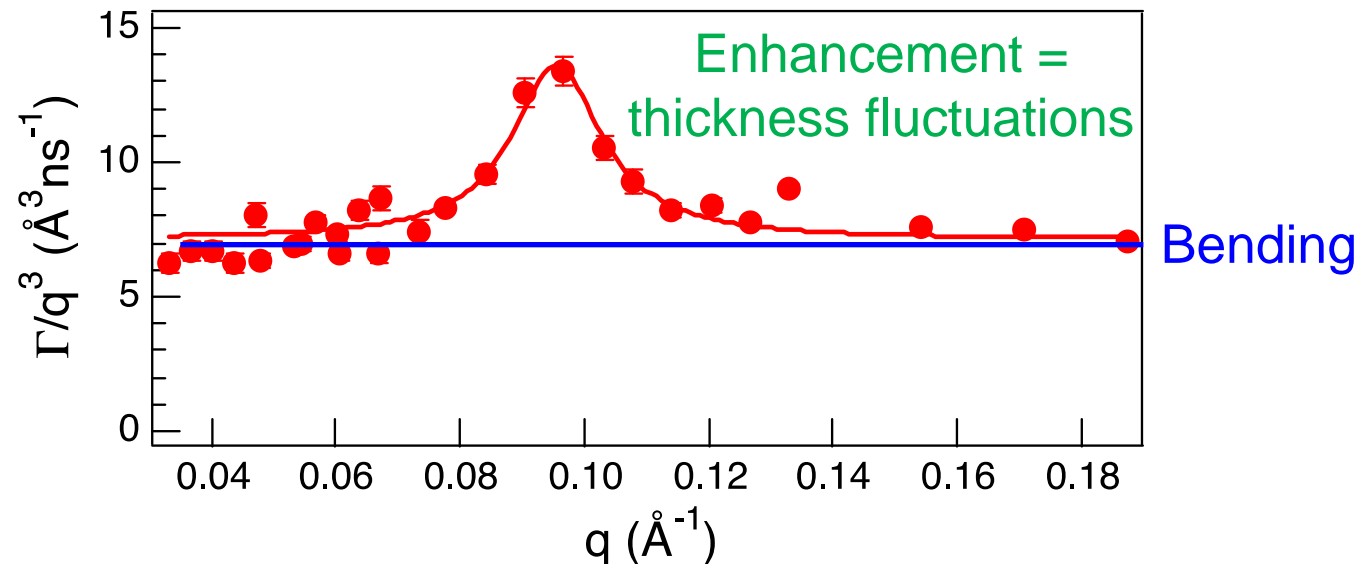
Lipid	$T_m$ ( $^{\circ}C$ )	$T$ ( $^{\circ}C$ )	$(T - T_m)$ ( $^{\circ}C$ )	$\kappa_c / k_B T$
14:0 PC	24	22	-2	$100.0 \pm 4.99$
		24	0	$20.9 \pm 0.61$
		28	+4	$13.9 \pm 0.24$
		35	+11	$15.3 \pm 0.31$
		45	+21	$13.9 \pm 0.44$
16:0 PC	41	30	-11	$8.2 \pm 0.12$
		41	0	$49.6 \pm 2.78$
		60	+19	$36.1 \pm 1.49$
18:0 PC	54	40	-14	$9.5 \pm 0.18$
		60	+6	$79.1 \pm 3.23$
				$13.6 \pm 0.24$

# $q$ dependence of the decay rate: *Below vs. Above $T_m$*

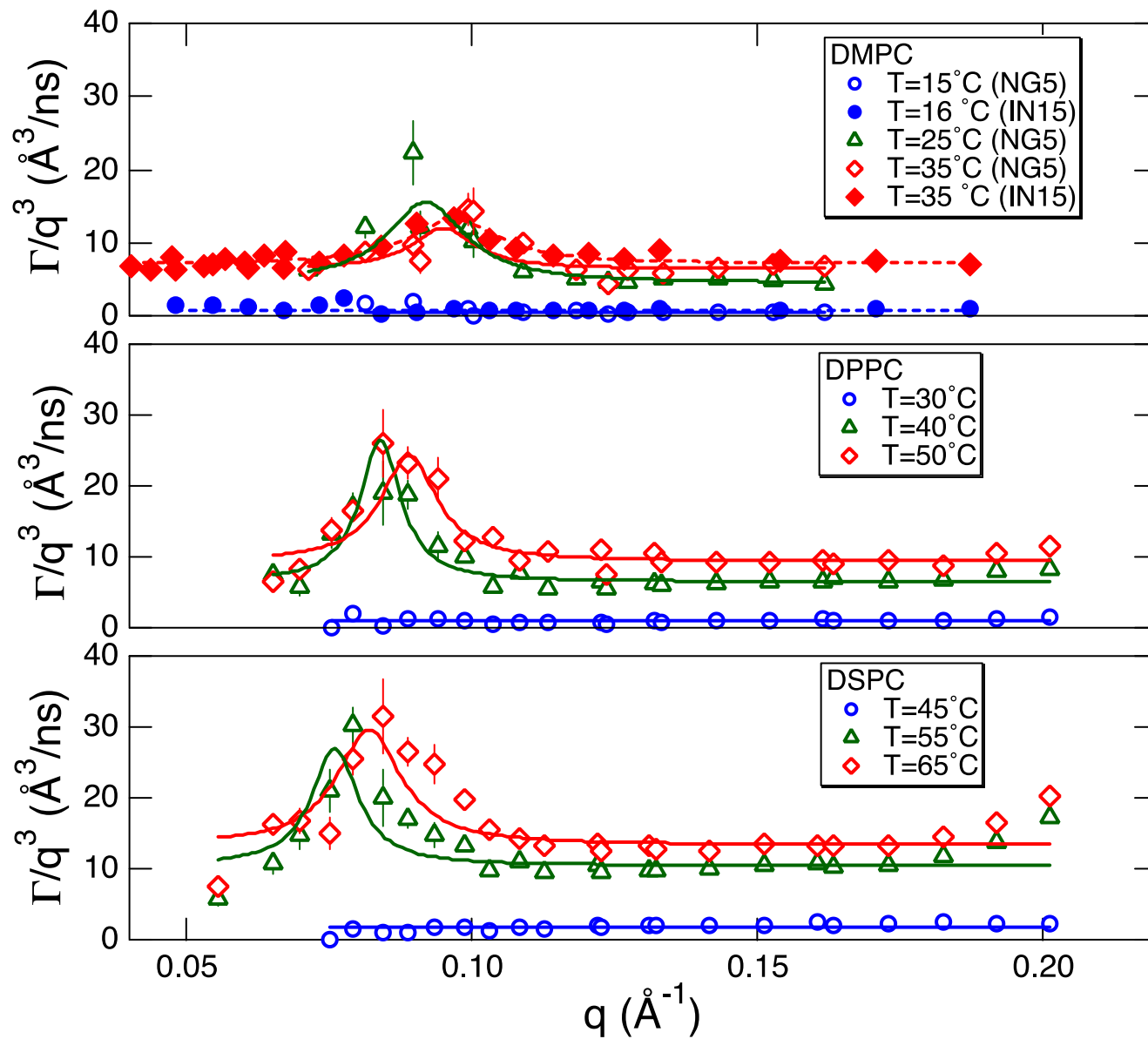
Below  $T_m$   
DMPC  
 $T = 16^\circ\text{C}$



Above  $T_m$   
DMPC  
 $T = 35^\circ\text{C}$

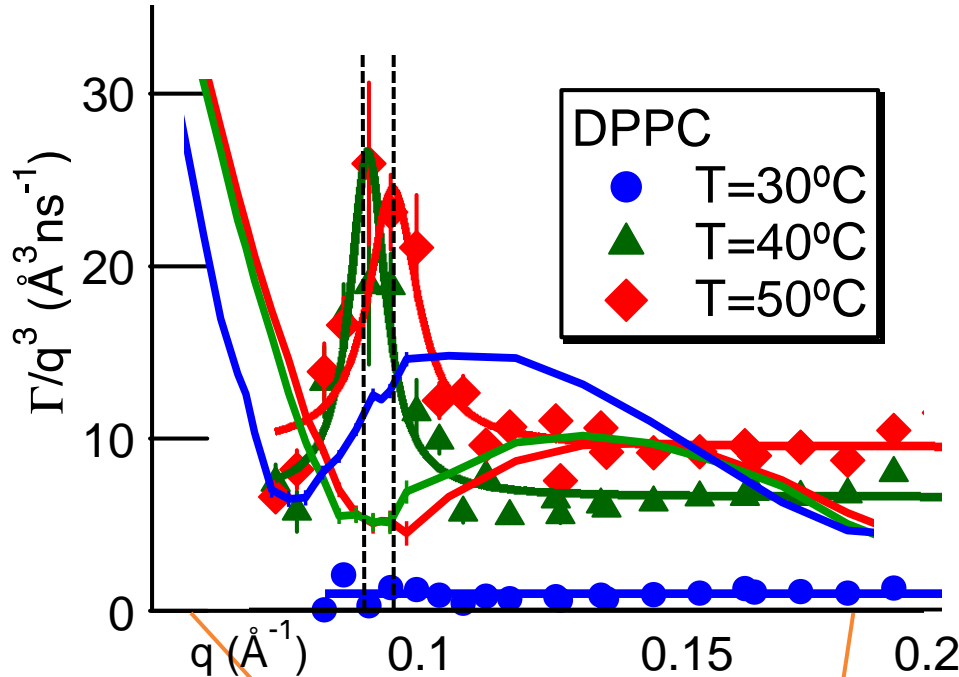


# $q$ dependence of the decay rate: *Lipid Tail Length*



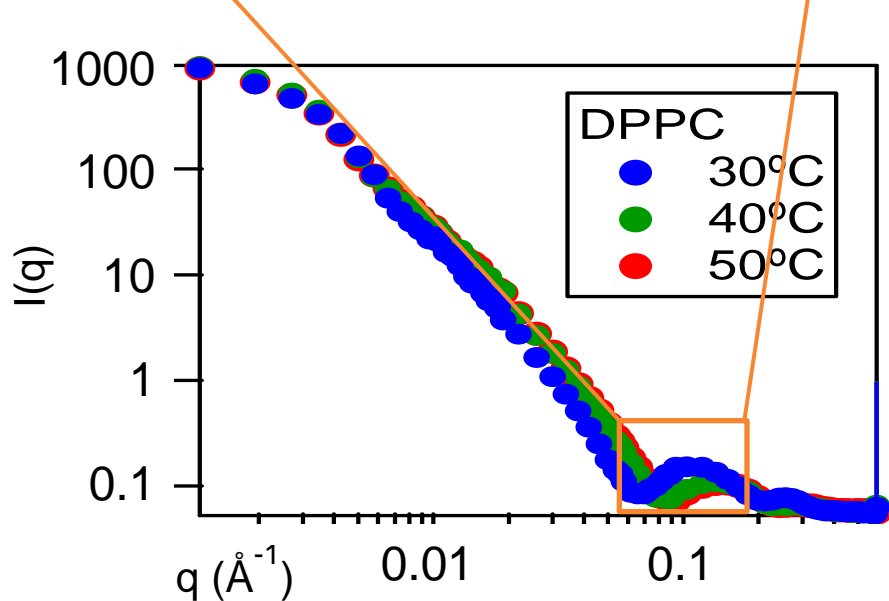


# NSE : Thickness fluctuations ( $\Gamma / q^3$ )



$$\frac{\Gamma}{q^3} = \frac{\Gamma_{BEND}}{q^3} + \frac{\Gamma_{TF}}{q^3} \frac{1}{1 + (q - q_0)^2 \xi^2}$$

M. Nagao



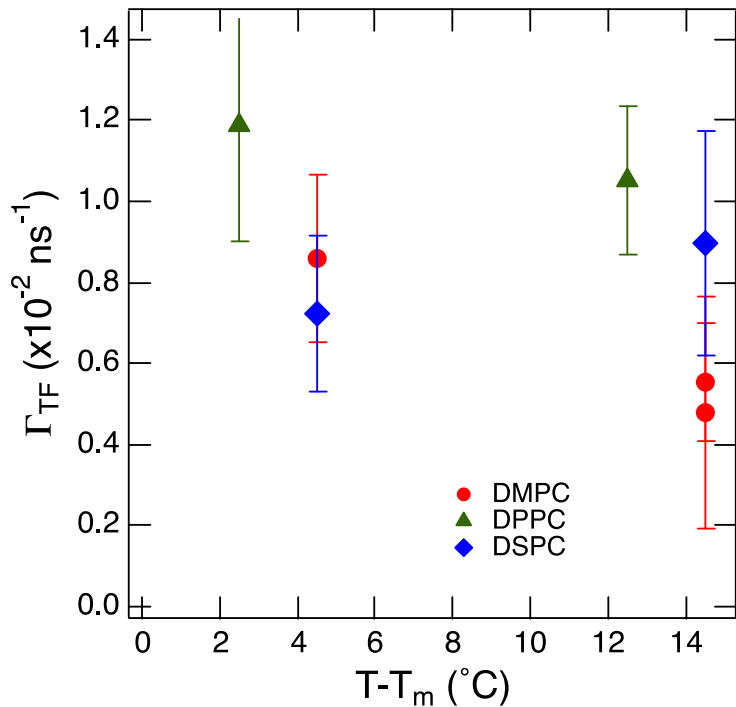
$\Gamma_{BEND} / q^3$ : accounts for bending motions

$\Gamma_{TF}$ : damping frequency of thickness fluctuations

$\xi^{-1}$ : Proportional to the amplitude of thickness fluctuations

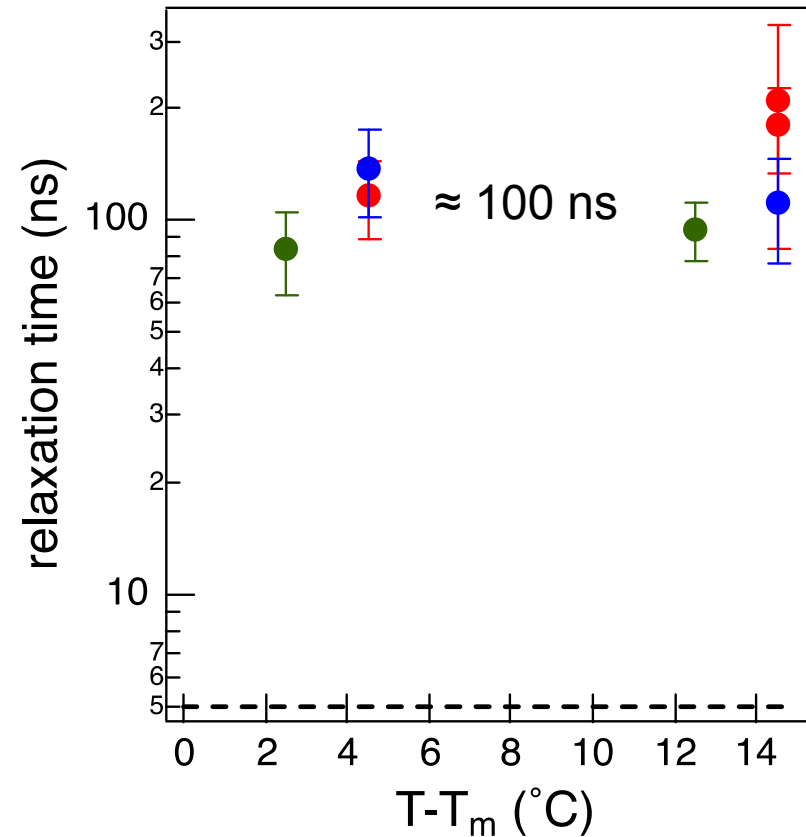
$q_0$ : SANS dip position (Lorentzian peak position)

# Membrane Thickness Fluctuations Time scale



Above  $T_m$ : independent of either temperature or tail length

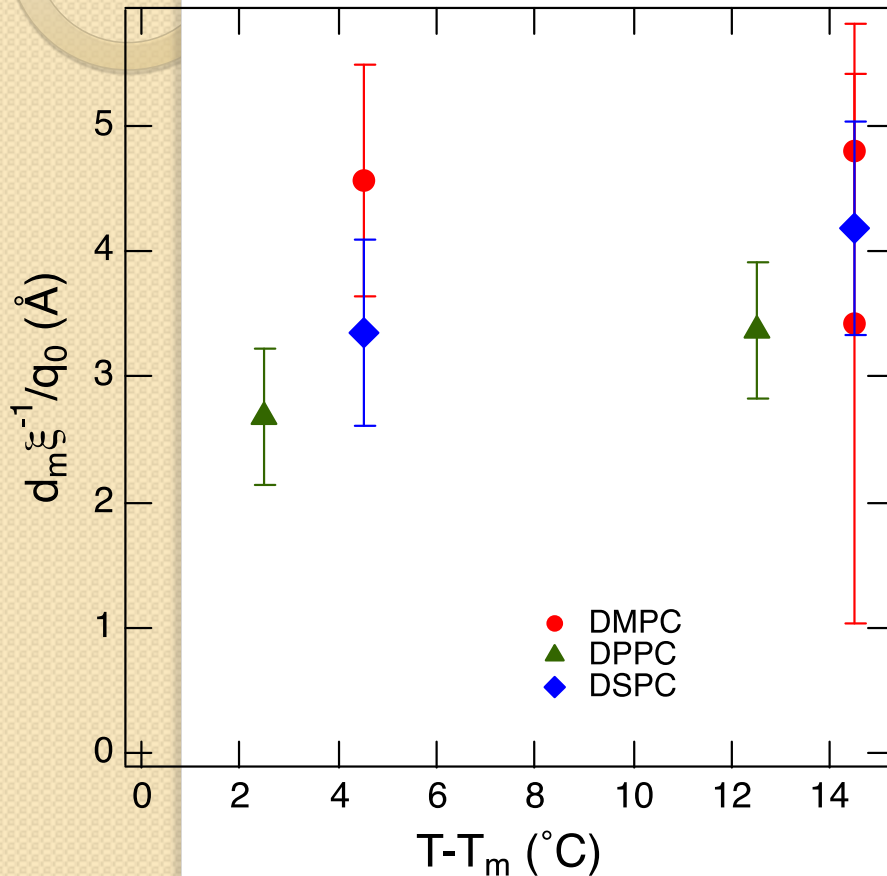
$$\frac{\Gamma}{q^3} = \frac{\Gamma_{BEND}}{q^3} + \frac{\Gamma_{TF}}{q^3} \frac{1}{1 + (q - q_0)^2 \xi^2}$$



An order of magnitude slower (than surfactant membranes)

# Membrane Thickness Fluctuations Amplitude

Width of Lorentzian peak relates to the fluctuation amplitude<sup>(1)</sup>



Experiment:

Mean amplitude =  $3.7 \text{ \AA} \pm 0.7 \text{ \AA}$

Theory:

Huang's mean amplitude  $\approx 4.5 \text{ \AA}$ <sup>(2)</sup>

Simulation:

Lindahl & Edholm's amplitude  $\approx 5 \text{ \AA}$ <sup>(3)</sup>

$\approx 8 \%$  of the membrane thickness;  
close to the value seen in surfactant  
membranes ( $\approx 12 \%$ )

Suggests amplitude is defined by physical constraints, like volume conservation

# Thickness Fluctuation Theory

Although membrane thickness fluctuations have not been previously measured Huang<sup>(7)</sup> has proposed a theory for thickness fluctuations in a lipid bilayer under the consideration of deformation free energy:

$$\langle D^2 \rangle = \frac{k_B T}{2\pi a K_1 C_2} \left\{ \tan^{-1} \left[ \frac{\left( \left( \frac{2\pi}{\lambda_0} \right)^2 + C_1 \right)}{C_2} \right] - \tan^{-1} (C_1 / C_2) \right\}$$

$$C_1 = \frac{\gamma}{2aK_1} \quad C_2 = \left( \frac{\bar{B}}{a^2 K_1} \right)^{1/2}$$

$$K_1 = \frac{\kappa}{d_m}$$

$$d_m = 2a \quad \lambda_0 \approx a$$

$$D \approx 4.5 \text{ \AA}$$

$D$  = thickness fluctuation amplitude

$B^\#$  = membrane

compressibility

$\lambda_0$  = wavelength cut off

$d_m^*$  = membrane thickness

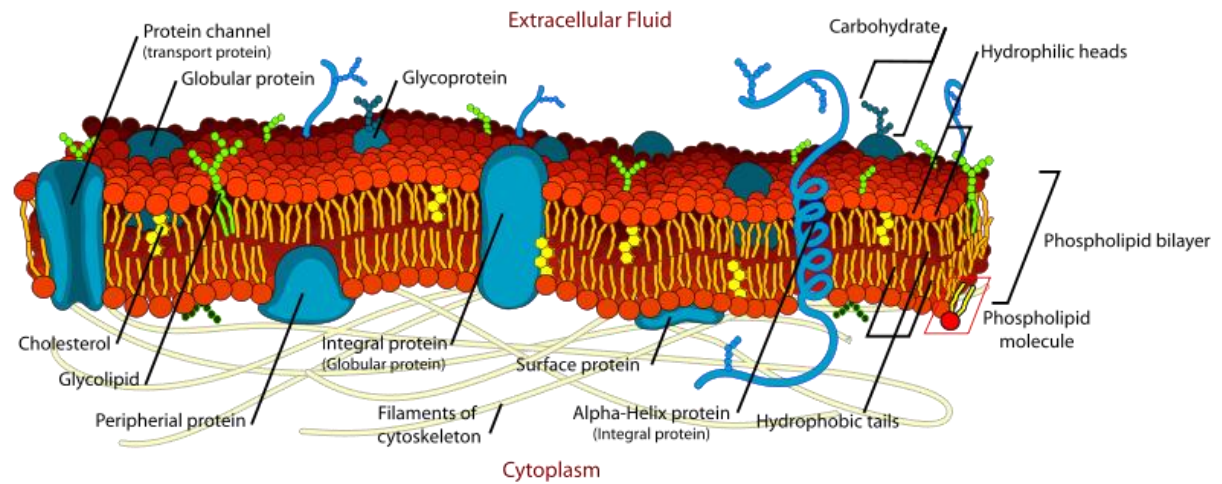
$\kappa^+$  = bending modulus

$\gamma^\#$  = surface tension

+ from NSE measurements    \* from SANS measurements    # from literature

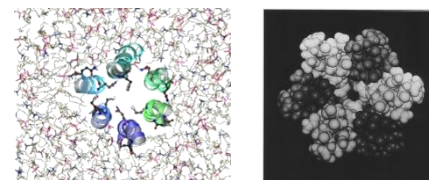
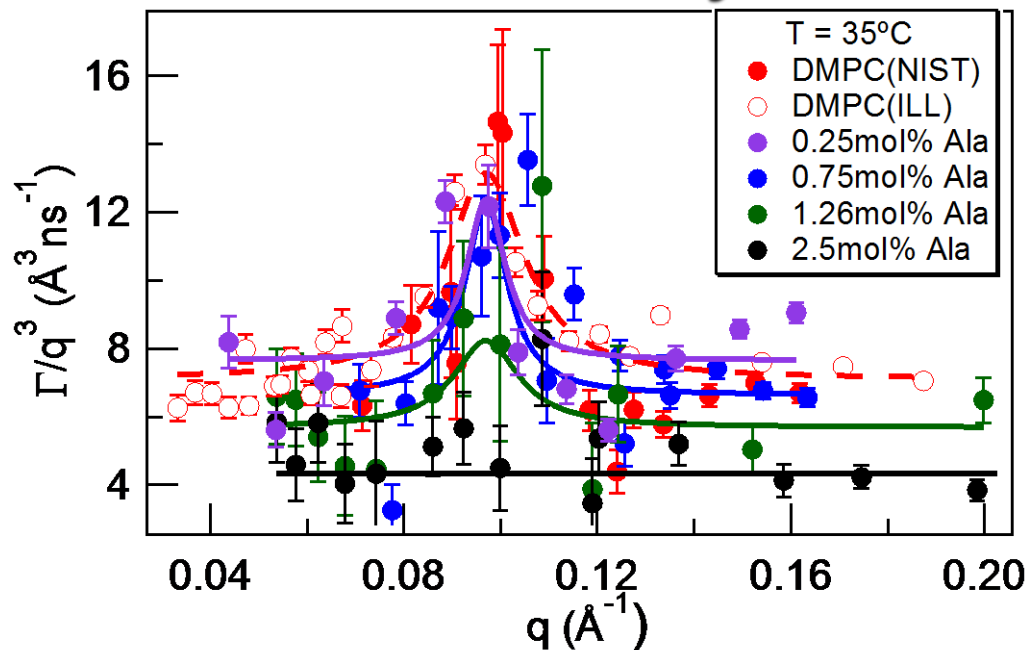
# Conclusions

- NSE was used to successfully measure lipid membrane thickness fluctuations
- From SANS it is clear that these fluctuations appear at the length scale of the membrane thickness.
- The relaxation time  $\approx 100$  ns and is independent of temperature and tail length.
  - An order slower than that observed in surfactant membranes
- Amplitude is  $\approx 8\%$  of the thickness, consistent with surfactant membranes (12%).
  - Volume conservation may define the fluctuation amplitude.
- Below  $T_m$ , thickness fluctuations are not observed, suggesting total suppression of the mode or much slower relaxation times which are not accessible by the current setup.
- The experimental amplitude agrees well with both theory and simulation
- FUTURE DIRECTION: What kind of effects do membrane associated molecules have on membrane dynamics such as thickness fluctuations?

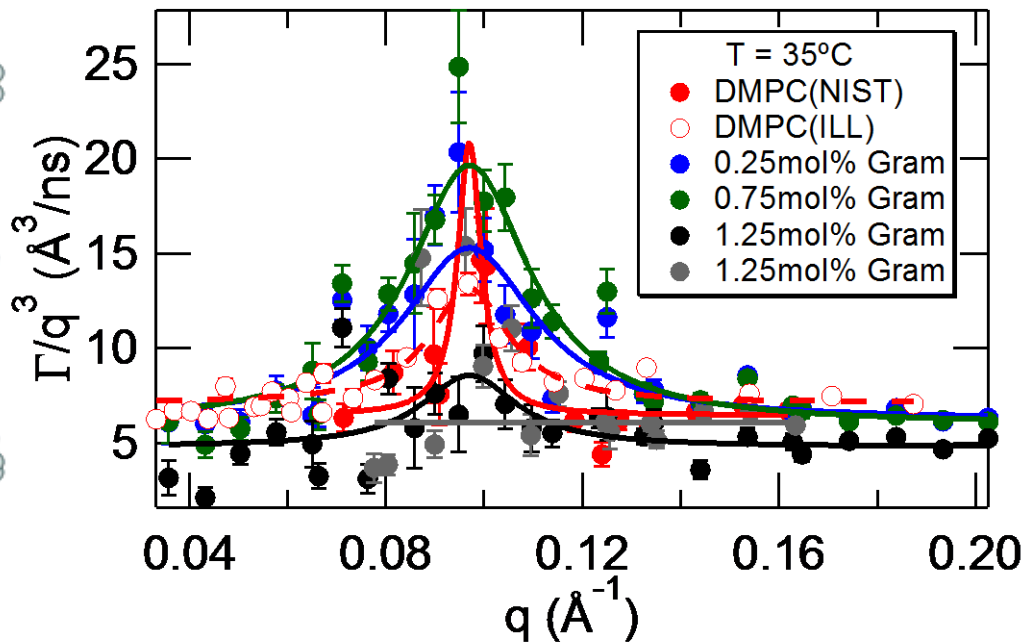
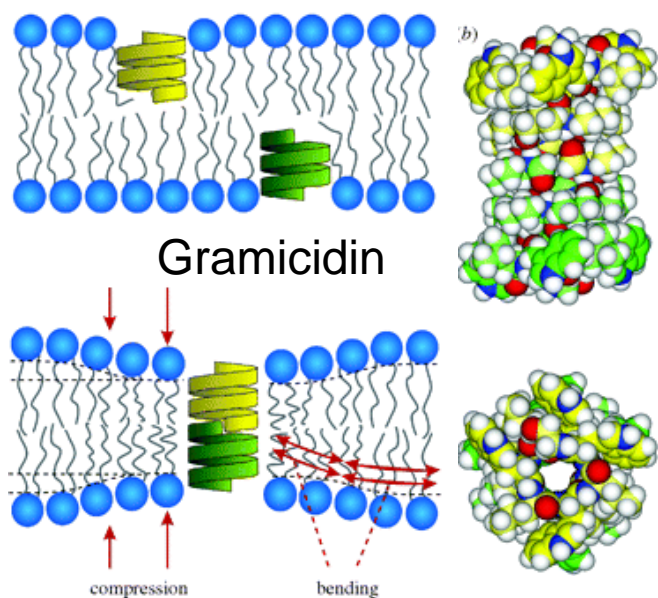
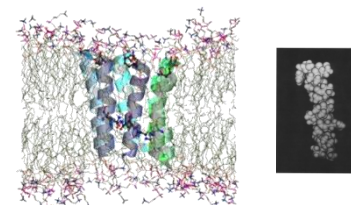




# Preliminary Data w/ Protein



Alamethicin



# Preliminary Data w/ Protein

