

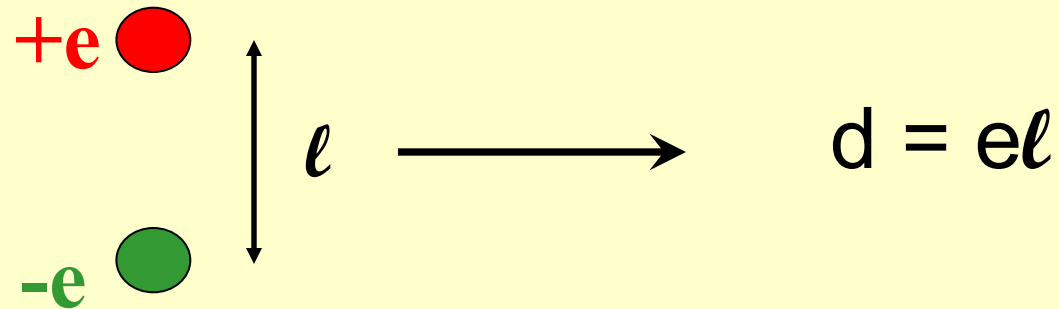
# Neutron Electric Dipole Moment (EDM)

- Why is it interesting? (recall S. Gardner)
- How do we measure it?
- What is the present limit?
- How can we significantly improve the sensitivity (& discover neutron EDM!!)?

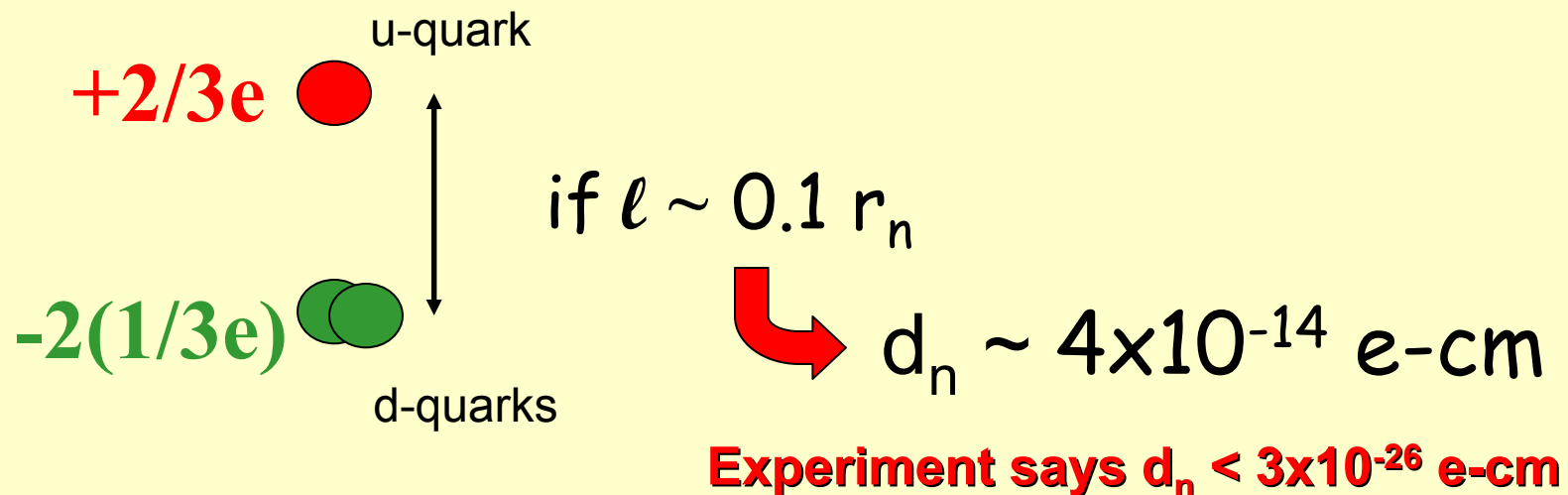


**Brad Filippone**  
**NIST Summer School**  
**Fundamental Neutron Physics**  
**June 26, 2009**

# What is an EDM?

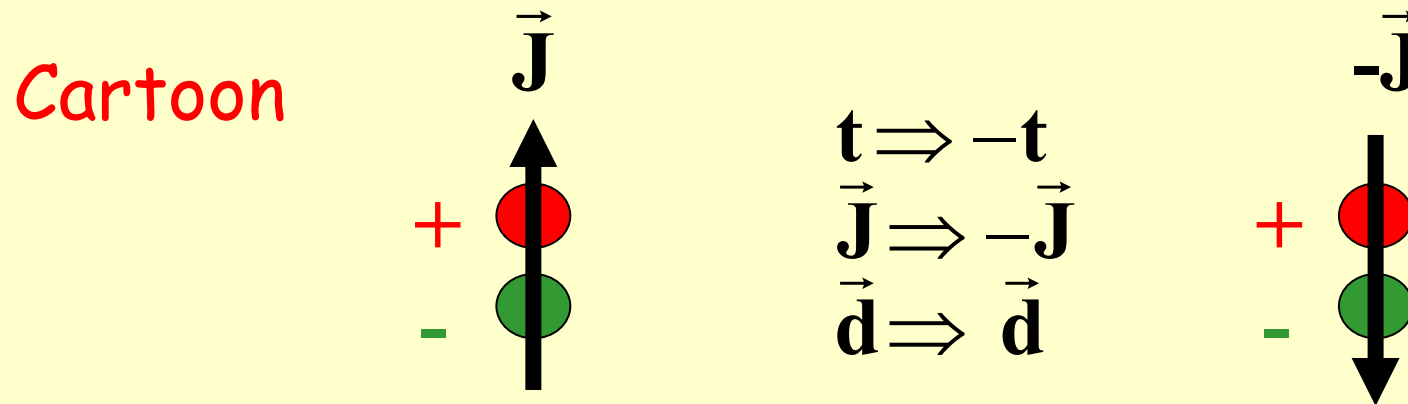


## How big is the neutron EDM?



# Why Look for EDMs?

- Existence of EDM implies violation of Time Reversal Invariance



- Time Reversal Violation seen in  $K^0-\bar{K}^0$  system
- May also be seen in early Universe
  - Matter-Antimatter asymmetry

but the Standard Model effect is too small !

# Quantum Picture - Discrete Symmetries

(08 Nobel Prize)

Charge Conjugation :  $\hat{C} \bullet \psi_n \Rightarrow \psi_{\bar{n}}$

Parity :  $\hat{P} \bullet \psi(x, y, z) \Rightarrow \psi(-x, -y, -z)$

Time Reversal :  $\hat{T} \bullet \psi(t) \Rightarrow \psi(-t)$

Assume  $\vec{\mu} = \mu \frac{\vec{J}}{J}$  and  $\vec{d} = d \frac{\vec{J}}{J}$

Non-Relativistic Hamiltonian

$$H = \underbrace{\vec{\mu} \cdot \vec{B}}_{\text{C-even}} + \underbrace{\vec{d} \cdot \vec{E}}_{\text{C-even}}$$

C-even	C-even
P-even	P-odd
T-even	T-odd

**Non-zero d violates T and CP**  
(Field Theories generally preserve CPT)

	C	P	T
$\uparrow \mu \uparrow B$	-	+	-
$\uparrow d \uparrow E$	-	+	-
$\uparrow \mu \uparrow E$	-	-	+
$\uparrow d \uparrow B$	-	+	-
$\uparrow J \uparrow$	+	+	-

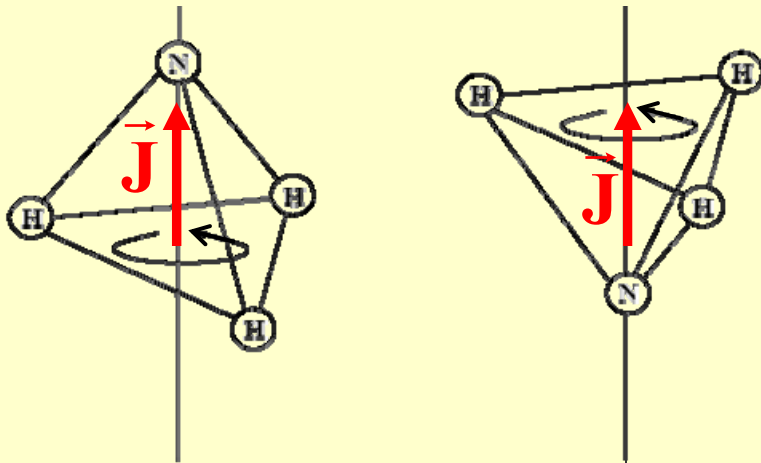
# But some molecules have HUGE EDMs!

H<sub>2</sub>O:  $d = 0.4 \times 10^{-8} \text{ e-cm}$

NaCl:  $d = 1.8 \times 10^{-8} \text{ e-cm}$

NH<sub>3</sub>:  $d = 0.3 \times 10^{-8} \text{ e-cm}$

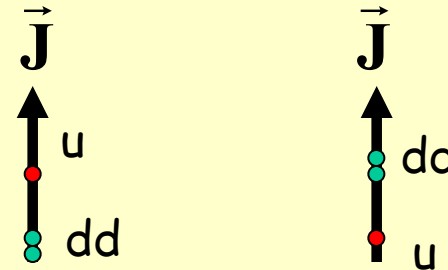
Note: n-EDM  $< 3 \times 10^{-26} \text{ e-cm}$



But NH<sub>3</sub> EDM is not T-odd or CP-odd since

$$\vec{d} \neq d \frac{\vec{J}}{J}$$

If Neutron had degenerate state



it would not violate T or CP

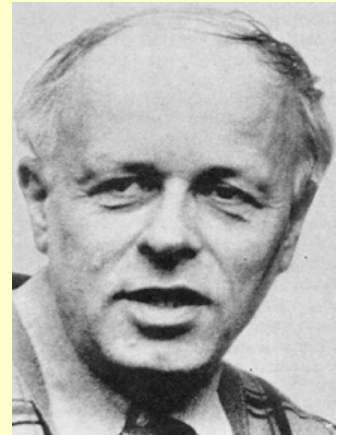
**both  $\vec{d} = +d \frac{\vec{J}}{J}$  and  $\vec{d} = -d \frac{\vec{J}}{J}$  exist!** Ground state is actually a superposition

# Role of CP Violation in the Matter/Antimatter Asymmetry of the Universe

- **Sakharov Criteria**

- Particle Physics can produce matter/antimatter asymmetry in the early universe *IF* there is:

- Baryon Number Violation
- CP & C violation
- Departure from Thermal Equilibrium

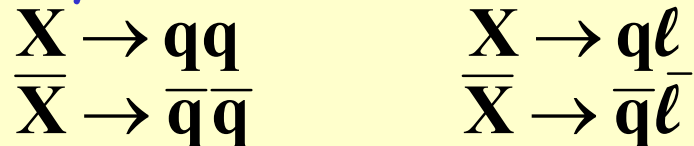


# Baryogenesis

- **Plausibility Argument**

- Consider heavy boson -  $X$

- Baryon number violation:



- C-Violation & CP-Violation

$$\begin{array}{ll} \Gamma_{X \rightarrow qq} = (1 + \Delta_q)\Gamma_q; & \Gamma_{X \rightarrow q\ell_L} = (1 - \Delta_\ell)\Gamma_\ell \\ \Gamma_{\bar{X} \rightarrow \bar{q}\bar{q}} = (1 - \Delta_q)\Gamma_q; & \Gamma_{\bar{X} \rightarrow \bar{q}\bar{\ell}_R} = (1 + \Delta_\ell)\Gamma_\ell \end{array}$$

**but**  $\Gamma_X^{\text{Tot}} = \Gamma_{\bar{X}}^{\text{Tot}}$  (CPT conservation!!) **if**  $\Delta_q\Gamma_q = \Delta_\ell\Gamma_\ell$

- Out of Thermal Equilibrium

Otherwise, in Equilibrium the reverse reactions:  
(e.g.  $qq \rightarrow X$ ,  $\bar{q}\bar{q} \rightarrow \bar{X}$ ) will smooth out any matter/antimatter excess

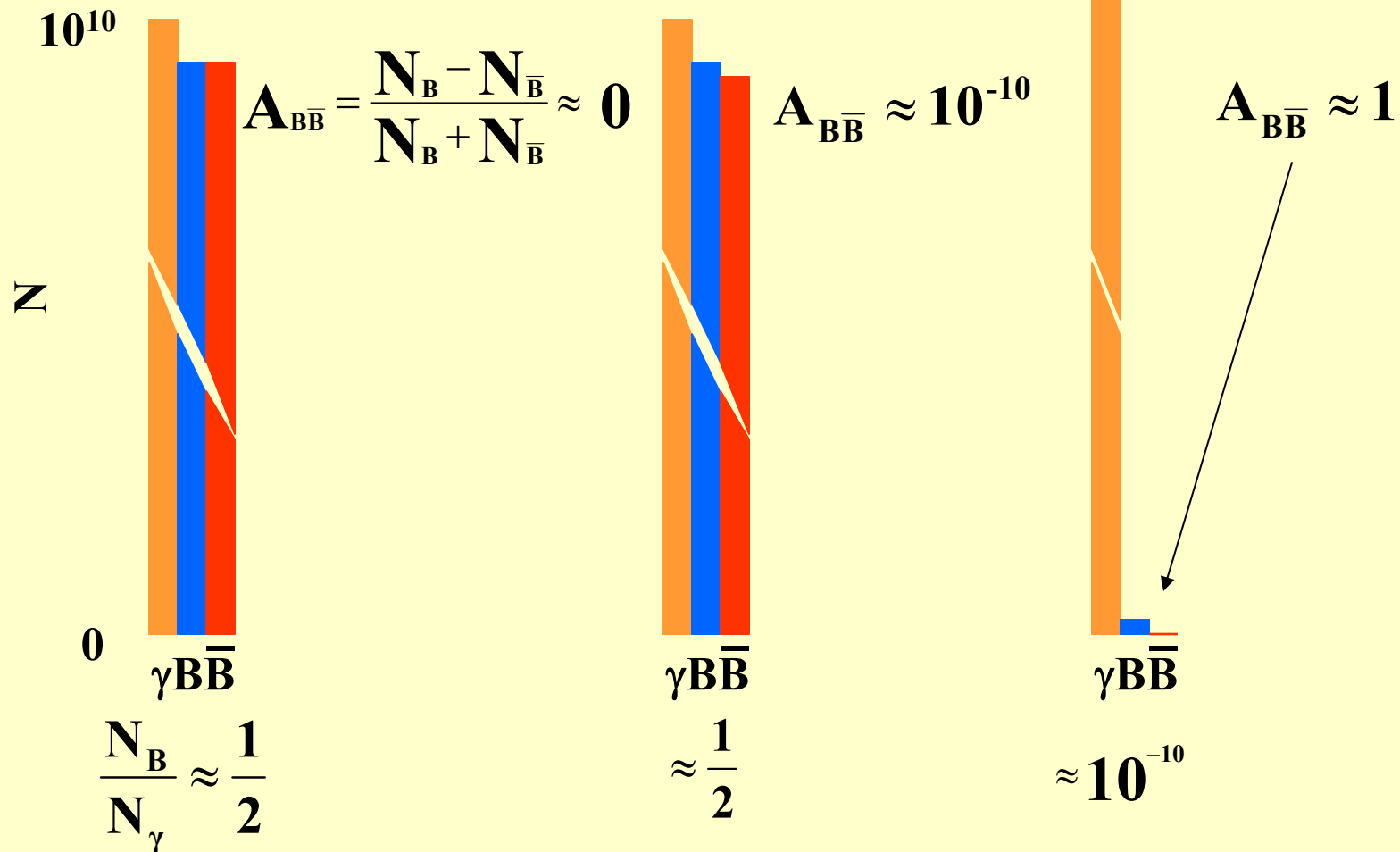
# Electroweak Baryogenesis

## Possible source of Matter-Antimatter Asymmetry

Before Electroweak Phase Transition

After EW Phase Transition

Today





# But Standard Model CP violation (CKM matrix) is Insufficient

- Must search for new sources of CP
  - B-factories, Neutrinos, EDMs
- Quarks/Gluons
  - Allows production of matter-antimatter asymmetry via "Baryogenesis"
- Neutrino mixing suggests possibility of new CP violation in leptons
  - Allows production of matter anti-matter asymmetry via "Leptogenesis"



# What's in SUSY?

- Great Names:

- Squarks, sleptons, gauginos, winos, binos, neutralinos,...

- In MSSM

- 124 parameters - 19 from Standard Model & 105 new parameters (from SUSY and also from SUSY breaking)
  - 36 mixing angles for squarks & sleptons
  - 40 CP-violating phases for squarks & sleptons
  - 21 squark & slepton masses
  - 5 couplings and 3 phases from gauginos/higgsinos

# SUSY, CP-Violation and EDMs

- New physics (e.g. SuperSymmetry = SUSY) has additional CP violating phases in added couplings
  - New phases: ( $\phi_{CP}$ ) should be  $\sim 1$  (why not?)
- Contribution to EDMs depends on masses of new particles

$$d_n \sim 10^{-24} \text{ e-cm} \times \sin\phi_{CP} (200 \text{ GeV}/M_{\text{SUSY}})^2$$

Note: experimental limit:  $d_n < 0.03 \times 10^{-24} \text{ e-cm}$   
Standard Model Prediction:  $d_n < 10^{-31} \text{ e-cm}$

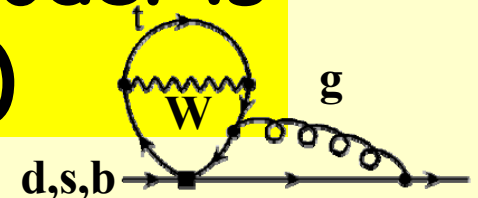
# Origin of EDMs

- **Standard Model EDMs are due to CP violation in the quark weak mixing matrix CKM (e.g. the  $K^0/B^0$ -system) but...**
  - $e^-$  and quark EDM's are zero in 1<sup>st</sup> & 2<sup>nd</sup> order
  - Need at least three Feynman diagram "loops" to get EDM's (electron actually requires 4 loops!)
    - Thus EDM's are VERY small in standard model

**Neutron EDM in Standard Model is**

$$\sim 10^{-32} \text{ e-cm } (=10^{-19} \text{ e-fm})$$

**Experimental neutron limit:  $< 3 \times 10^{-26} \text{ e-cm}$**



**Electron EDM in Standard Model is**

$$< 10^{-40} \text{ e-cm}$$

# Origin of Hadronic EDMs

- Hadronic (strongly interacting particles) EDMs are from
  - $\theta_{\text{QCD}}$  (a special parameter in Quantum Chromodynamics - QCD)
  - or from the quarks themselves

$$\begin{aligned}
 \mathcal{L}_{eff}^{\text{CP}} = & \frac{g_s^2}{32\pi^2} \bar{\theta} G_{\mu\nu}^a \tilde{G}^{\mu\nu,a} + \frac{1}{3} w f^{abc} G_{\mu\nu}^a \tilde{G}^{\nu\beta,b} G_{\beta\mu,c} \\
 & - \frac{i}{2} \sum_{i=e,u,d,s} d_i \bar{\psi}_i (F \cdot \sigma) \gamma_5 \psi - \frac{i}{2} \sum_{i=u,d,s} \tilde{d}_i \bar{\psi}_i g_s (G \cdot \sigma) \gamma_5 \psi + \dots
 \end{aligned}$$

$\theta_{\text{QCD}}$  (Weinberg 3-gluon term)

$e^-$ , quark EDM (quark color EDM (chromo-EDM))

# EDM from $\theta_{\text{QCD}}$

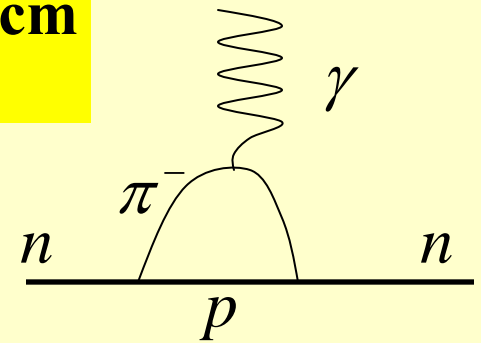
- This is the strong-CP problem in QCD

$$\mathcal{L}_{\text{QCD}} = -\theta \left( \frac{\alpha_s}{8\pi} \right) \tilde{G}_a^{\mu\nu} G_{\mu\nu}^a$$

- $\theta_{\text{QCD}}$  should be naturally about  $\sim 1$
- This gives a neutron EDM of

$$d_n = \frac{g_{\pi NN}}{4\pi^2} \left( \frac{e}{m_p f_\pi} \right) \ln \left( \frac{m_\rho}{m_\pi} \right) \left( \frac{m_u m_d}{m_u + m_d} \right) \theta \approx (-10^{-15}) \theta \text{ e-cm}$$

but  $d_n^{\text{exp}} < 10^{-25} \text{ e-cm}$   
 $\therefore \theta < 10^{-10}$   
 Why so small??



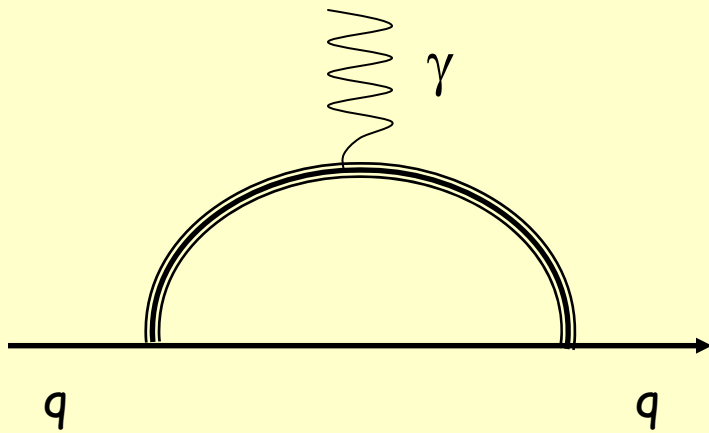
# EDM from $\theta_{\text{QCD}}$

- Small  $\theta_{\text{QCD}}$  does not provide any new symmetry for  $\mathcal{L}_{\text{QCD}}$ 
  - Popular solution is "axions" (Peccei-Quinn symmetry) - new term in  $\mathcal{L}_{\text{QCD}}$ 
    - No Axions observed yet
  - Extra dimensions might suppress  $\theta_{\text{QCD}}$
  - Remains an unsolved theoretical "problem"

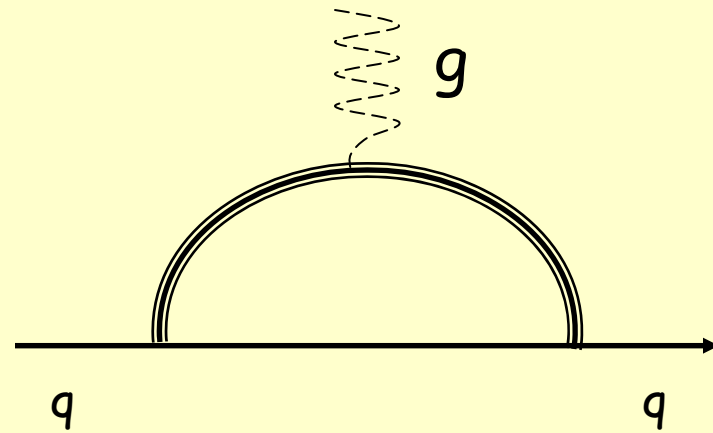


# Hadronic EDM from Quarks

- Quark EDM contributes via



$d_q$   
Quark EDM



$\tilde{d}_q$   
Quark  
ChromoEDM

# Relative EDM Sensitivities

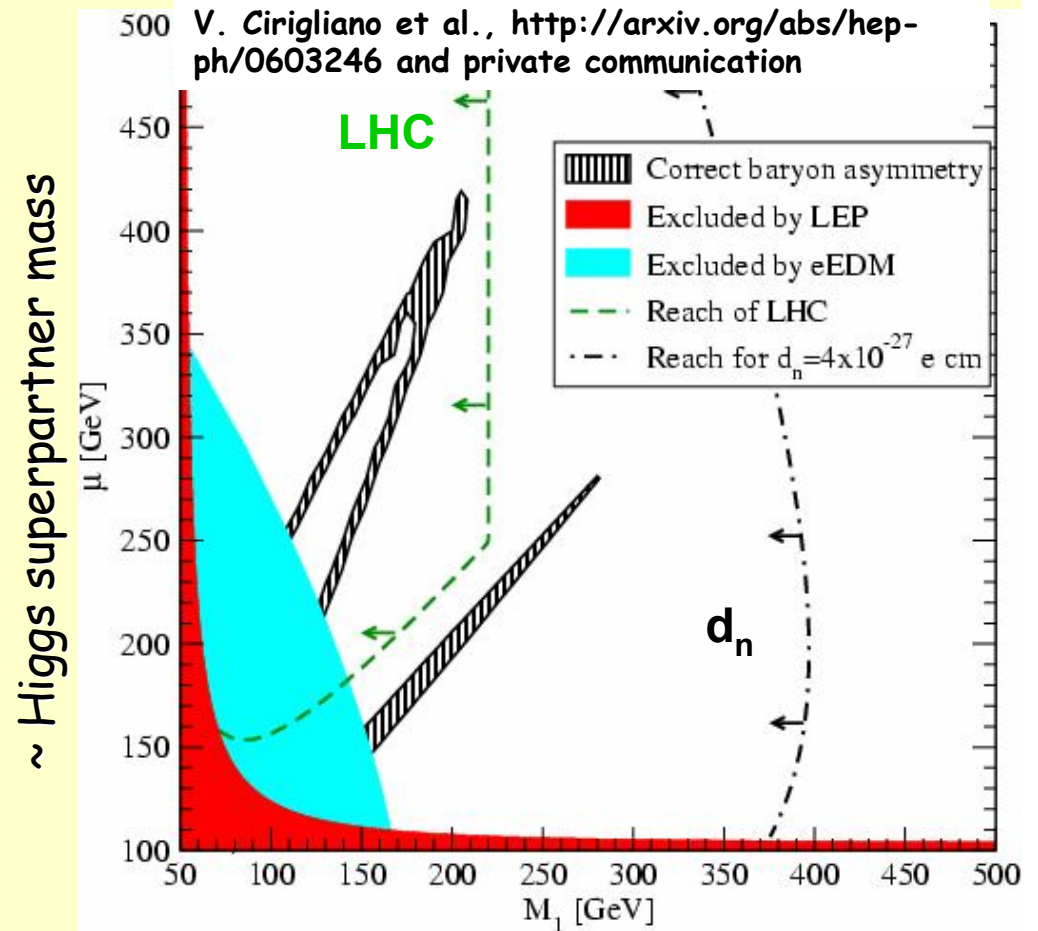
System	Dependence	Present Limit (e-cm)	Future (e-cm)
n	$d_n \sim (3 \times 10^{-16}) \theta_{\text{QCD}} +$ $0.7(d_d - \frac{1}{4} d_u) + 0.6e(\tilde{d}_d + \frac{1}{2} \tilde{d}_u)$	$< 3 \times 10^{-26}$	$10^{-28}$
d	$d_d \sim (-1 \times 10^{-16}) \theta_{\text{QCD}} +$ $6e(\tilde{d}_d - \tilde{d}_u)$	?	$10^{-27}(?)$
$^{199}\text{Hg}$	$d_{\text{Hg}} \sim (0.007 \times 10^{-16}) \theta_{\text{QCD}} -$ $0.007e(\tilde{d}_d - \tilde{d}_u)$	$< 7 \times 10^{-29}$	$10^{-29}(?)$

# Possible impacts of non-zero EDM

- Must be new Physics
- Sharply constrains models beyond the Standard Model (especially *with* LHC data)

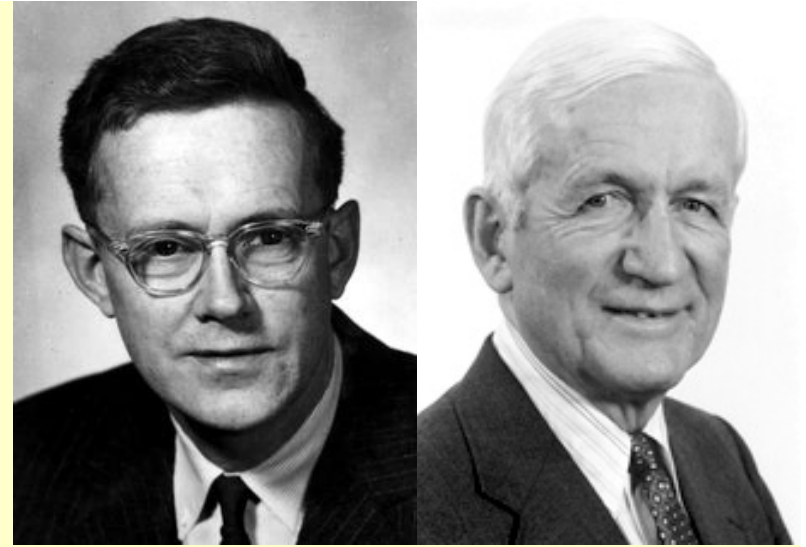
Large Hadron Collider

- May account for matter-antimatter asymmetry of the universe



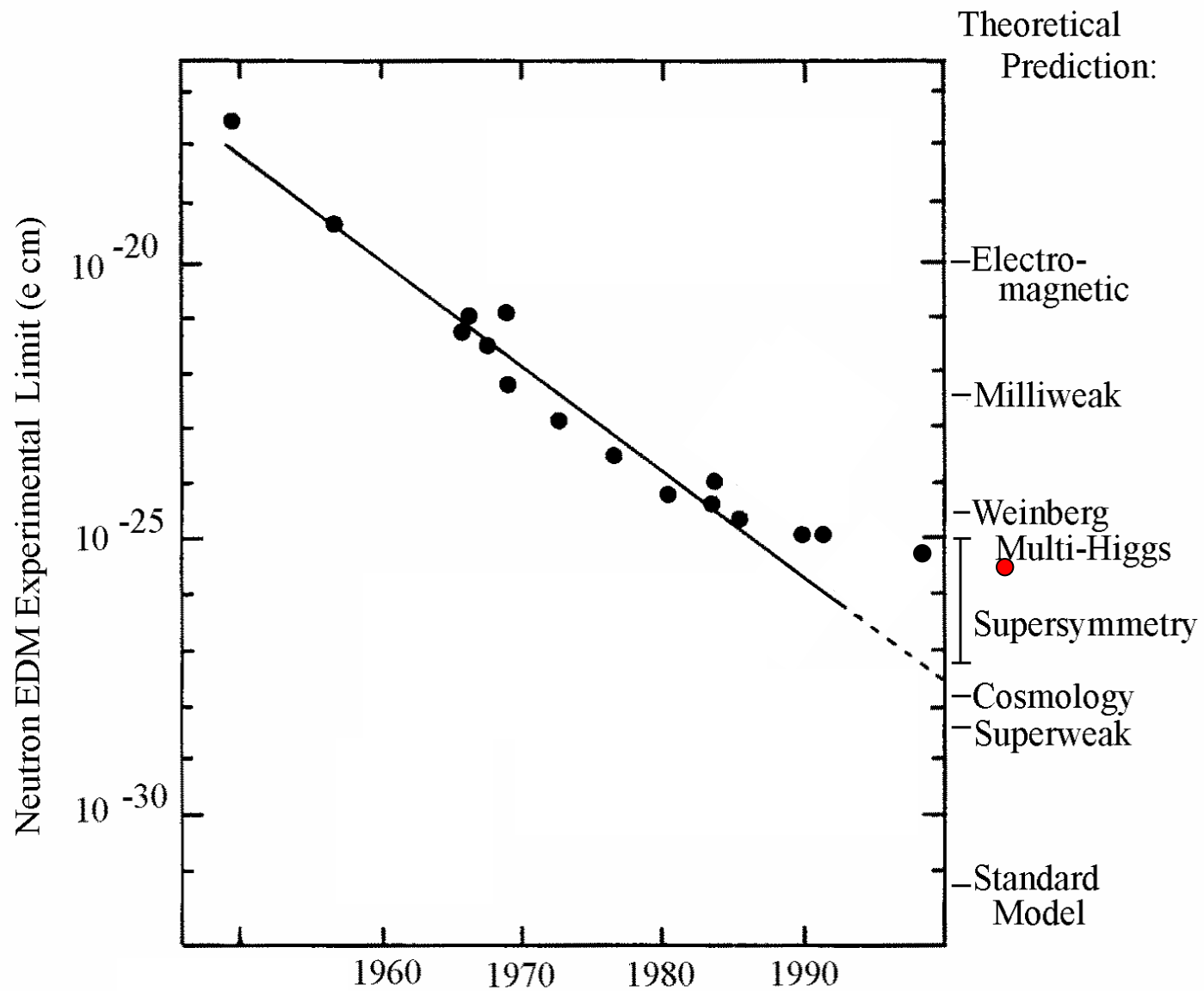
gauge boson superpartner mass

# First result for neutron EDM



- E.M. Purcell and N.F. Ramsey, *Phys. Rev.* **78**, 807 (1950)
  - Neutron Scattering
  - Searching for Parity Violation
  - Pioneered Neutron Beam Magnetic Resonance

# n-EDM vs Time



# How to measure an EDM?

Recall magnetic moment in B field:

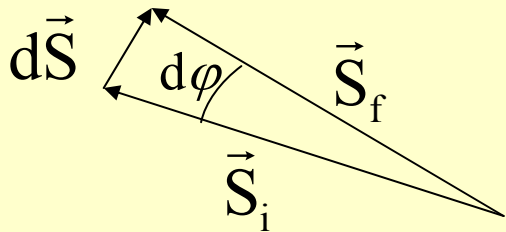
$$\hat{H} = \vec{\mu} \cdot \vec{B}; \quad \vec{\mu} = 2 \left( \frac{\mu_N}{\hbar} \right) \vec{S} \quad ; \text{ for spin } \frac{1}{2}$$

$$\vec{\tau} = \frac{d\vec{S}}{dt} = \vec{\mu} \times \vec{B} \Rightarrow 2 \left( \frac{\mu_N}{\hbar} \right) \|\vec{S}\| \|\vec{B}\|; \quad \text{if } \vec{S} \perp \vec{B}$$

Classical Picture:

- If the spin is not aligned with B there will be a precession due to the torque
- Precession frequency  $\omega$  given by

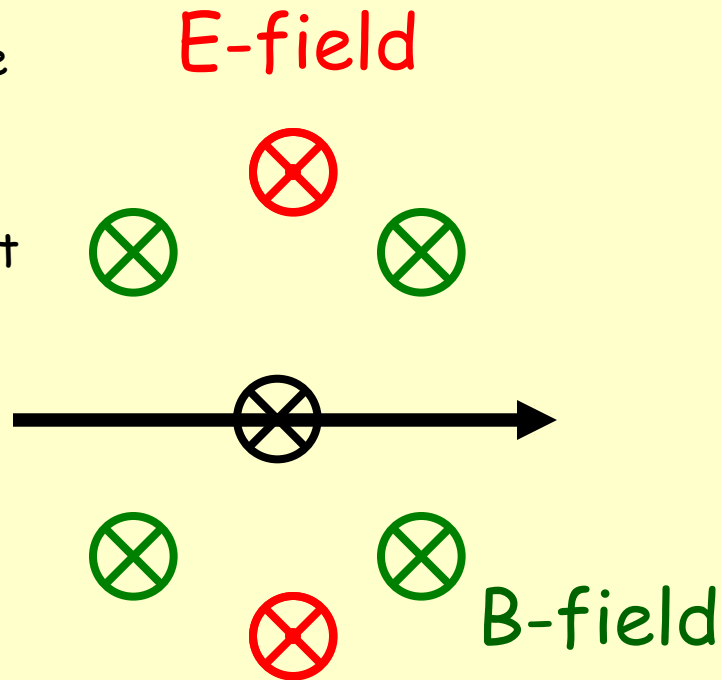
$$\omega = \frac{d\phi}{dt} = \frac{1}{S} \frac{dS}{dt}$$



$$= \frac{2\mu_N B}{\hbar} \quad ; \quad \text{or} \quad \frac{2d_N E}{\hbar} \quad \text{for a } \vec{d}_N \text{ in } \vec{E}$$

# Simplified Measurement of EDM

1. Inject polarized particle
2. Rotate spin by  $\pi/2$
3. Flip E-field direction
4. Measure frequency shift



$$\nu = \frac{2\vec{\mu} \cdot \vec{B} \pm 2\vec{d} \cdot \vec{E}}{h}$$

Must know B very well

# What is the precision in EDM measurement?

$$\mathcal{E} = \hbar\omega = \vec{\mathbf{d}} \cdot \vec{\mathbf{E}}$$

Using Uncertainty Principle:

$$\Delta\mathcal{E}\Delta t \sim \hbar$$

Precise energy measurement requires long measurement time, giving

$$\sigma_d \sim \frac{\hbar}{|\vec{\mathbf{E}}| T_m}$$

But must include counting statistics  $\propto \frac{1}{\sqrt{N}}$

**Sensitivity:** 
$$\sigma_d \cong \frac{\hbar}{|\vec{\mathbf{E}}| T_m \sqrt{mN}}$$

**E** – Electric Field

**T<sub>m</sub>** – Time for measurement

**m** – total # of measurements

**N** – Total # of counts/meas.

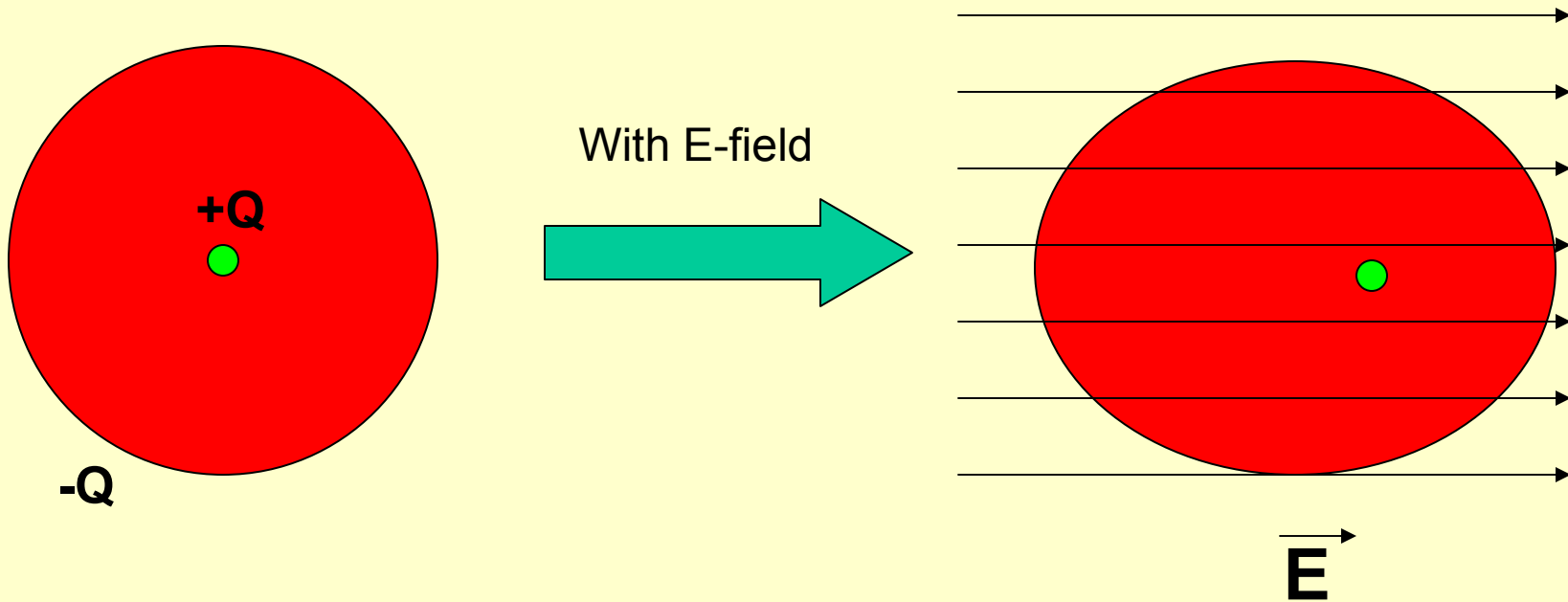


# What particles can be measured?

- Charged particle is difficult
  - Electric field accelerates
  - May work for storage ring
- Neutral particle is easier
  - Atoms (for electron EDM)
    - Also can work for quark EDM
  - Free Neutrons (for quark EDM)

# Atomic EDMs

- Schiff Theorem
  - Neutral atomic system of point particles in Electric field readjusts itself to give zero E field at all charges



# But ...

- Magnetic effects and finite size of nucleus can break the symmetry (relativistic effects can also enhance)
  - Enhancement for  $d_e$  in paramagnetic atoms (**unpaired electrons**)  
(magnetic effect with mixing of opposite parity atomic states)

$$\text{Thus } d_{\text{Tl}} \sim -585 d_e \quad \& \quad |d_e| < 1.5 \times 10^{-27} \text{ e-cm}$$

- Suppression for hadronic EDMs in Diamagnetic atoms (**paired electrons**) - but Schiff Moment is non-zero  
(due to finite size of nucleus and nuclear force)

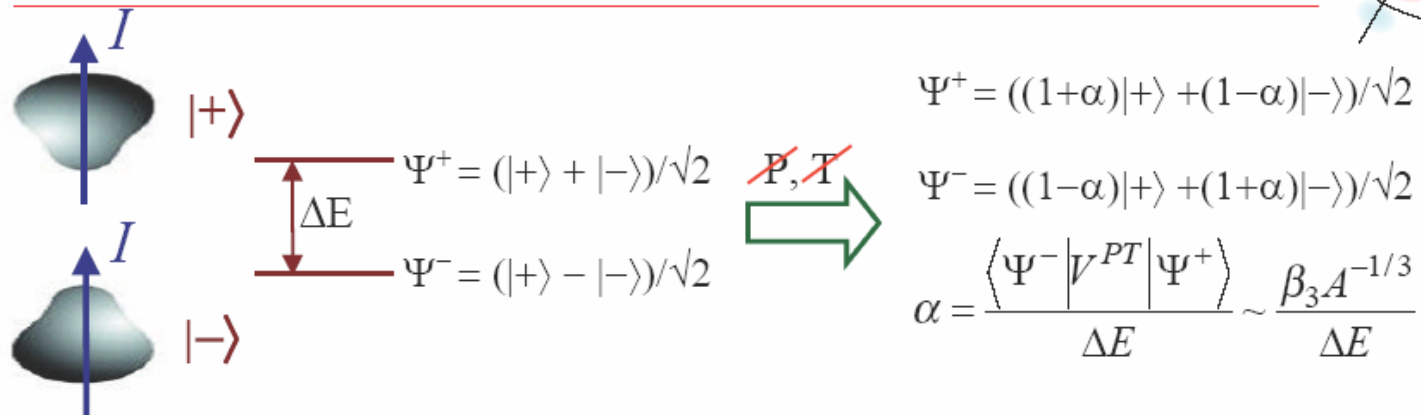
$$\text{Naively expect } d_A \sim \left( \frac{R_{\text{Nucleus}}}{r_{\text{Atom}}} \right)^2 d_{n,p} \sim \left( \frac{A^{1/3} R_0}{a/Z} \right)^2 d_{n,p} \sim 10^{-4} d_{n,p}$$

**for  $^{199}\text{Hg}$**

# But, but, ...

## Can enhance heavy atom EDMs via nuclear deformation

### Octupole deformations



$$S_{\text{intr}} \sim eZA\beta_2\beta_3$$

$$S_{\text{lab}} \sim eZA^{2/3}\beta_2\beta_3^2/\Delta E$$

$$\beta_2, \beta_3 \sim 0.1$$

Haxton & Henley; Auerbach, Flambaum & Spevak; Hayes, Friar & Engel

	$^{223}\text{Rn}$	$^{223}\text{Ra}$	$^{225}\text{Ra}$	$^{223}\text{Fr}$	$^{225}\text{Ac}$	$^{229}\text{Pa}$	$^{199}\text{Hg}$	$^{129}\text{Xe}$
$t_{1/2}$	23.2 m	11.4 d	14.9 d	22 m	10.0 d	1.5 d		
$I$	7/2	3/2	1/2	3/2	3/2	5/2	1/2	1/2
$\Delta e_{\text{th}}$ (keV)	37	170	47	75	49	5		
$\Delta E_{\text{exp}}$ (keV)	--	50.2	55.2	160.5	40.1	0.22		
$10^5 S$ (efm <sup>3</sup> )	1000	400	300	500	900	12000	-1.4	1.75
$10^{28} d_A$ (e cm)	<b>2000</b>	<b>2700</b>	<b>2100</b>	<b>2800</b>			-5.6	0.8

# Experimental EDMs

- Present best limits come from atomic systems and the free neutron
  - Paramagnetic like  $^{205}\text{Tl}$  are primarily sensitive to  $d_e$
  - Diamagnetic atoms (e.g.  $^{199}\text{Hg}$ ) and the free neutron are primarily sensitive to  $\theta_{\text{QCD}}$ ,  $d_q$ ,  $\tilde{d}_q$
- Future best limits may come from
  - Molecules ( $\text{ThO}$ ,  $\text{YbF}$ )
  - Liquids ( $^{129}\text{Xe}$ )
  - Solid State systems (high density)
  - Storage Rings (Muons, Deuteron)
  - Radioactive Atoms ( $^{225}\text{Ra}$ ,  $^{223}\text{Rn}$ )
  - New Technology for Free Neutrons (PSI, ILL, SNS)

# Present and Future EDMs

particle	Present Limit (90% CL) (e-cm)	Laboratory	Possible Sensitivity (e-cm)	Standard Model (e-cm)
$e^-$ (TI)	$1.6 \times 10^{-27}$	Berkeley	$10^{-29}$ $10^{-29}$ $10^{-30}$	$< 10^{-40}$
$e^-$ (PbO)		Yale		
$e^-$ (YbF)		Sussex		
$e^-$ (GGG)		Yale/Indiana		
$\mu$	$9.3 \times 10^{-19}$	CERN	$< 10^{-24}$	$< 10^{-36}$
$\mu$		BNL		
$n$	$3 \times 10^{-26}$	ILL	$1.5 \times 10^{-26}$	$\sim 10^{-32}$
$n$		ILL	$\sim 2 \times 10^{-28}$	
$n$		PSI	$\sim 7 \times 10^{-28}$	
$n$		SNS	$< 1 \times 10^{-28}$	
$^{199}\text{Hg}$	$3 \times 10^{-29}$ (if interpreted as $d_n < 6 \times 10^{-26}$ )	Seattle	$1 \times 10^{-29}$	$\sim 10^{-33}$
$^{129}\text{Xe}$		Princeton	$10^{-31}$	$\sim 10^{-34}$
$^{225}\text{Ra}$		Argonne	$10^{-28}$	
$^{223}\text{Rn}$		TRIUMF	$1 \times 10^{-28}$	
$d$		BNL/JPARC?	$< 10^{-27}$	

# Non-neutron EDMs

- Atomic EDMs for electron EDM
- Atomic EDMs for quark chromo-EDM
- Possible storage ring experiments:
  - In particle rest frame see an electric field

$$\vec{\mathbf{E}} = \frac{\vec{\mathbf{v}} \times \vec{\mathbf{B}}}{c} \quad (\text{Can be large if } \beta \sim 1)$$

Rotates a longitudinally polarized particle into the vertical direction

# $^{205}\text{Tl}$ EDM

Phys. Rev. Lett. 88, 071805 (2002)

B. C. Regan, E. D. Commins,  
C. J. Schmidt, & D. DeMille

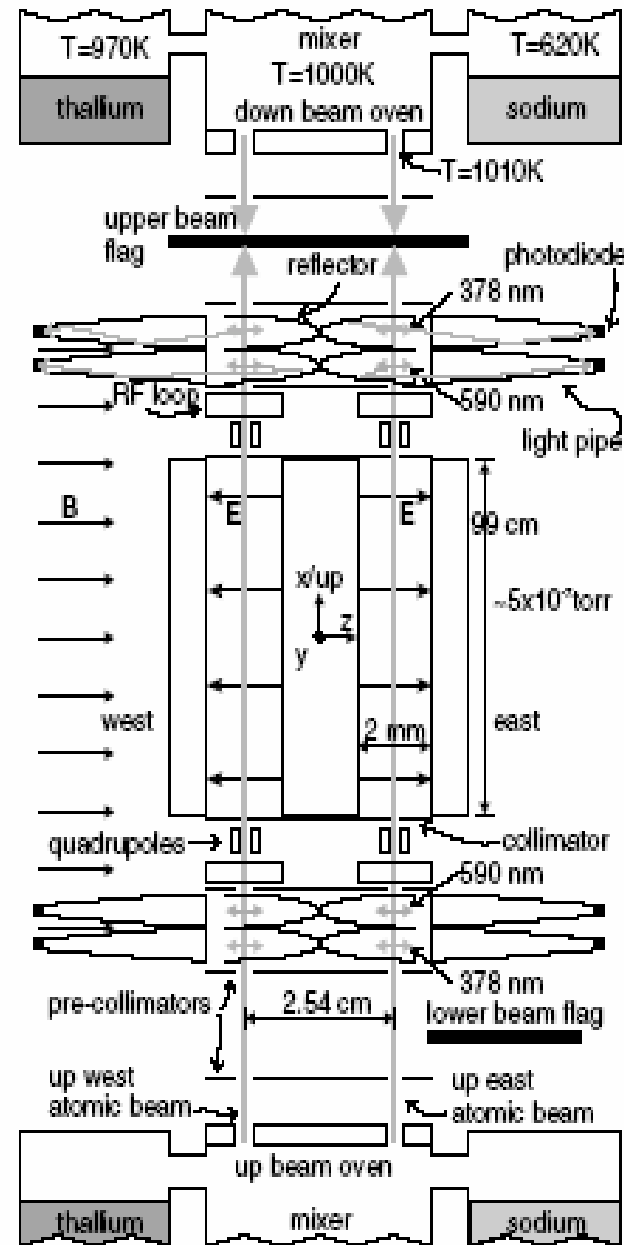
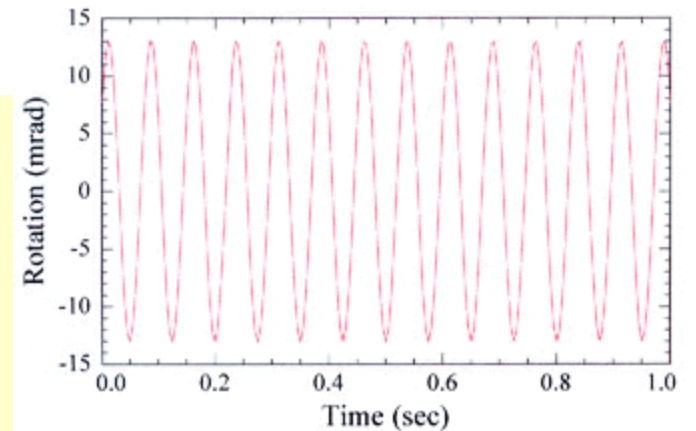
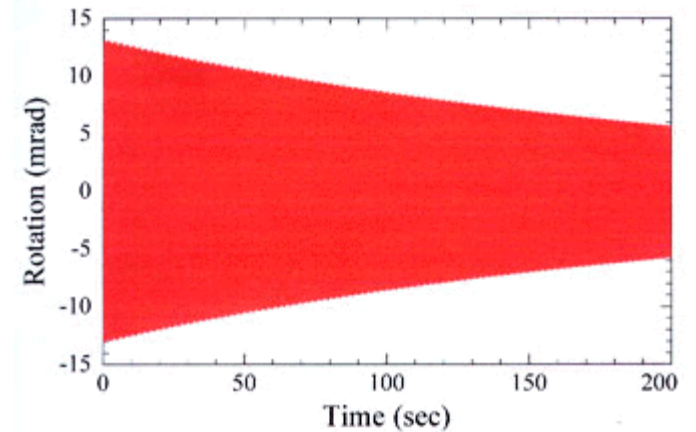
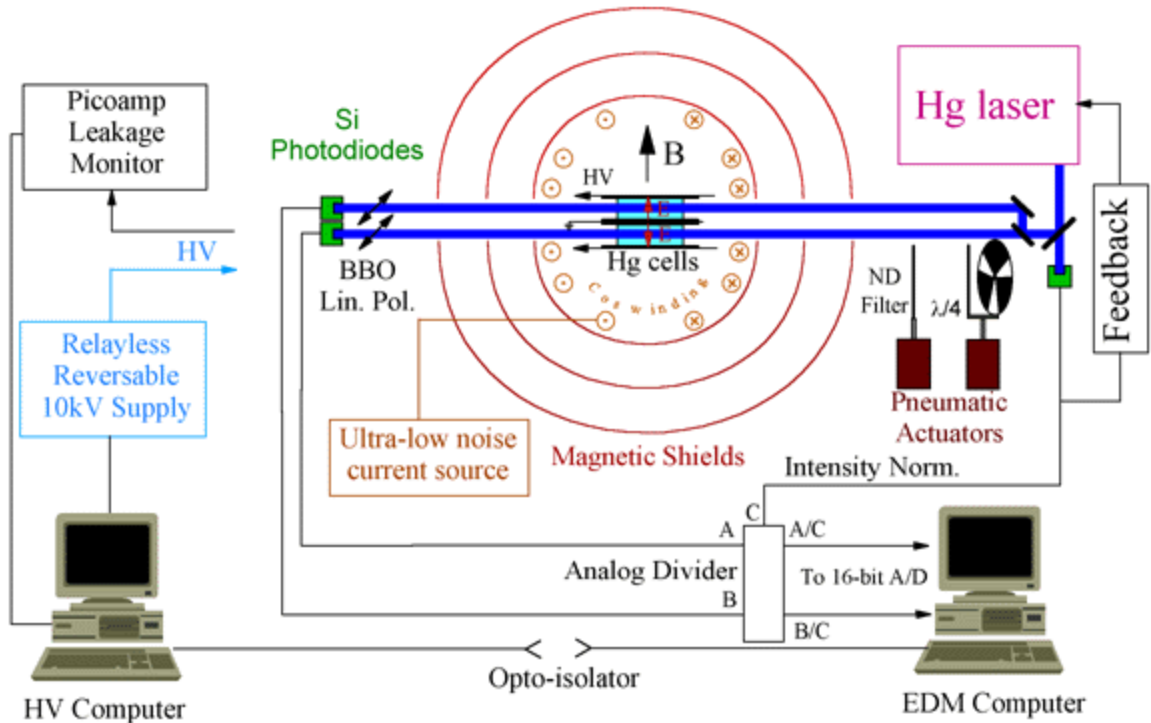


FIG. 1. Schematic diagram of the experiment; not to scale.



# $^{199}\text{Hg}$ EDM

## $^{199}\text{Hg}$ EDM Experimental Setup



Phys. Rev. Lett. **102**, 101601 (2009)  
W. C. Griffith, M. D. Swallows, T. H. Loftus,  
M. V. Romalis, B. R. Heckel, and E. N. Fortson

# EDM with Trapped Radium Atoms

Irshad Ahmad, Roy J. Holt, Zheng-Tian Lu, Elaine C. Schulte, Physics Division, Argonne National Laboratory

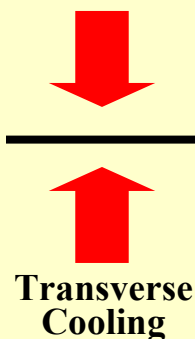
## Advantages of an EDM measurement on $^{225}\text{Ra}$ atoms in a trap

- In  $^{225}\text{Ra}$  the EDM effect is enhanced by two orders of magnitude due to nuclear quadrupole and octupole deformation.
- Trap allows a long coherence time ( $\sim 300$  s).
- Cold atoms result in a negligible “ $v \times E$ ” systematic effect.
- Trap allows the efficient use of the rare and radioactive  $^{225}\text{Ra}$  atoms.
- Small sample in an UHV allows a high electric field ( $> 100$  kV/cm).

10 mCi  
 $^{225}\text{Ra}$  sample



Oven

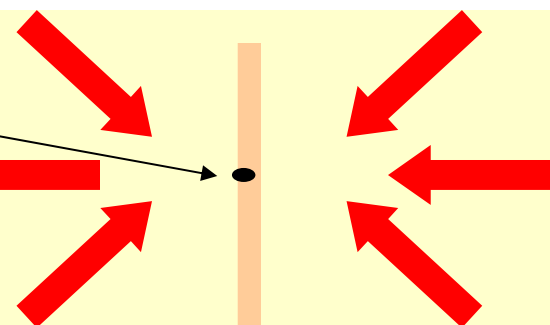


Atomic Beam

Magneto-Optical Trap



EDM-probing region



$\text{CO}_2$ -Laser  
Optical Dipole Trap

$^{225}\text{Ra}$

Nuclear Spin =  $\frac{1}{2}$

Electronic Spin = 0

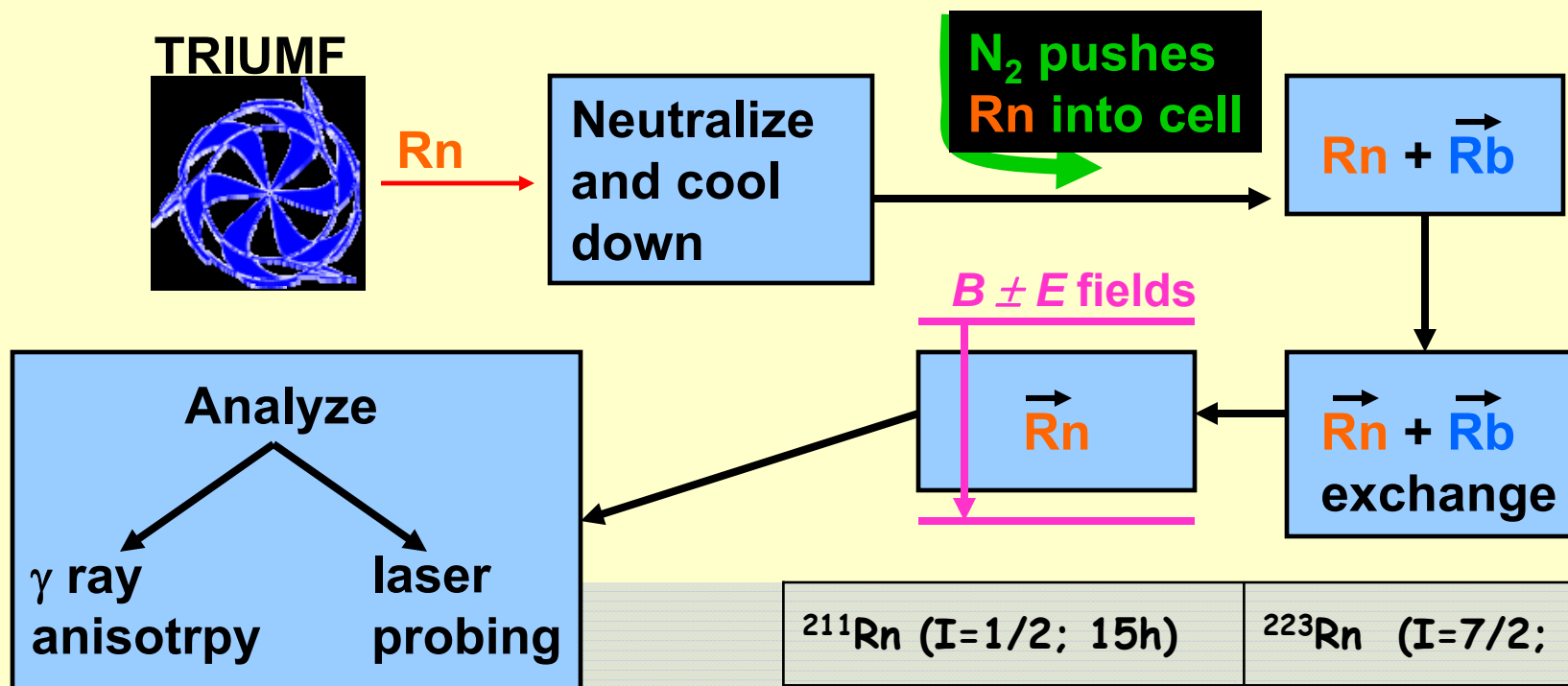
$t_{1/2} = 15$  days

# EDM in Rn

Spokesmen: Timothy Chupp<sup>2</sup> and Carl Svensson<sup>1</sup>

Sarah Nuss-Warren<sup>2</sup>, Eric Tardiff<sup>2</sup>, Kevin Coulter<sup>2</sup>, Wolfgang Lorenzon<sup>2</sup>, Timothy Chupp<sup>2</sup>

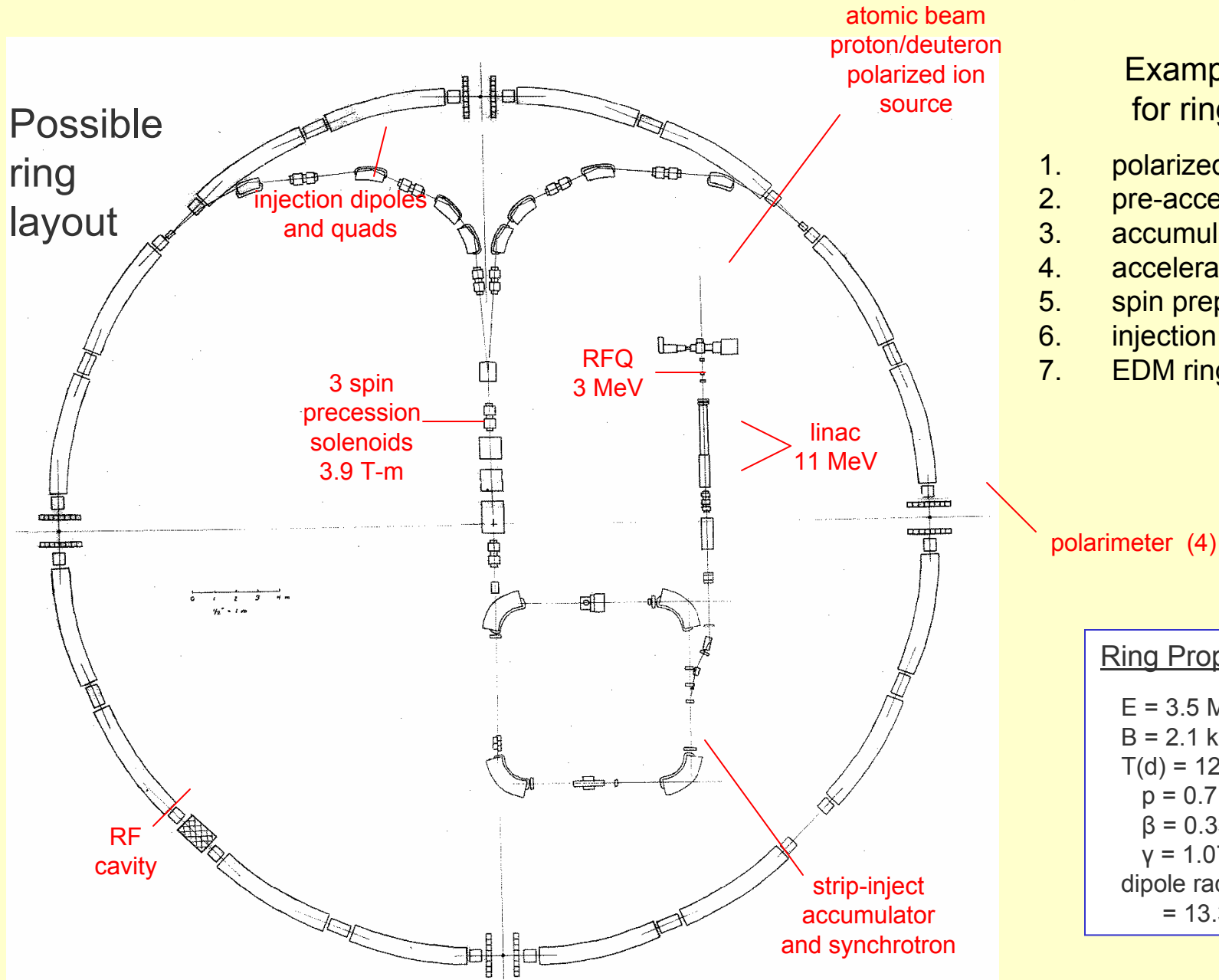
John Behr<sup>4</sup>, Matt Pearson<sup>4</sup>, Peter Jackson<sup>4</sup>, Mike Hayden<sup>3</sup>, Carl Svensson<sup>1</sup>  
 University of Guelph<sup>1</sup>, University of Michigan<sup>2</sup>, Simon Fraser University<sup>3</sup>, TRIUMF<sup>4</sup>



	<sup>211</sup> Rn (I=1/2; 15h)		<sup>223</sup> Rn (I=7/2; 23 m)	
	Rate	$\sigma_d$ ( $\sqrt{\text{day}}$ )	Rate	$\sigma_d$ ( $\sqrt{\text{day}}$ )
TRIUMF	$2 \times 10^9$	$6 \times 10^{-29}$	$1 \times 10^7$	$2 \times 10^{-27}$
RIA	$1 \times 10^{10}$	1 $3 \times 10^{-29}$	$5 \times 10^8$	$1 \times 10^{-28}$

# Deuteron (and Muon) EDM in Storage Ring

BNL, BU, Cornell, Illinois, Indiana, Massachusetts, Oklahoma & Foreign Institutions



Example for ring:

1. polarized ion source
2. pre-accelerator
3. accumulator
4. accelerator
5. spin preparation
6. injection
7. EDM ring

## Ring Properties:

$E = 3.5 \text{ MV/m}$   
 $B = 2.1 \text{ kG}$   
 $T(d) = 126 \text{ MeV}$   
 $p = 0.7 \text{ GeV}/c$   
 $\beta = 0.35$   
 $\gamma = 1.07$   
 dipole radius  
 $= 13.3 \text{ m}$

# Neutron EDM Experiments

- Most recent published result
  - (from ILL)
  - Experiment limited by new systematic effect "discovered" during measurement
- Future experiments
  - 3xILL, PSI, SNS, TRIUMF(?), NIST(?)

# ILL-Grenoble neutron EDM Experiment

Harris et al. Phys. Rev. Lett. 82, 904 (1999)

Baker et al. Phys. Rev. Lett. 97, 131801 (2006)

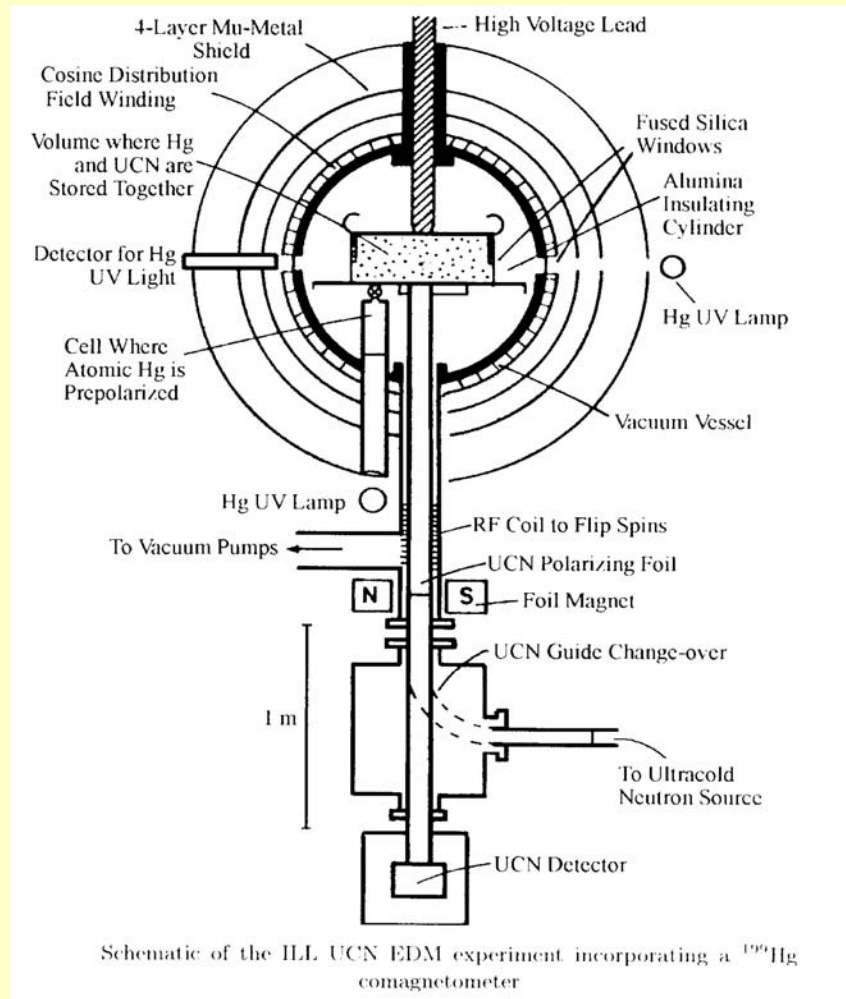
Trapped Ultra-Cold Neutrons (UCN) with  
 $N_{\text{UCN}} = 0.5 \text{ UCN/cc}$

$|E| = 5 - 10 \text{ kV/cm}$

100 sec storage time

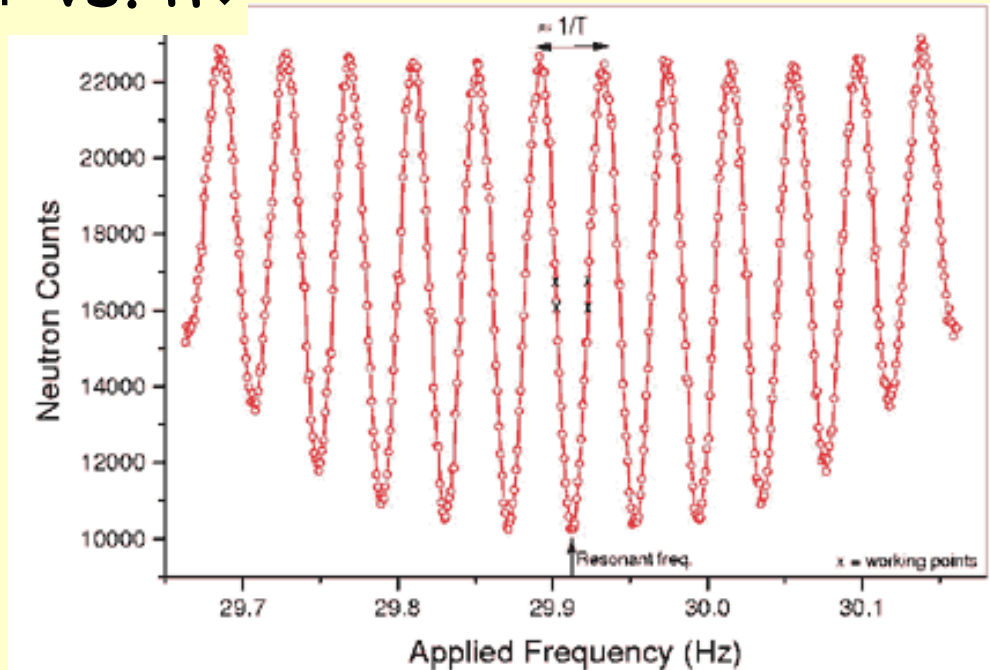


$$\sigma_d = 3 \times 10^{-26} \text{ e cm}$$

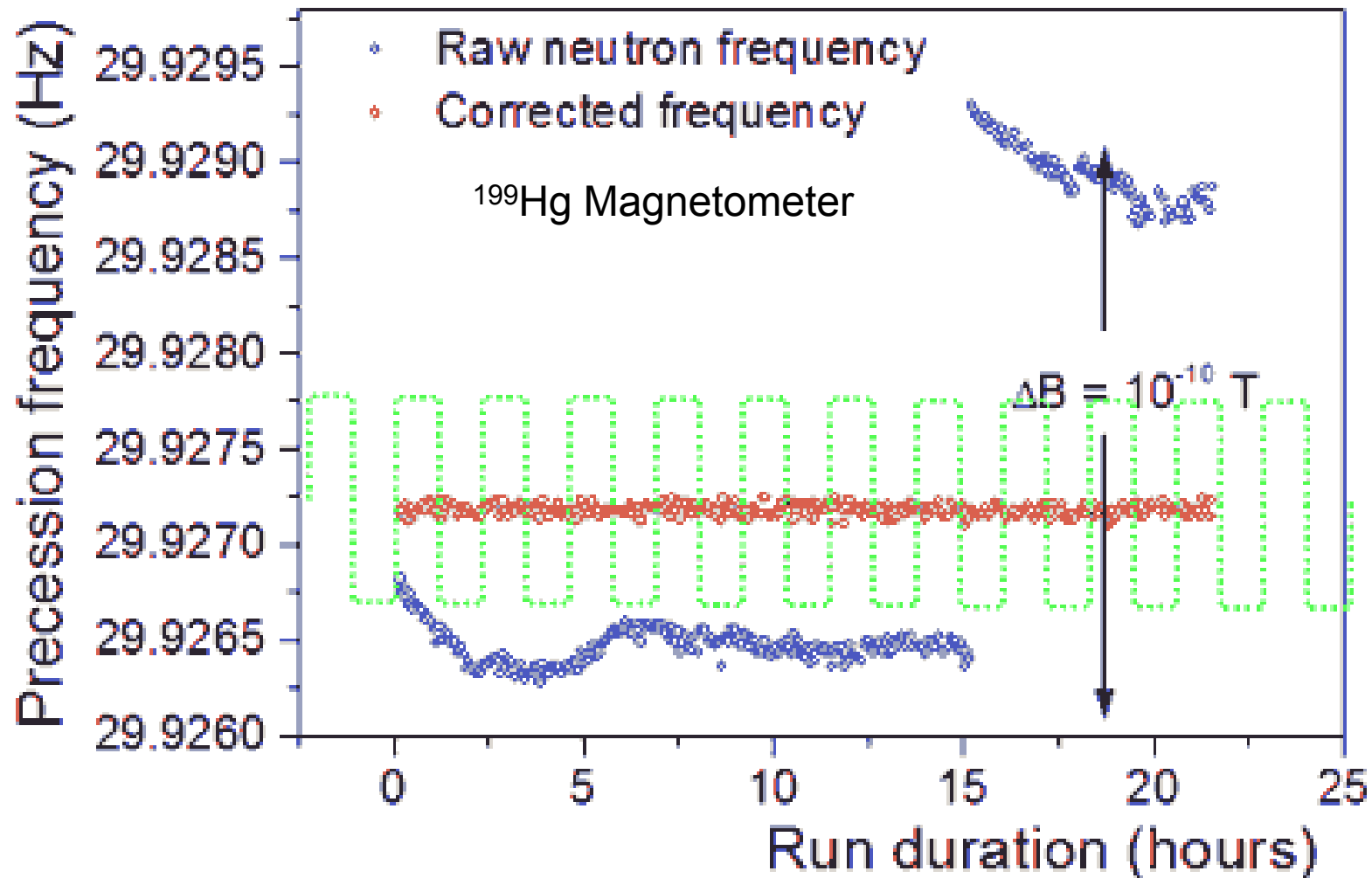


# Measurement of frequency difference

- ILL uses Ramsey separated-oscillatory field technique
  - Inject  $n$
  - Rotate and precess for  $\Delta t$
  - Spin rotates by  $\Delta\omega\Delta t$  (assuming  $\ll 1$ )
  - Measure how many  $n\uparrow$  vs.  $n\downarrow$



# Careful magnetometry is essential !



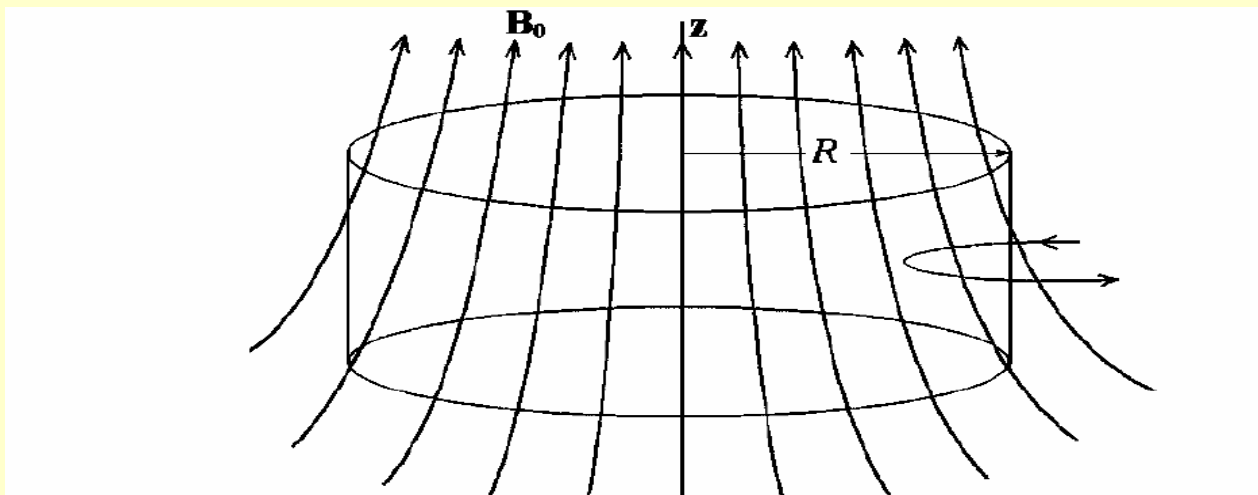


# Experiment limited by systematic effect "Geometric Phase"

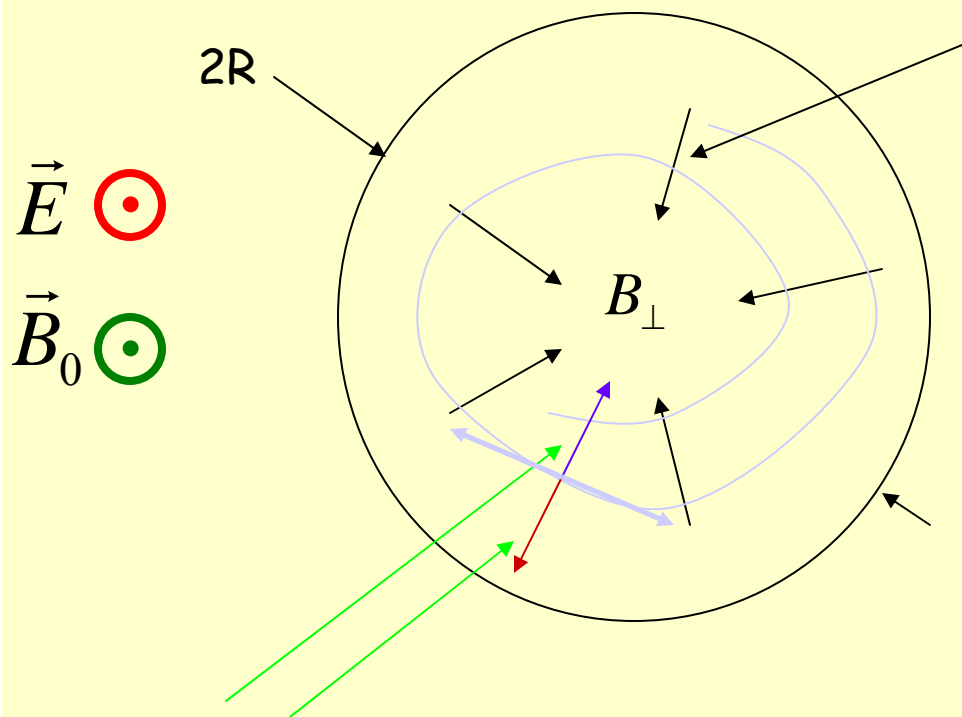
- For slow particles:
  - Path-dependent phase
  - E.g. Parallel transport of vector on sphere
  - In Quantum Mechanics often called Berry's phase
  - Actually a *relativistic* effect!

# False EDM from Geometric phase

- Commins Am J Phys 59, 1077 (91)
- Pendlebury et al PRA 70 032102 (04)
- Lamoreaux and Golub PRA 71 032104 (05)
- Motional ( $v \times E$ ) B-fields can add to radial B fields perpendicular to  $B_0$  (These result e.g. from  $dB_0/dz$ ) giving a false EDM



# Geometric phase with $B_E = v \times E$ field



$v \times E$  field  
changes sign with  
neutron direction

Radial B-field due to gradient

- Motion in B – field shifts the precession frequency -  $\omega_0$  :

$$\Delta\omega \cong \frac{\omega_{\perp}^2 \left[ 1 \pm (2\vec{v}_n \times \vec{E})/cB_{\perp} \right]}{2(\omega_0 \mp v_n/R)}$$

- $\pm, \mp$  due to different trajectories
- Does NOT average to 0
- Is proportional to  $\vec{E}$
- Gives :

$$d_n \approx \frac{v_{\perp}^2 \left| \frac{\partial \mathbf{B}}{\partial \mathbf{z}} \right|}{B_0^2}$$

# Observed in ILL Experiment

$\omega$  depends on E-field  
For neutrons and  
magnetomers

ILL exp. oriented vertically!!

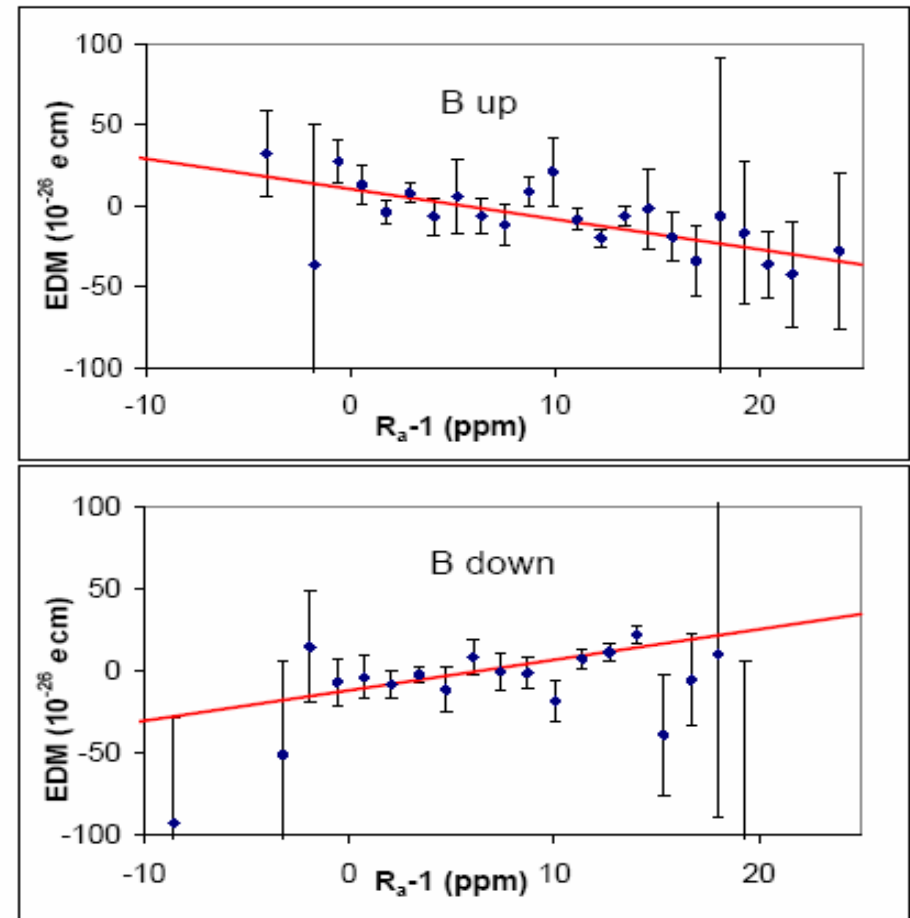


FIG. 2: (Color online) Measured EDM as a function of the relative frequency shift of neutrons and mercury. For clarity, data are binned.

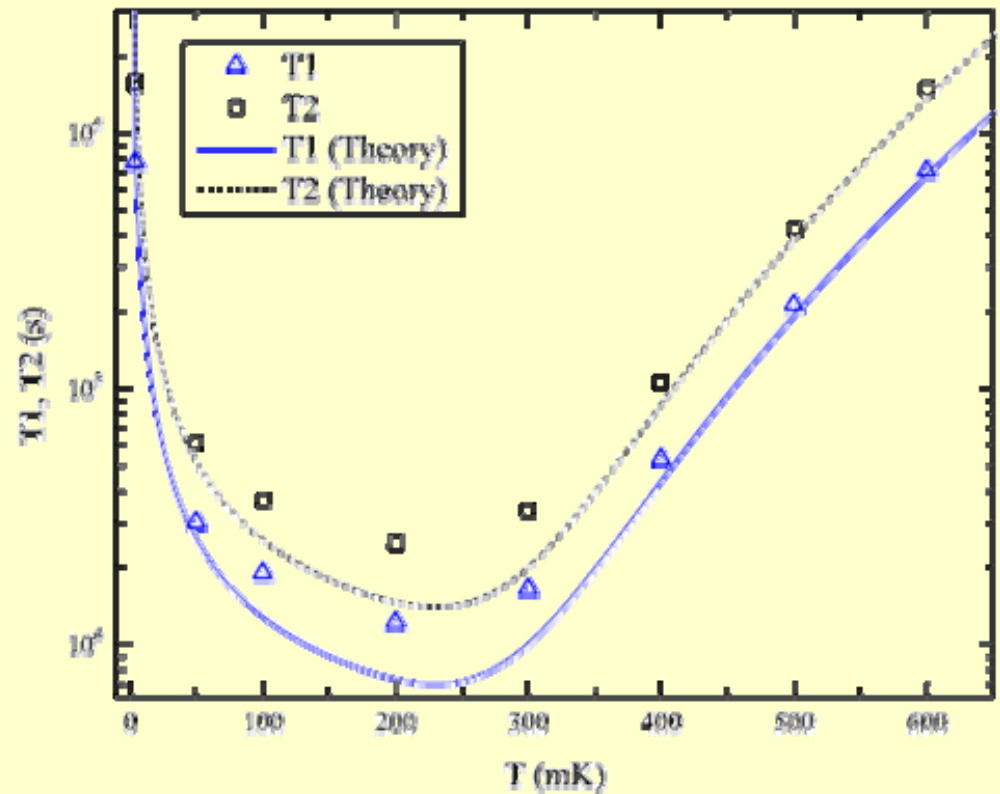
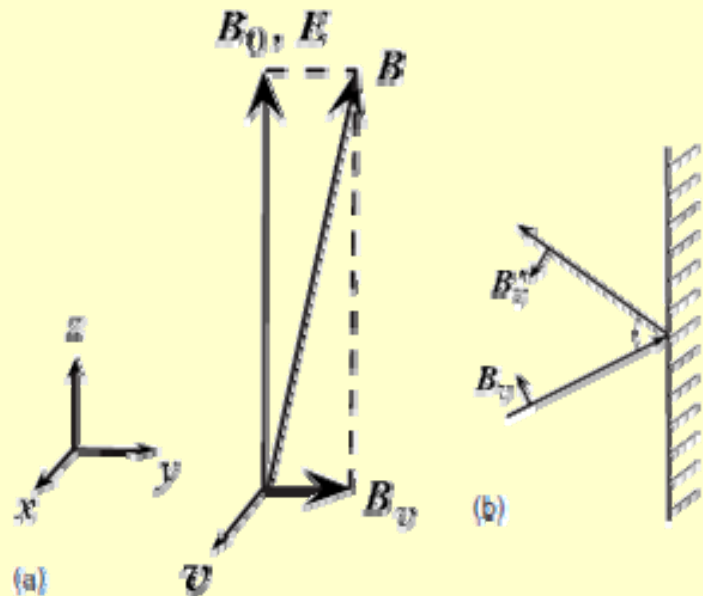
# Related Relativistic Issue

PHYSICAL REVIEW A 78, 023401 (2008)

## Motional spin relaxation in large electric fields

Riccardo Schmid, B. Plaster, and B. W. Filippone  
*California Institute of Technology, Pasadena, California 91125, USA*  
(Received 16 May 2008; published 1 August 2008)

For Polarized  $^3\text{He}$



# To further improve search for neutron EDM, need new techniques

- Enhance number of stored neutrons
- Increase Electric field
- Minimize key systematic effects

# Active worldwide effort to improve neutron EDM sensitivity

- ILL - Grenoble
  - CryoEDM at ILL (superfluid  $^4\text{He}$ )
  - Multiple cell
  - Crystal diffraction of neutron beam
- Paul-Scherrer Institute (PSI) - Switzerland
  - Large Solid  $\text{D}_2$  UCN source
- TRIUMF (possible continuation at JPARC)
  - Superfluid  $^4\text{He}$  source
- Spallation Neutron Source (SNS)  
@ Oak Ridge National Lab
  - Superfluid  $^4\text{He}$

# Example of future Neutron EDM Sensitivity

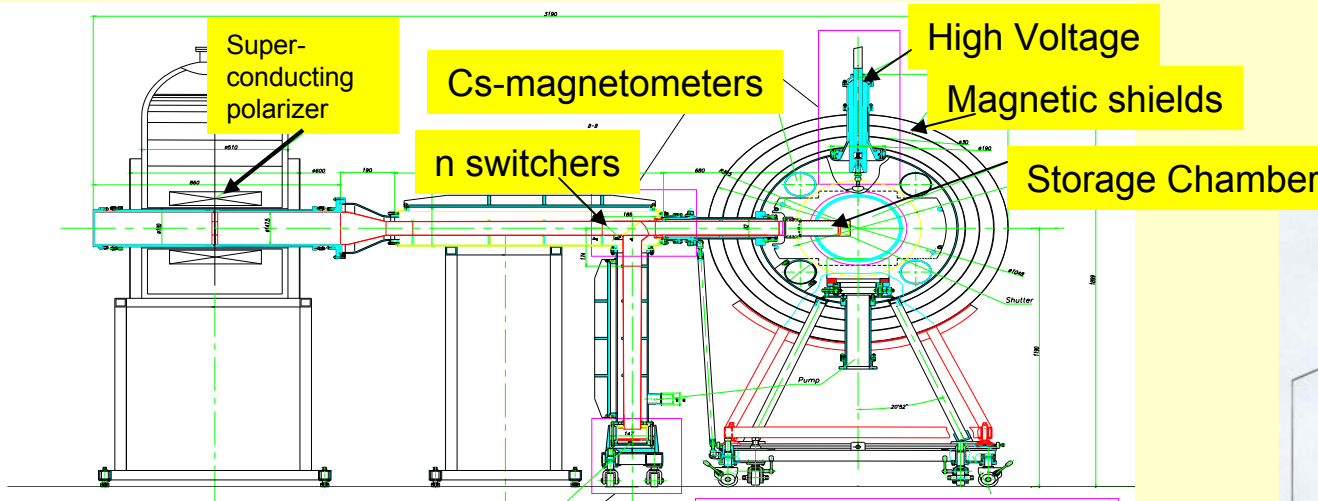
	EDM @ ILL	EDM @ SNS
$N_{\text{UCN}}$	$1.3 \times 10^4$	$4 \times 10^5$
$ \vec{E} $	10 kV/cm	50 kV/cm
$T_m$	130 s	500 s
m (cycles/day)	270	30
$\sigma_d$ (e-cm)/day	$3 \times 10^{-25}$	$8 \times 10^{-27}$

$$\sigma_d \cong \frac{\hbar}{|\vec{E}| T_m \sqrt{m N_{\text{UCN}}}}$$

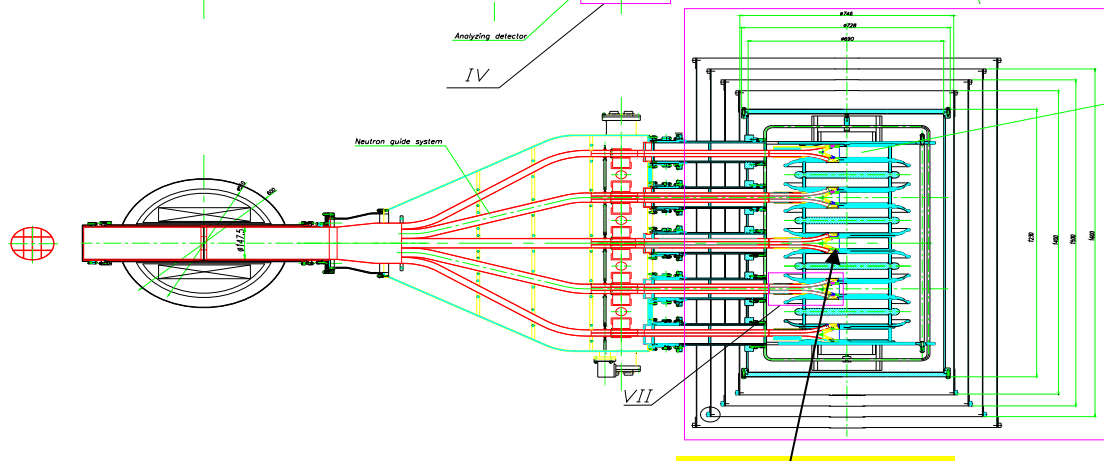
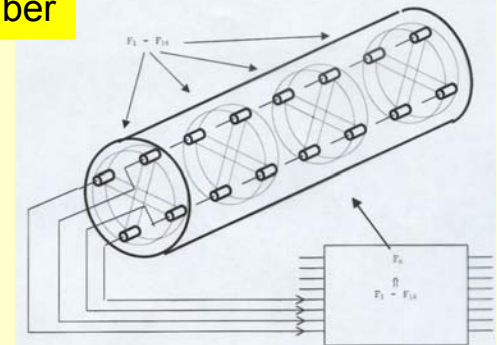


# Scheme of PNPI-ILL multi-chamber EDM spectrometer

Side view



System of 16 Cs-magnetometers

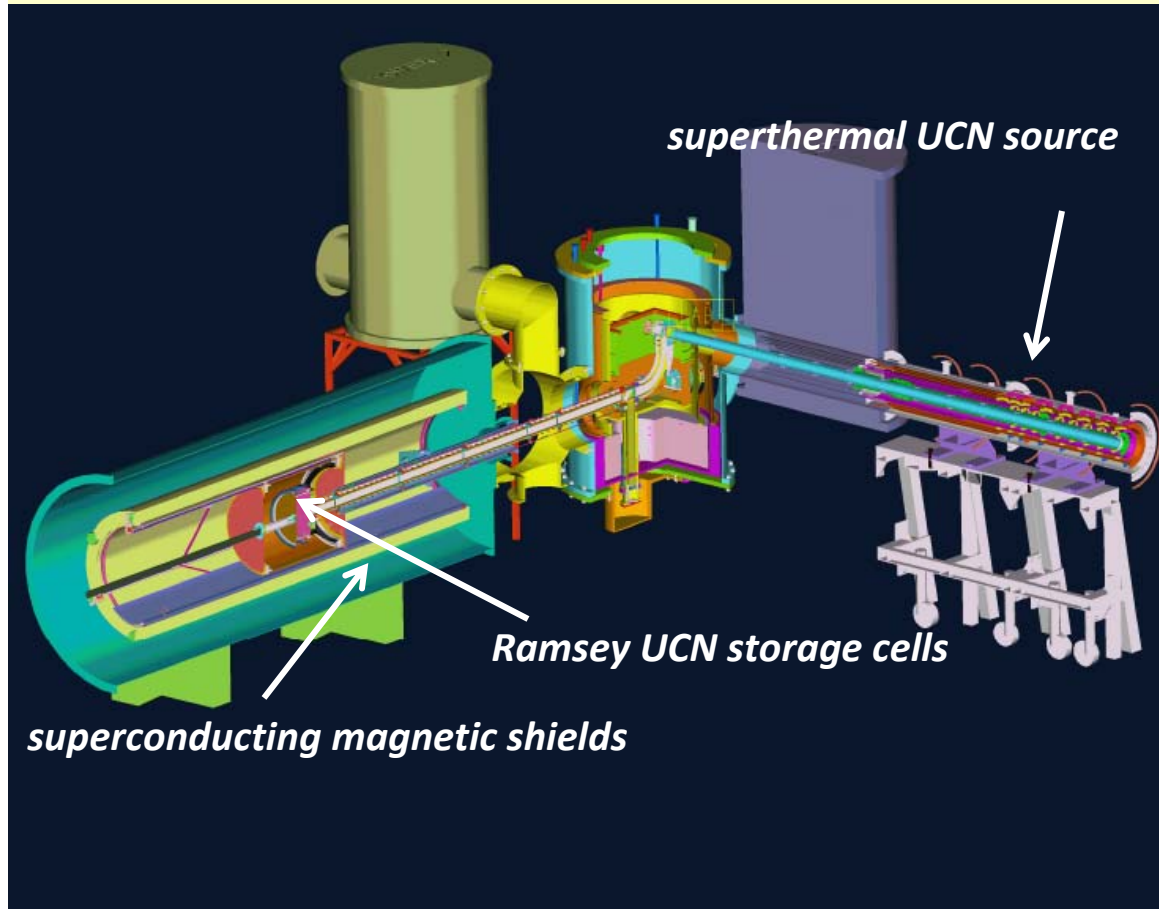


Top view

5 storage cells

Could move experiment to Solid  $D_2$  UCN source at PNPI

# CryoEDM @ ILL

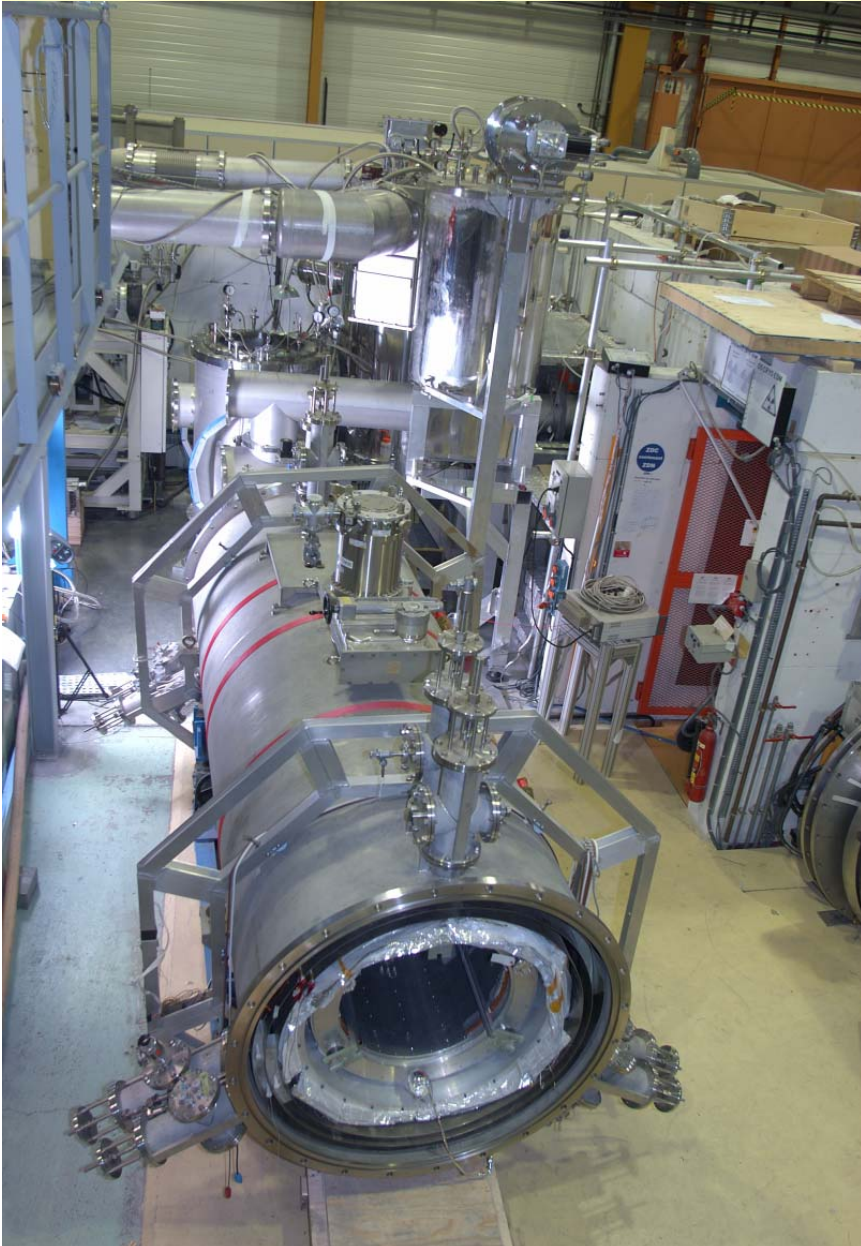


Rutherford Appleton  
Laboratory  
University of Sussex  
University of Oxford  
Institut Laue Langevin  
University of Kure (Japan)

**whole experiment in superfluid He at 0.5 K**

- production of UCN
- storage & Larmor precession of UCN
- SQUID magnetometry
- detection of UCN

# CryoEDM



- Completed constructed, beginning commissioning/start of exploitation
- still requires tuning to deliver a competitive EDM measurement
- apparatus in a position to make an EDM measurement first half 2009 and deliver improved limits

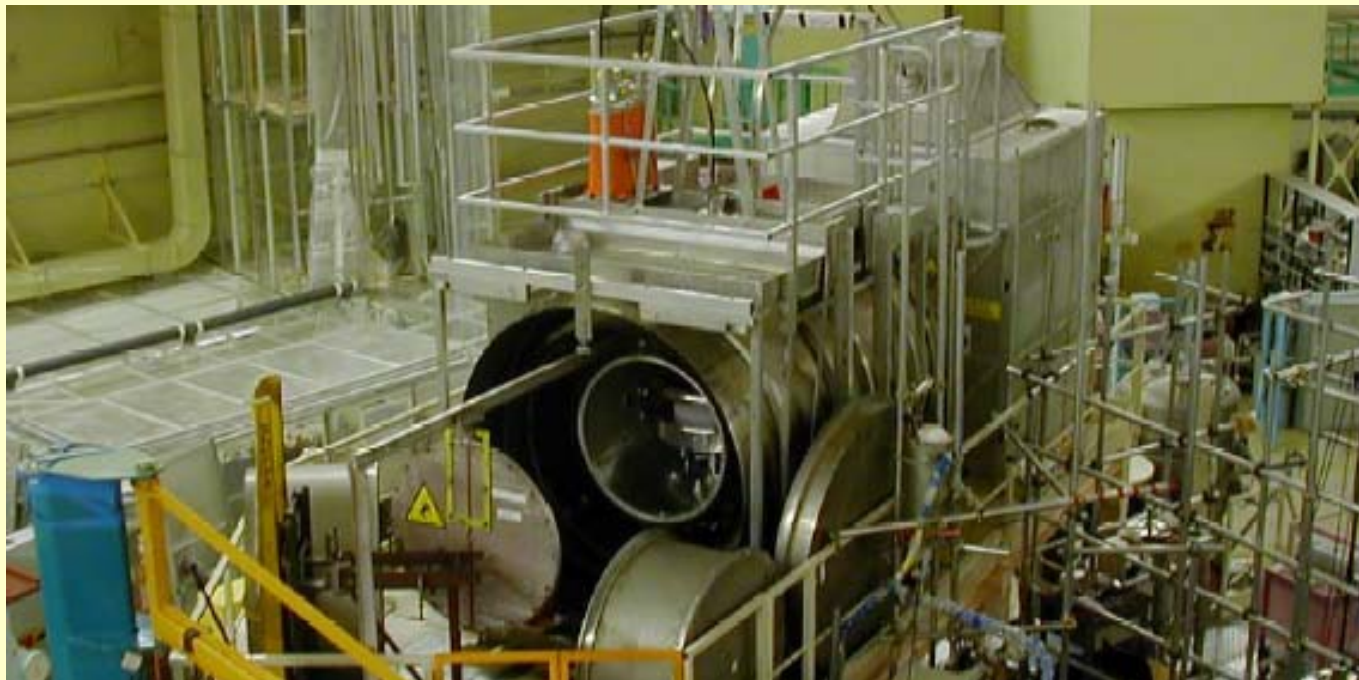
On H53 beam: sensitivity  $\sim 10^{-27} e \text{ cm}$

On new beam: sensitivity  $\sim 10^{-28} e \text{ cm}$

# Neutron EDM at PSI

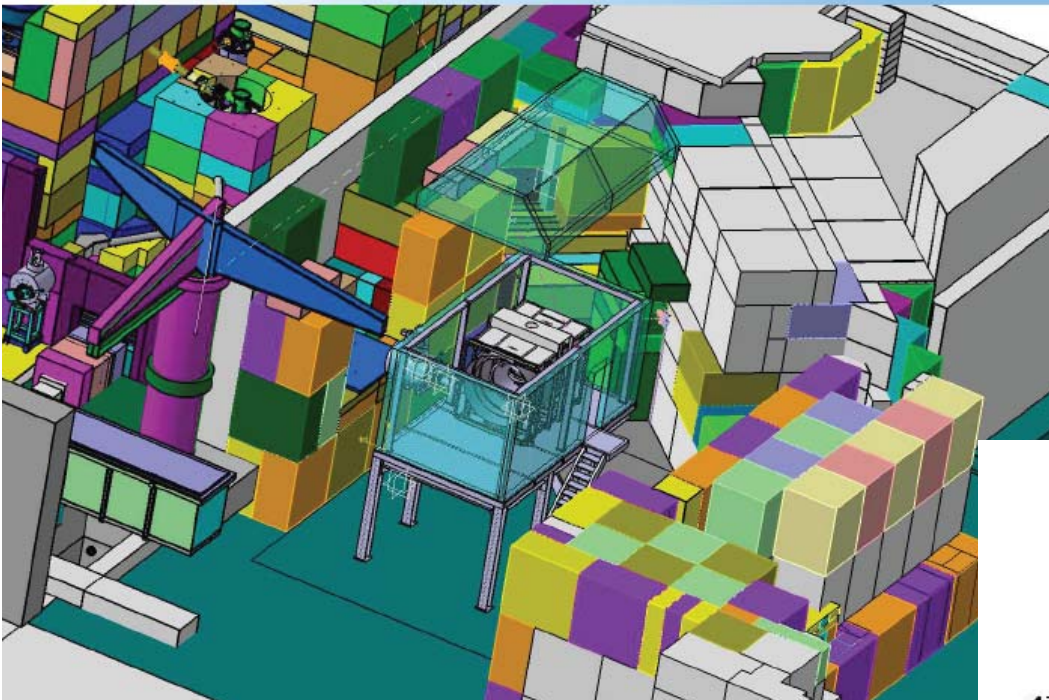
Paul Scherrer Institut

Using new PSI UCN Facility using Solid  $D_2$   
(Based on Los Alamos-et al Concept for UCNA)

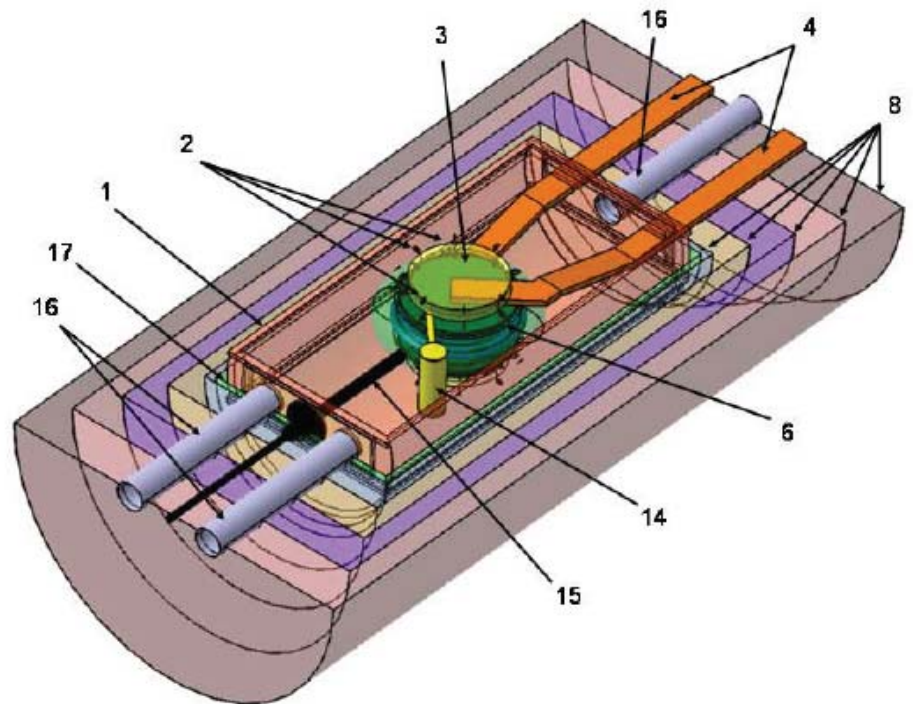


# Neutron EDM at PSI

## PSI UCN area south



- Initial data will use original apparatus from ILL with magnetic upgrades
- New apparatus being designed for higher sensitivity



# Crystal-diffraction neutron EDM project @ ILL

## PNPI

V.V. Fedorov,  
E.G. Lapin,  
I.A. Kusnetsov,  
S.Yu. Semenikhin,  
V.V. Voronin  
Yu.P. Braginets

## ILL

M. Jentschel,  
E. Lelievre-Berna,  
V. Nesvizhevsky,  
A. Petoukhov,  
T. Soldner

Sensitivity

$$\sigma^{-1} \sim E\tau\sqrt{N}$$

Max for UCN method

$$E \sim 10^3 \text{ V/cm}$$

$$\tau \sim 1000 \text{ s (time of life)}$$

$$E\tau \sim 10^7 \text{ (V}\cdot\text{s)/cm}$$

Max for Crystal-diffraction

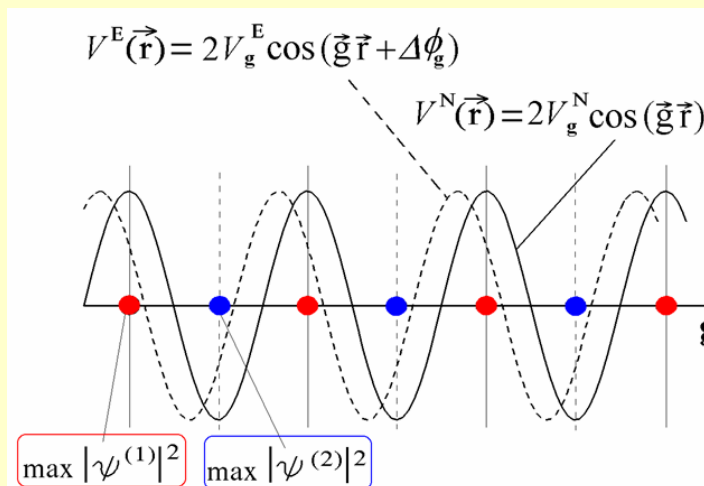
$$E \sim 10^9 \text{ V/cm}$$

$$\tau \sim 10^{-2} \text{ s (time of absorption)}$$

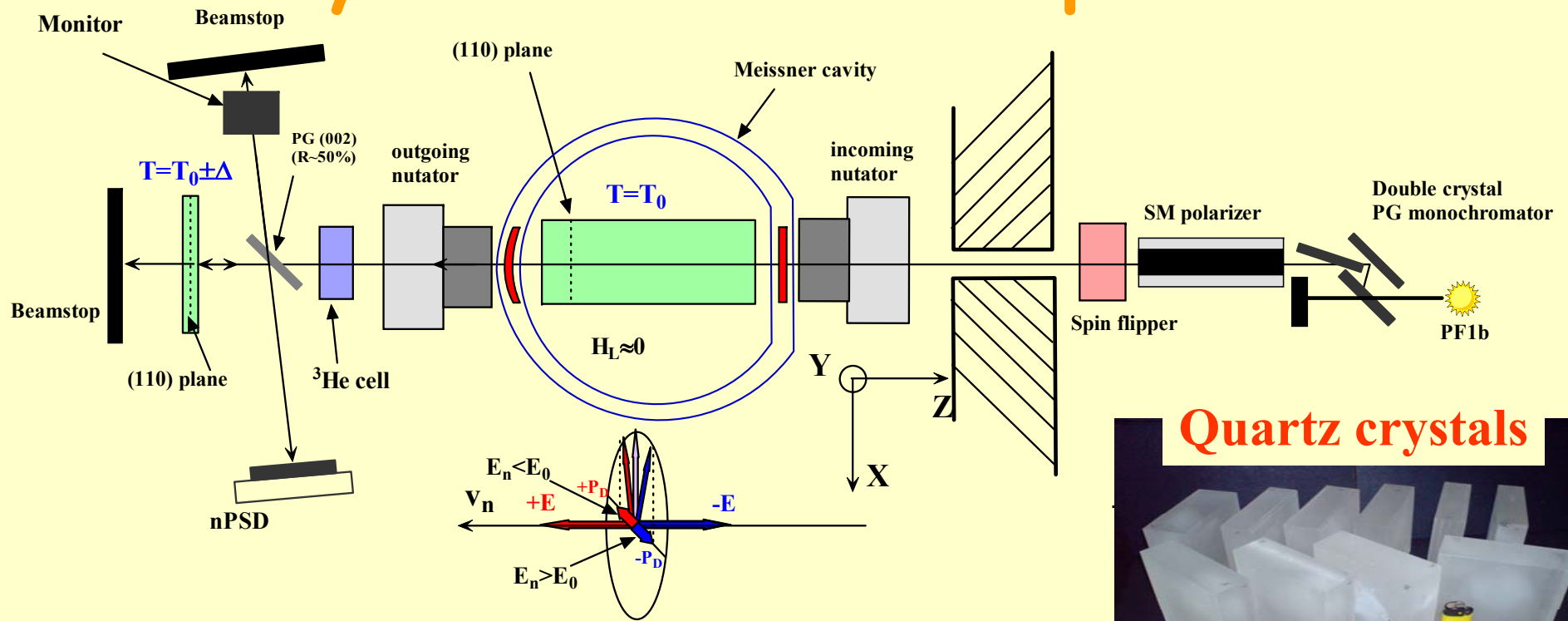
$$E\tau \sim 10^7 \text{ (V}\cdot\text{s)/cm}$$

## In the non-centrosymmetric crystal

neutron is moving under strong electric field if the electric planes deviate from the nuclear ones spatially, because of the neutron concentration on (or between) the nuclear planes



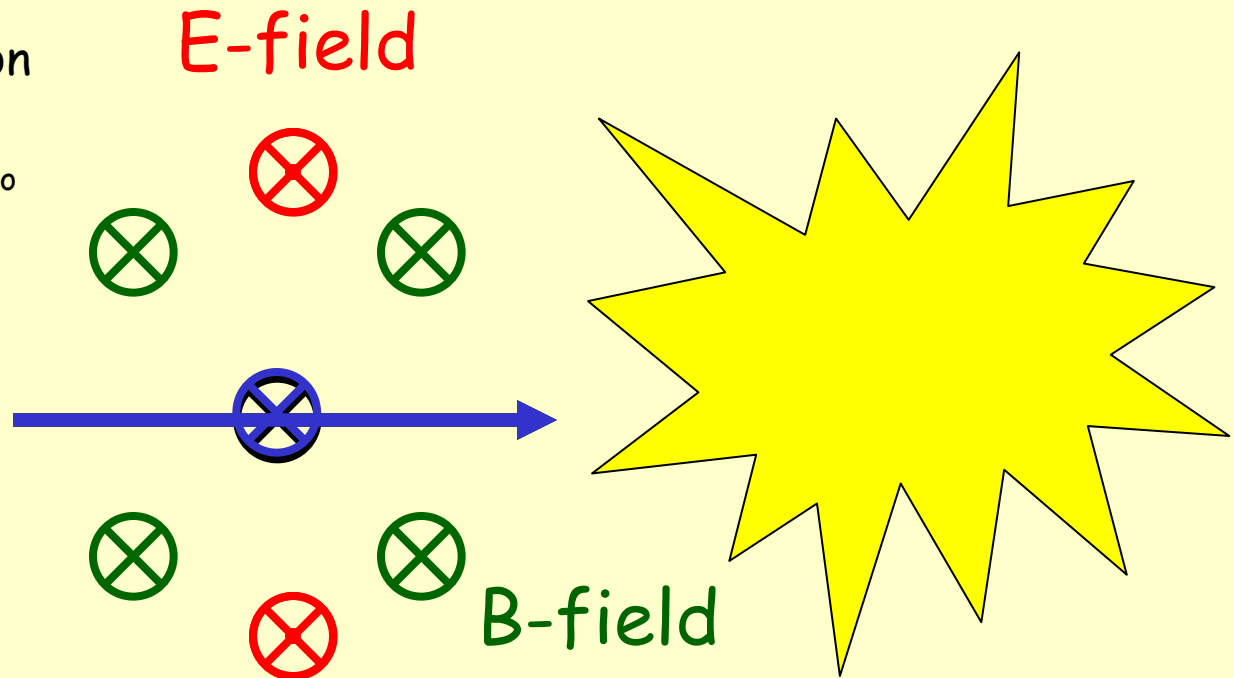
# Layout of the ILL experiment



**NOTE:**  
Centro-symmetric experiment being studied for NIST

# New Technique for n-EDM

1. Inject polarized neutron & polarized  $^3\text{He}$
2. Rotate both spins by  $90^\circ$
3. Measure  $n+^3\text{He}$  capture vs. time  
(note:  $\sigma_{\downarrow\uparrow} \gg \sigma_{\uparrow\uparrow}$ )
4. Flip E-field direction



$^3\text{He}$  functions as "co-magnetometer"



# New Technique for n-EDM:

R. Golub & S. K. Lamoreaux, Phys. Rep. 237, 1 (1994)

- Use Superthermal (non-equilibrium) system to produce UCN
  - Superfluid  $^4\text{He}$  can yield  $\sim 1000$  more UCN than conventional UCN source
- Higher Electric fields in  $^4\text{He}$ 
  - Breakdown voltage may be 10x vacuum breakdown
- $^3\text{He}$  comagnetometer measures B-field at same location as neutrons
  - *Very* small amount of  $^3\text{He}$  in  $^4\text{He}$
  - Use SQUIDS to measure  $^3\text{He}$  precession - calibrates B-field since  $\omega_3 \propto |\vec{B}|$
- $\vec{n} + ^3\vec{\text{He}} \Rightarrow t + p$  has  $\sigma_{\uparrow\downarrow} \gg \sigma_{\uparrow\uparrow}$
- - Detect capture via scintillation of  $^4\text{He}$ 
    - UV photons converted to visible in tetraphenyl butadiene - TPB)
    - Measures difference of  $\omega_n$  and  $\omega_3$
- "Dressed" spin technique suppresses sensitivity to fluctuations in B-field
  - Additional RF field can match  $^3\text{He}$  and neutron precession frequency

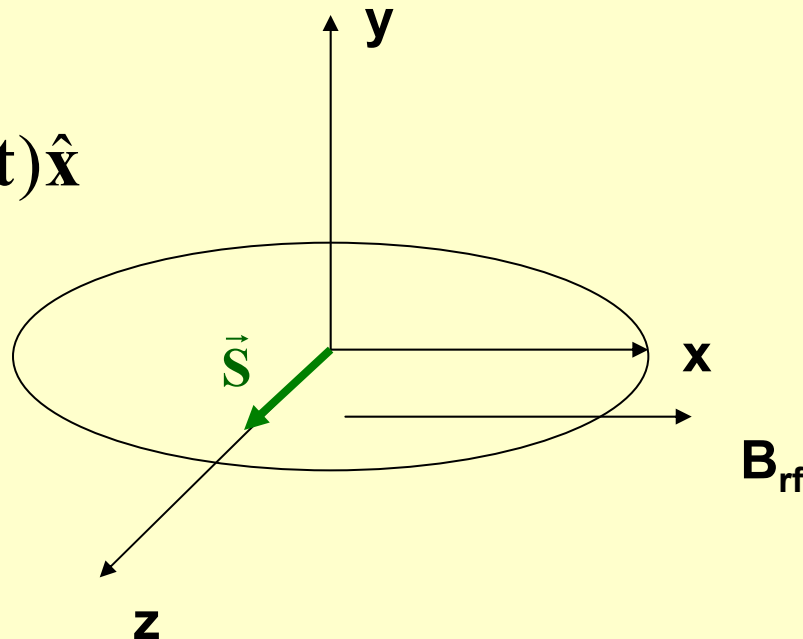
# "Dressed Spins"

- By applying a strong non-resonant RF field, the effective precession frequencies can be modified or "dressed"

# Classical spin in AC B-field

Consider

$$\vec{\mathbf{B}} = \mathbf{B}_{\text{rf}} \sin(\omega_{\text{rf}} \mathbf{t}) \hat{\mathbf{x}}$$



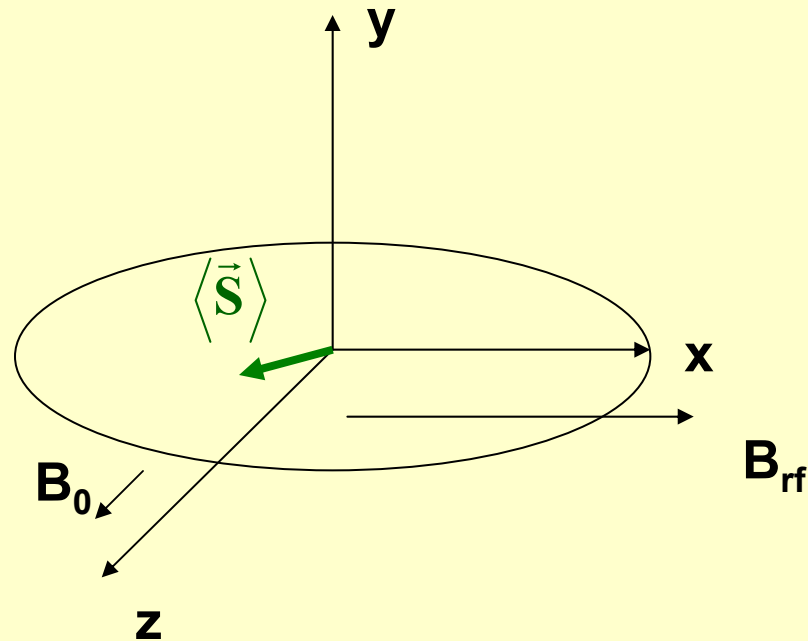
Solution:

$$\mathbf{S}_z = \cos \left[ \left( \frac{\gamma \mathbf{B}_{\text{rf}}}{\omega_{\text{rf}}} \right) \sin(\omega_{\text{rf}} \mathbf{t}) \right] \quad ; \quad \& \text{ averaging over time :}$$

$$\langle \mathbf{S}_z \rangle = \frac{\omega_{\text{rf}}}{2\pi} \int_0^T \mathbf{S}_z \, d\mathbf{t} \equiv \frac{1}{2\pi} \int_0^{2\pi} \cos \left[ \frac{\gamma \mathbf{B}_{\text{rf}}}{\omega_{\text{rf}}} \sin(\theta) \right] d\theta \equiv \mathbf{J}_0 \left( \frac{\gamma \mathbf{B}_{\text{rf}}}{\omega_{\text{rf}}} \right) = \mathbf{J}_0(\mathbf{x})$$

# Classical spin in AC B-field

- Now apply a very small B-field along  $z = B_0$



- Reduced spin  $\langle S_z \rangle = J_0(\mathbf{x})$  begins to precess about z-axis with reduced frequency  $\sim \gamma B_0 J_0(\mathbf{x})$

# Classical spin in AC B-field

- For particular values of the dressing field, the neutron and  $^3\text{He}$  precession frequencies are equal

$$\gamma_3 \mathbf{J}_0 \left( \frac{\gamma_3 \mathbf{B}_{\text{rf}}}{\omega_{\text{rf}}} \right) = \gamma_n \mathbf{J}_0 \left( \frac{\gamma_n \mathbf{B}_{\text{rf}}}{\omega_{\text{rf}}} \right)$$

- Can modulate the dressing field around a relative precession of zero.
  - Reducing effect of external B-fields
  - Measure this parameter vs. direction of E-field
- Challenging technical issues must be overcome
  - Uniformity of the RF field must be better than 0.1%
  - Eddy currents will heat conductors

# The SNS nEDM Collaboration

## Expertise:

Nuclear

Atomic

Condensed Matter

Low Temperature

Polarized  $^3\text{He}$

UCN

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G. Greene

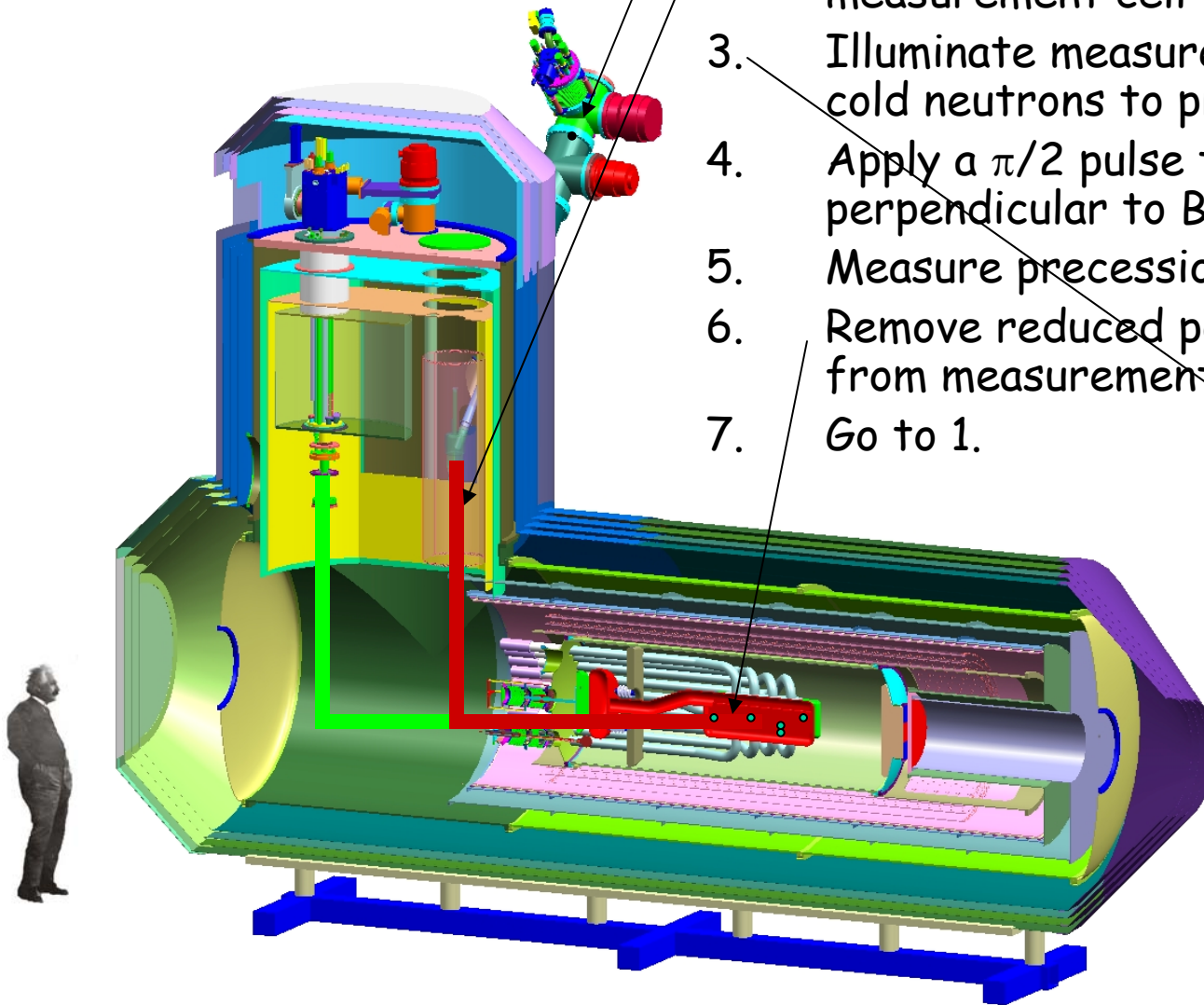
*University of Tennessee, Knoxville, TN, USA*

S. Lamoreaux, D. McKinsey, A. Sushkov

*Yale University, New Haven, CT 06520, USA*

# SNS nEDM Measurement cycle

1. Load collection volume with polarized  $^3\text{He}$  atoms
2. Transfer polarized  $^3\text{He}$  atoms into the measurement cell
3. Illuminate measurement cell with polarized cold neutrons to produce polarized UCN
4. Apply a  $\pi/2$  pulse to rotate spins perpendicular to  $B_0$
5. Measure precession frequency
6. Remove reduced polarization  $^3\text{He}$  atoms from measurement cell
7. Go to 1.



# Systematic Effects in EDM

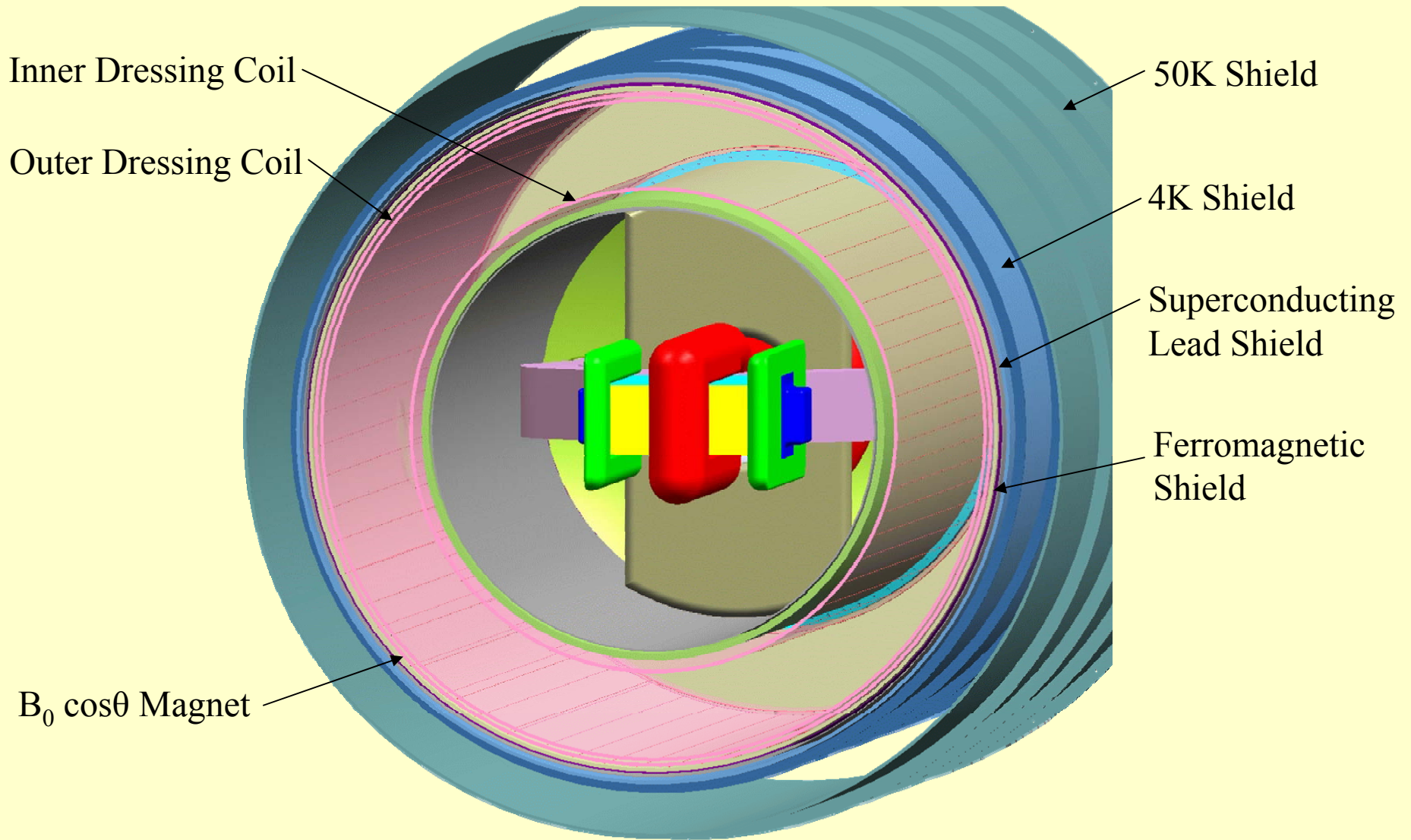
- Variation of B-field
  - Comagnetometer cancels B-field variations
- Leakage currents from Electric Field
  - These produce B-fields that change with E-field (must be less than picoAmps)
- Gravitational offset of n and  ${}^3\text{He}$  ( $\sim 10^{-29}$  e-cm)
- $\vec{v} \times \vec{E}$  effects are the largest sources of systematic error in present ILL exp.
  - $\vec{B}_E = \vec{v} \times \vec{E} \rightarrow$  changes  $\vec{\mu}$  precession frequency
  - Geometric phase due to  $\vec{B}$  gradients



# Systematic Controls in new EDM experiment

- Highly uniform E and B fields
  - $\cos\theta$  coil in Ferromagnetic shield
  - Kerr effect measurement of E-field
- Two cells with opposite E-field
- Ability to vary influence of  $B_0$  field
  - via "dressed spins" (atomic physics trick)
- Control of central temperature
  - Can vary  $^3\text{He}$  diffusion

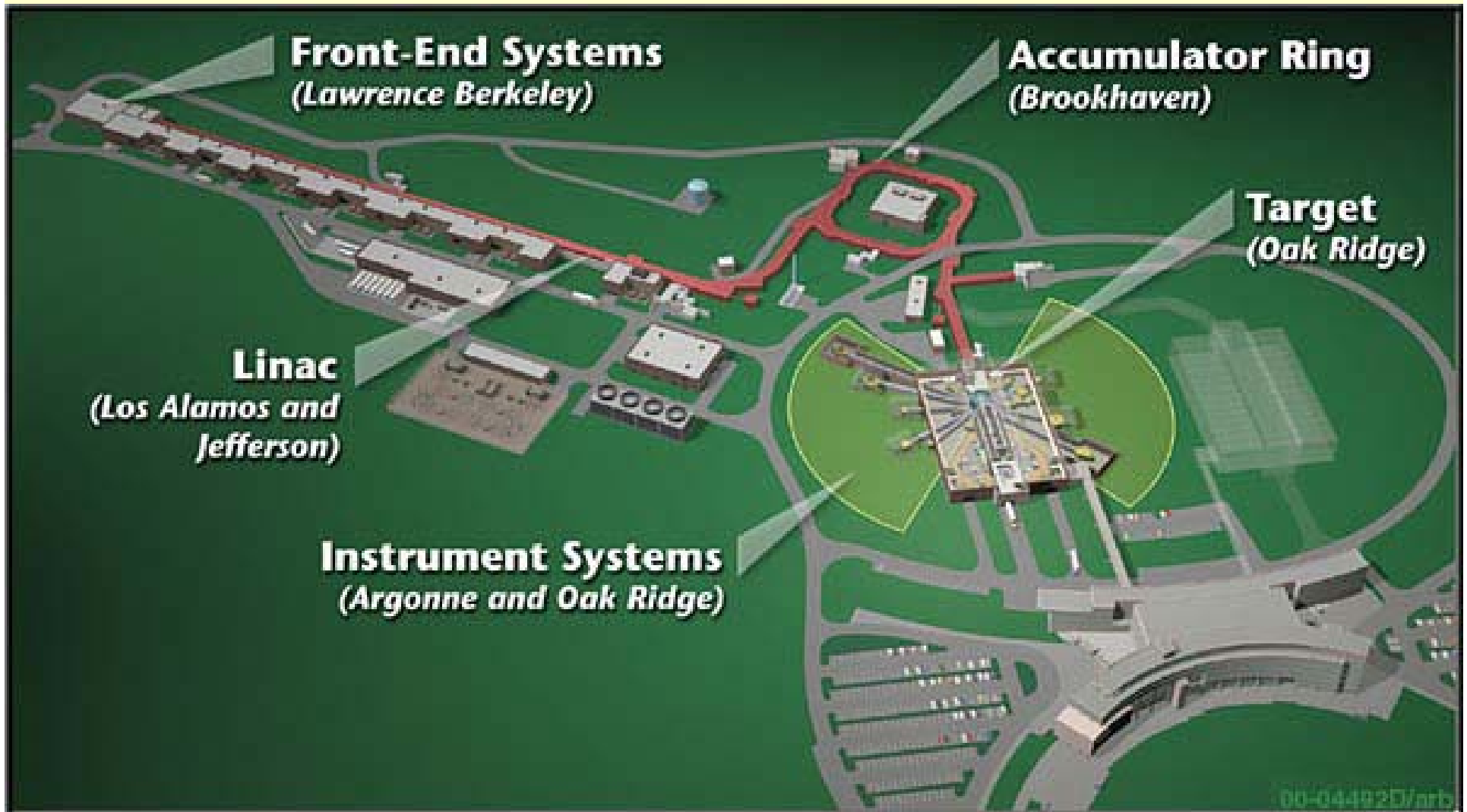
# Measurement Cell



**Neutrons come from Oak Ridge National Laboratory**

# Spallation Neutron Source (SNS) at ORNL

1 GeV proton beam with 1.4 MW on spallation target



# SNS Status



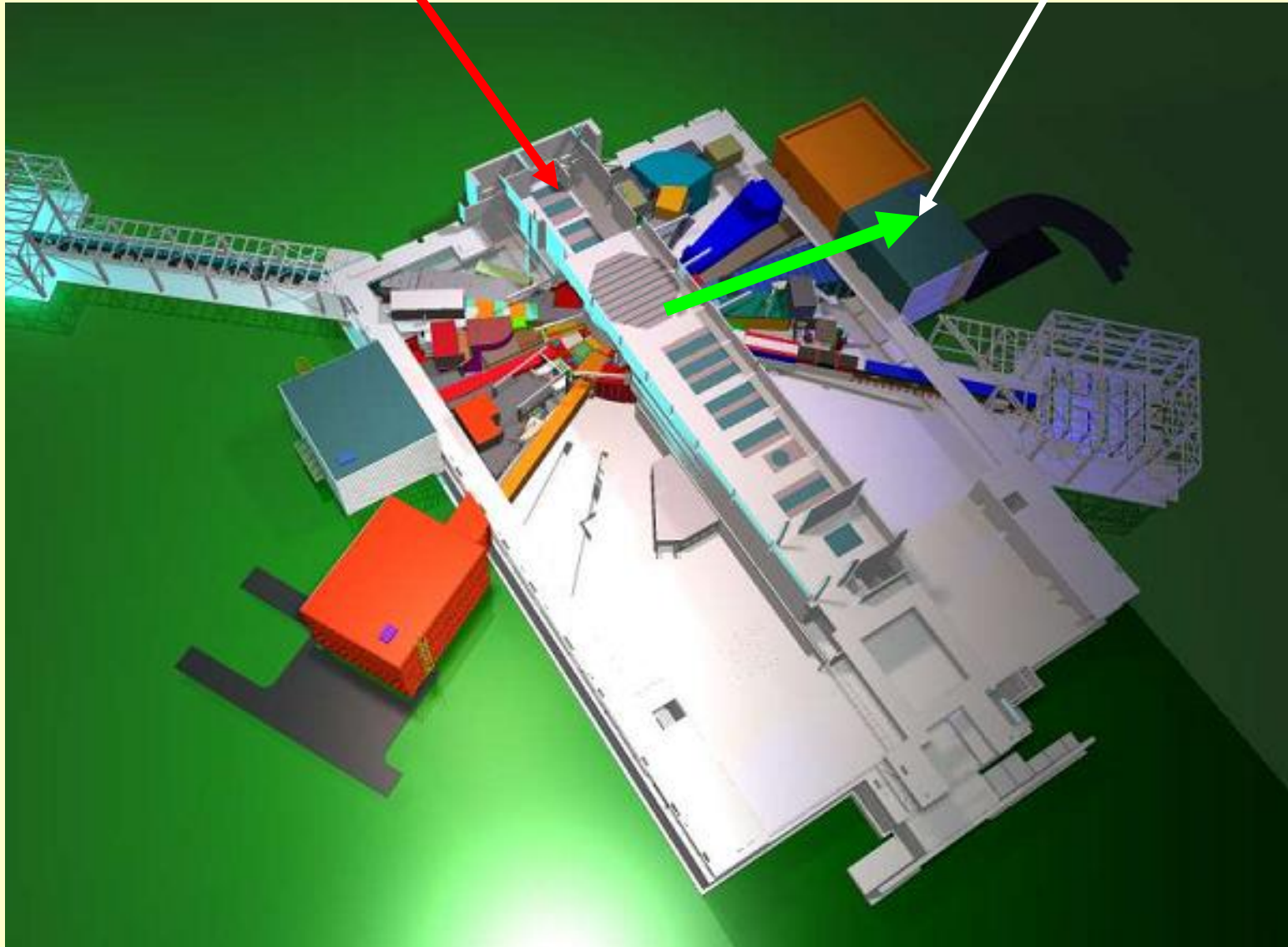
Photo courtesy of ORNL

- **SNS completed:** 2006
- **Beam line completed:** 2007
- **Full design flux:** 2009
- **SNS Total Project Cost:** 1.411B\$

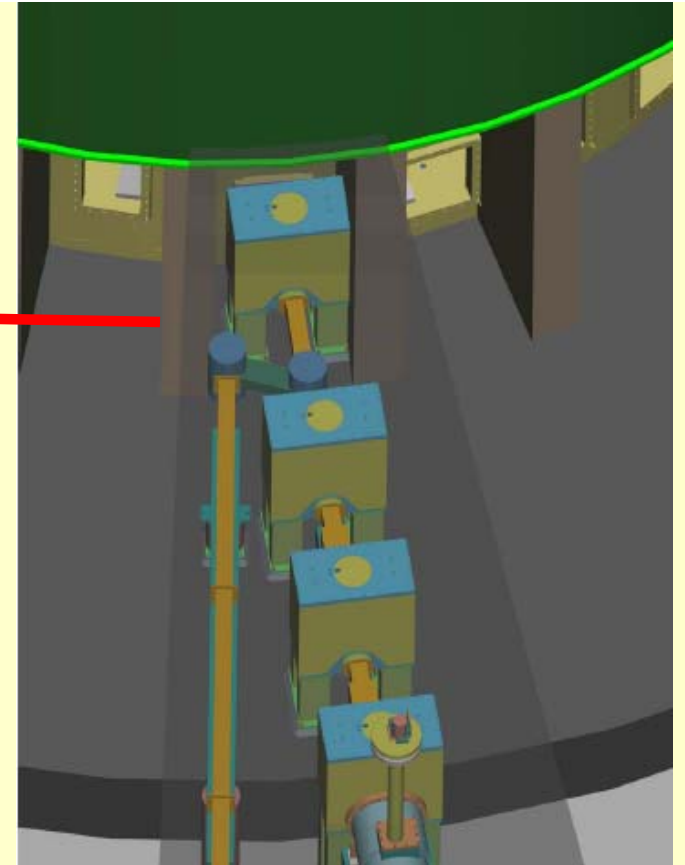
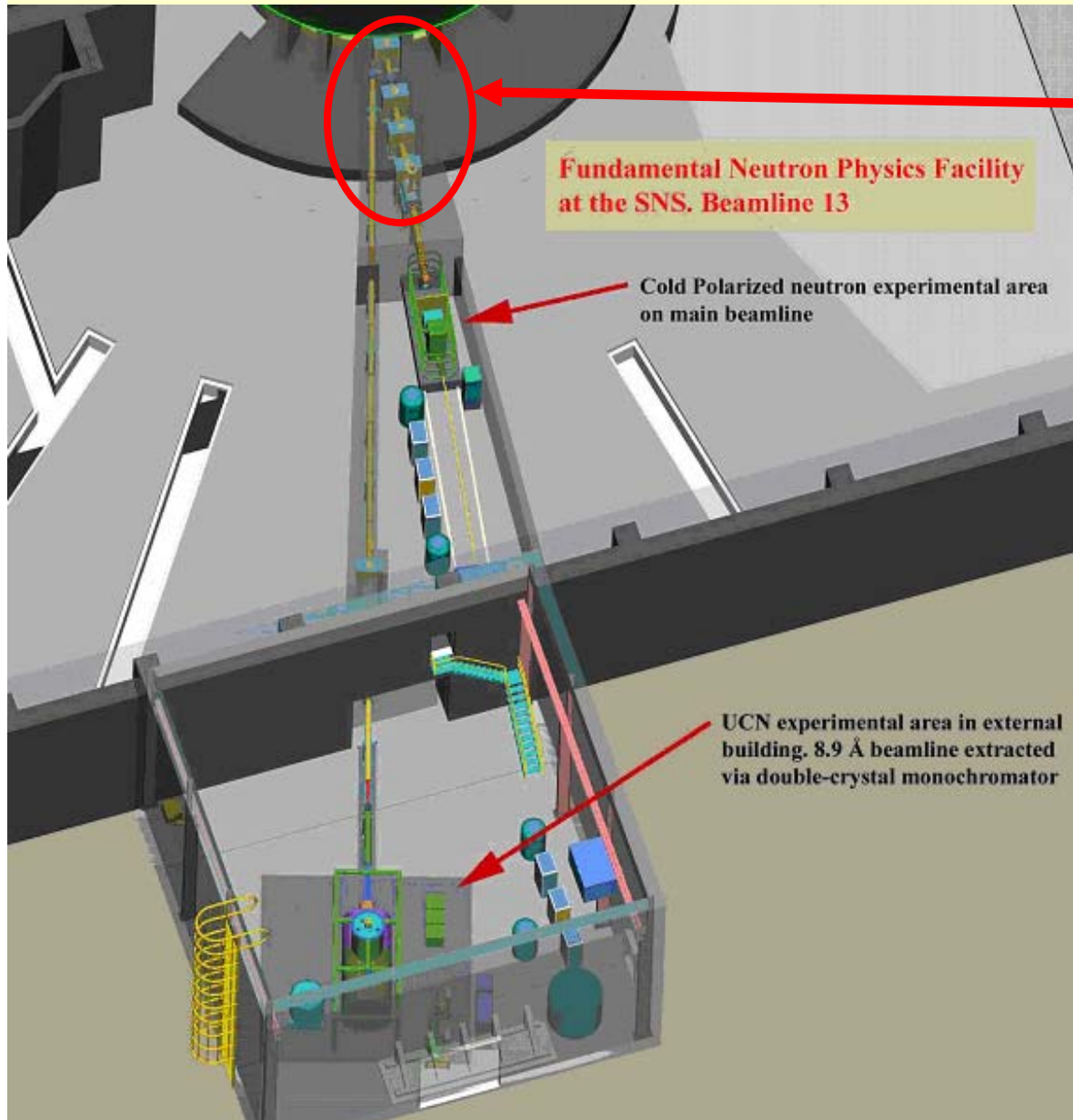
# SNS Target Hall

18 neutron beam ports with 1 for Nuclear Physics

p beam



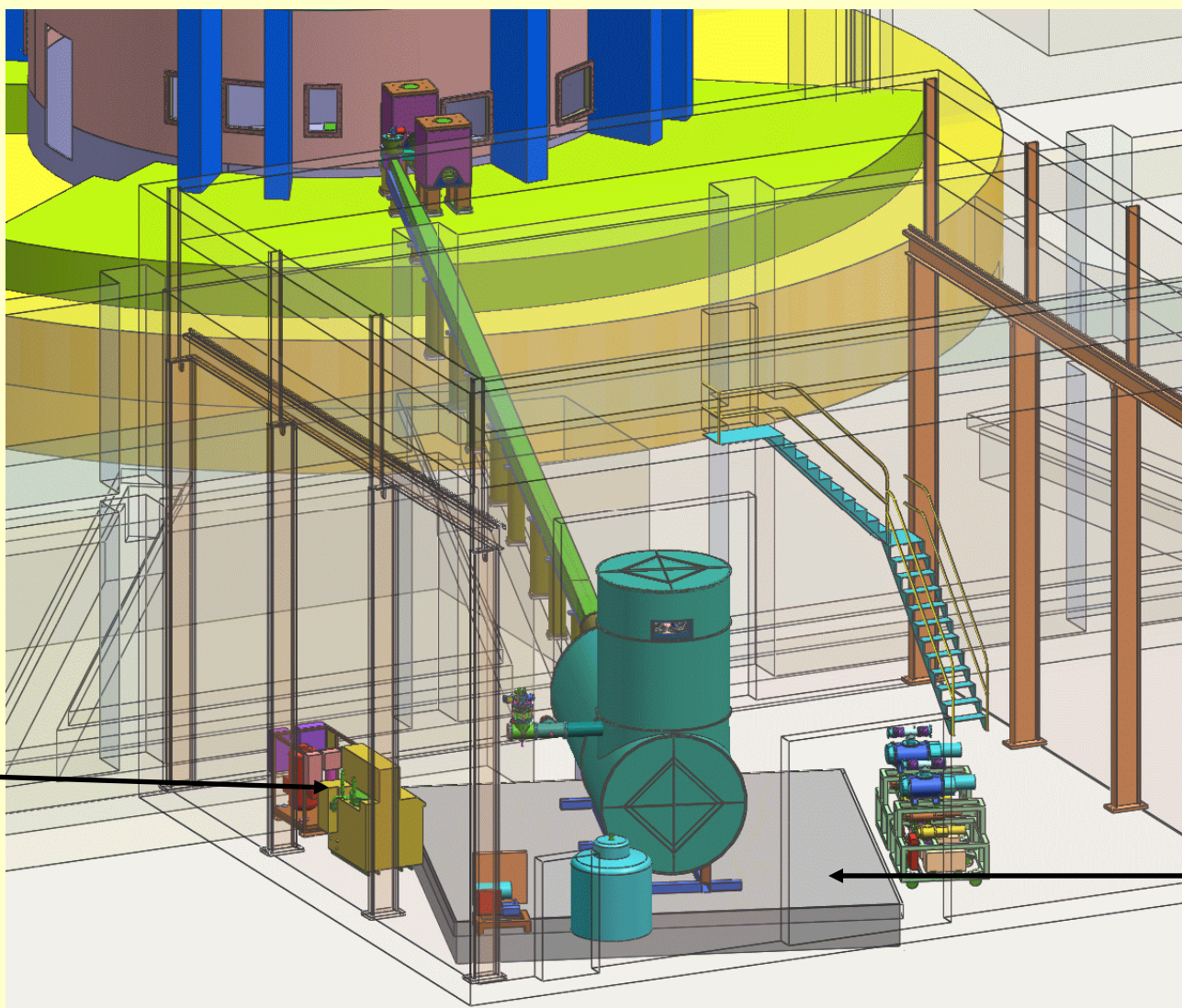
# Fundamental Neutron Physics Beamline



Double monochromator  
selects 8.9 Å neutrons

# EDM Experiment at SNS

He  
Liquifier



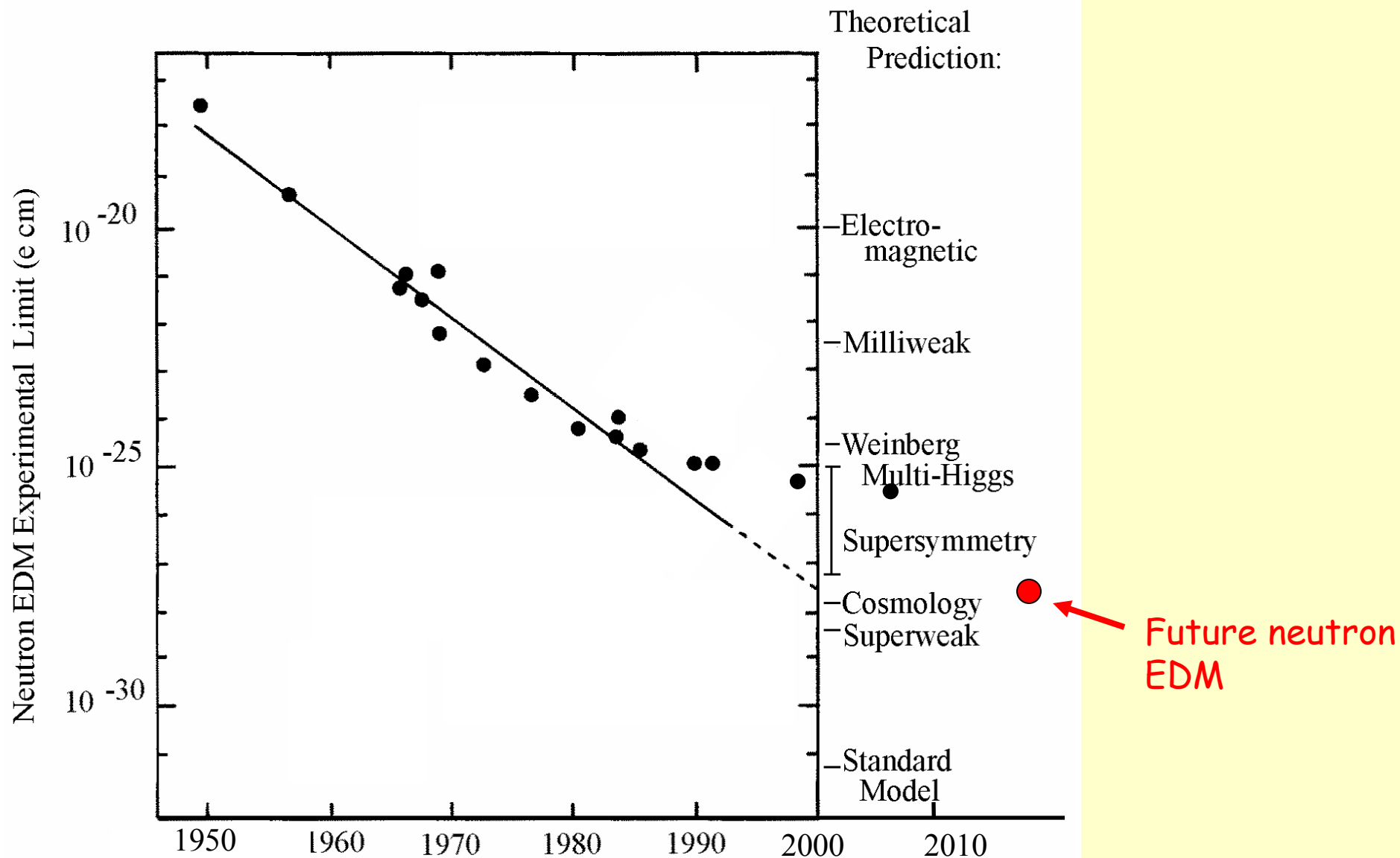
Isolated  
floor

# Summary of future neutron EDM experiments

Exp	UCN source	cell	Measurement techniques	$\sigma_d$ ( $10^{-28}$ e-cm)
ILL CryoEDM	Superfluid $^4\text{He}$	$^4\text{He}$	Ramsey technique for $\omega$ External SQUID magnetometers	$\sim 50$ $< 5$
PNPI - I LL - $\text{SD}_2$	ILL turbine PNPI/Solid $\text{D}_2$	Vac.	Ramsey technique for $\omega$ $\vec{E}=0$ cell for magnetometer	$< 100$ $< 10$
ILL Crystal	Cold n Beam			$< 100$
PSI EDM	Solid $\text{D}_2$	Vac.	Ramsey technique for $\omega$ External Cs & $^3\text{He}$ magnetom.	$\sim 50$ $\sim 5$
SNS EDM	Superfluid $^4\text{He}$	$^4\text{He}$	$^3\text{He}$ capture for $\omega$ $^3\text{He}$ comagnetometer SQUIDS & Dressed spins	$\sim 5$
TRIUMF/JPARC	Superfluid $^4\text{He}$	Vac.	Under Development	?



# New n-EDM Sensitivity



# Summary

- Physics reach of EDM measurements is significant (even after Large Hadron Collider)
  - New sources of CP violation likely in SUSY
- A new neutron EDM experiment with two orders of magnitude improvement
  - Allows possible discovery of new sources of CP violation