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Magnetic phase transition and spin fluctuations in the geometrically frustrated antiferromagnetic spinel CdCr₂O₄: An experiment using the SPINS cold-neutron triple axis spectrometer

<u>Group A</u>

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$CdCr_2O_4$: $Cr^{3+}(3d^3, S=3/2)$

Magnetic Cr^{3+} ions form a lattice of corner-shared tetrahedra Similar lattices are found in spinel (AB₂O₄) B-sites or pyrochlores (A₂B₂O₇)



Magnetic exchange energy between a pair of spins is $H = J_{ij} \mathbf{S}_i \cdot \mathbf{S}_j$. If $J_{ij} > 0$, then \mathbf{S}_i and \mathbf{S}_j will become antiparallel to each other.

But CdCr₂O₄ eventually orders at very low temperatures. How?

frustrated

 $ZnCr_2O_4$

unfrustrated

CdCr₂O₄

Magnetic phase transition in CdCr₂O₄

Magnetic susceptibility $\chi = dM/dH$

M: magnetization of the material H: applied magnetic field



SPINS cold neutron triple axis spectrometer



Why SPINS for this study? Because SPINS - Precisely access desired Q and $\hbar\omega$

- Covers $\hbar\omega$ in the range between 0.1 to 10 meV

- Perform diffraction measurement

- Provides a flexible choice of high resolution or high intensity

Looking for the degenerate ground states in the geometrically frustrate phase

Using simulations, we can try to determine the magnetic structure by comparing to the collected data.



Simple calculation of magnetic neutron scattering intensity Magnetic neutron scattering cross section $\frac{\mathrm{d}^2 \sigma}{\mathrm{d}\Omega \mathrm{d}\omega} = r_o^2 \frac{k_f}{k_i} S(\mathbf{Q}, \omega)$ where, $S(\mathbf{Q}, \omega) = \sum_{\alpha, \beta} (\delta_{\alpha\beta} - \tilde{Q}_{\alpha}\tilde{Q}_{\beta}) \sum_{\lambda \neq \lambda'} p_{\lambda} \sum_{l \neq l} \sum_{l' \neq l'} f_{d'}^{*}(\mathbf{Q}) f_{d'}(\mathbf{Q}) \exp\{i\mathbf{Q} \cdot (\mathbf{R}_{l'd'} - \mathbf{R}_{ld})\}$ $\times \left\langle \lambda \left| \hat{S}_{ld}^{\alpha} \left| \lambda' \right\rangle \right\rangle \left\langle \lambda' \left| \hat{S}_{l'd'}^{\beta} \right| \lambda \right\rangle \delta(\hbar \omega + \hbar \omega_{\lambda} - \hbar \omega_{\lambda'})$ But if we consider only up and down spins for diffuse quasi-elastic scattering, all we need is the following simple equation: $I(\mathbf{Q}) \propto \left| \sum_{\mathbf{R}} f_{\mathbf{R}}(\mathbf{Q}) \sigma_{\mathbf{R}} e^{i\mathbf{Q}\cdot\mathbf{R}} \right|^{2}$ f : magnetic form factor

$$\sigma = -1$$
, or 1

Q-dependence along the (h,0,0) and (0,k,0) crystal configurations

Experimental Data Theoretical Simulation 4 (a) $\hbar \omega = 0.6 \text{ meV}$ T = 15 K 3 3 (0 k 0) (r.l.u.) (0 k 0) (r.l.u.) 0 0 3 3 4 2 4 1 0 1 $(h \ 0 \ 0)$ (r.l.u.) (h 0 0) (r.l.u) T = 15 K

Q-dependence along the (h,h,0) and (0,0,l) crystal configurations

Experimental Data

Theoretical Simulation



Temperature Dependence



Summary

- Using the Triple-Axis Spectrometer, we studied the magnetic phase transitions in CdCr₂O₄, which is a spinel with antiferromagnetic interactions between Cr³⁺ ions on B sites.
- We examined the Q-dependence on different crystal planes.
- Using a theoretical simulations, we were able to determine the basic magnetic structure.
- Through the use of temperature dependence, we were able to find that the discovered peaks are consistent with magnetic excitations.
- The examination of the line width for the quasi-elastic fluctuation shows the lifetime to be on the pico-second scale.

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